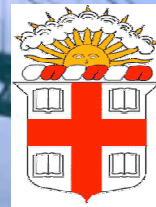


# Looking for Extra Dimensions and Black Holes with Early LHC Data

**Greg Landsberg**



**Brown University**

**LIP/CFTP Seminar**

**November 10, 2008**



# Outline

- ◆ The Hierarchy Problem
- ◆ Models with Extra Dimensions
- ◆ Gravity at Short Distances
- ◆ Astrophysical Constraints
- ◆ Collider Searches for Extra Dimensions
- ◆ Black Holes at the LHC
- ◆ Conclusions





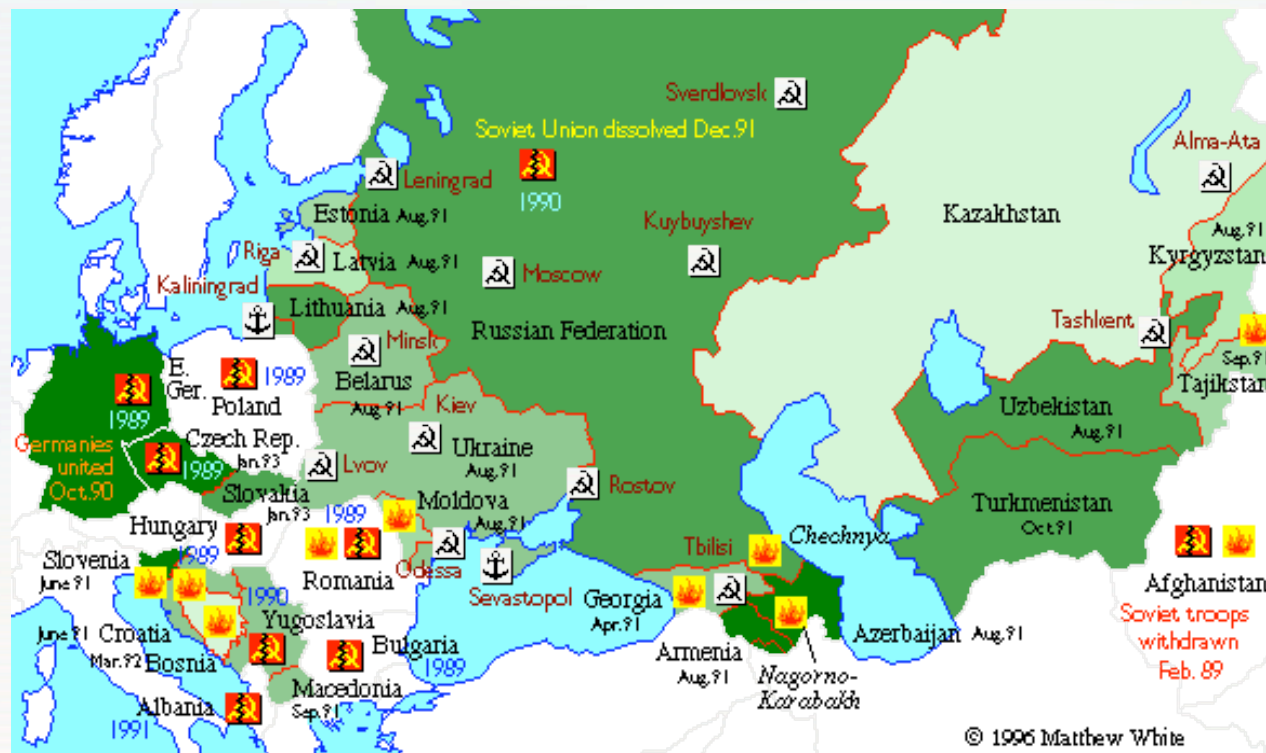
# Large Hierarchies Tend to Collapse...





# More Large Hierarchies

## Collapse of the Soviet Union



The nineties...



# Gravitational Hierarchy Collapse

With thanks to Chris Quigg and the  
B44 restaurant in San Francisco



- Human Castles in Catalonia





# But Keep in Mind...



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- **Alternative:** the *anthropic principle*
  - Properties of the universe are so special because we happen to exist and be able to ask these very questions
  - Not covered in this talk



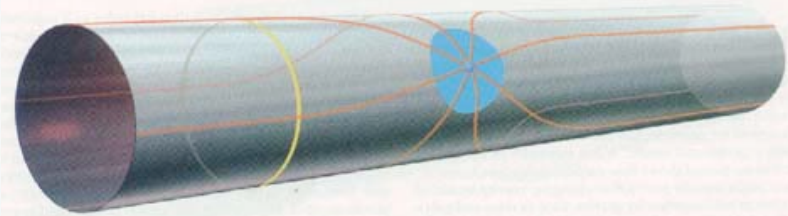
# 1998: Large Extra Dimensions

- But: **what if** there is no other scale, and SM model is correct up to  $M_{\text{Pl}}$ ?
  - Give up **naturalness**: inevitably leads to anthropic reasoning
  - Radically new approach – Arkani-Hamed, Dimopoulos, Dvali (ADD, 1998): maybe the fundamental Planck scale is only  $\sim 1 \text{ TeV}$ !!
- Gravity is made strong** at a TeV scale due to existence of **large** ( $r \sim 1 \text{ mm} - 1 \text{ fm}$ ) extra spatial dimensions:
  - SM particles are confined to a 3D “brane”
  - Gravity is the only force that permeates “bulk” space
- What about **Newton’s law**?

$$V(\rho) = \frac{1}{M_{\text{Pl}}^2} \frac{m_1 m_2}{\rho^{n+1}} \rightarrow \frac{1}{(M_{\text{Pl}}^{[3+n]})^{n+2}} \frac{m_1 m_2}{\rho^{n+1}}$$

- Ruled out for infinite ED**, but does not apply for compact ones:

$$V(\rho) \approx \frac{1}{(M_{\text{Pl}}^{[3+n]})^{n+2}} \frac{m_1 m_2}{r^n \rho}, \text{ for } \rho \gg r$$



- Gravity is fundamentally strong** force, but we do not feel that as it is diluted by the large volume of the bulk space  
 $G'_N = 1/(M_{\text{Pl}}^{[3+n]})^2 = 1/M_D^2$ ;  $M_D \sim 1 \text{ TeV}$

$$M_D^{n+2} \sim M_{\text{Pl}}^2 / r^n$$

- More precisely, from Gauss’s law:

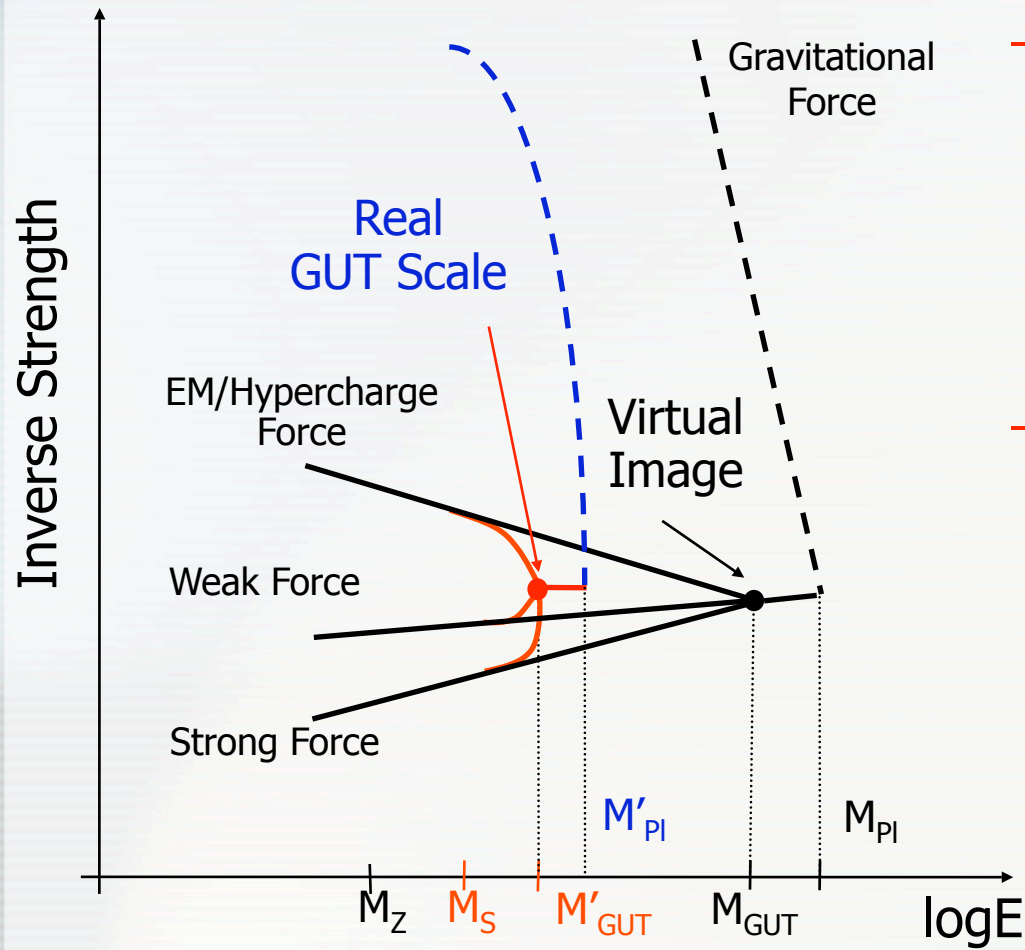
$$r = \frac{1}{\sqrt{4\pi} M_D} \left( \frac{M_{\text{Pl}}}{M_D} \right)^{2/n} \sim \begin{cases} 8 \times 10^{12} m, & n = 1 \\ 0.7 \text{ mm}, & n = 2 \\ 3 \text{ nm}, & n = 3 \\ 6 \times 10^{-12} m, & n = 4 \end{cases}$$

- Amazing as it is, but as of 1998 **no one** has tested Newton’s law to distances less than  $\sim 1 \text{ mm}$ !
- Thus, the fundamental Planck scale could be as low as 1 TeV for  $n > 1$



# 1998: $\text{TeV}^{-1}$ Extra Dimensions

- Simultaneously, another idea has appeared:

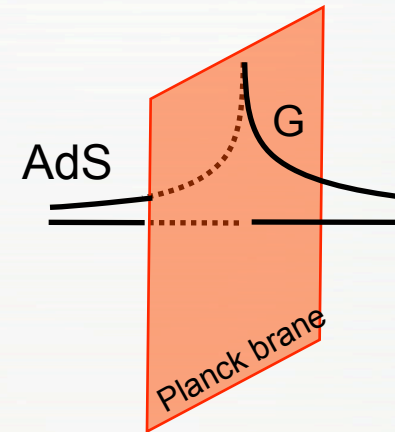


- Explore modification of force behavior in  $(3+n)$ -dimensions to achieve low-energy grand unification [Dienes, Dudas, Gherghetta, PL **B436**, 55 (1998)]
- To achieve that, allow other force carriers ( $g$ ,  $\gamma$ ,  $W$ , and  $Z$ ) to propagate in an extra dimension, which is “longitudinal” to the SM brane and compactified on a “natural” EW scale:
  - $R \sim 1 \text{ TeV}^{-1} \sim 10^{-19} \text{ m}$



# 1999: Randall-Sundrum Model

- Randall-Sundrum (RS) model [PRL **83**, 3370 (1999); PRL **83**, 4690 (1999)]
  - One + brane – no low energy effects
  - Two + and – branes – TeV Kaluza-Klein modes of graviton
  - Low energy effects on SM brane are given by  $\Lambda_\pi$ ; for  $kr \sim 10$ ,  $\Lambda_\pi \sim 1$  TeV and the hierarchy problem is solved naturally

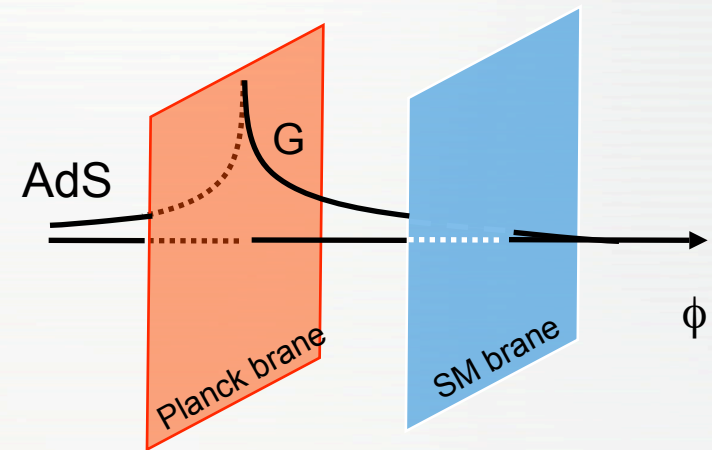






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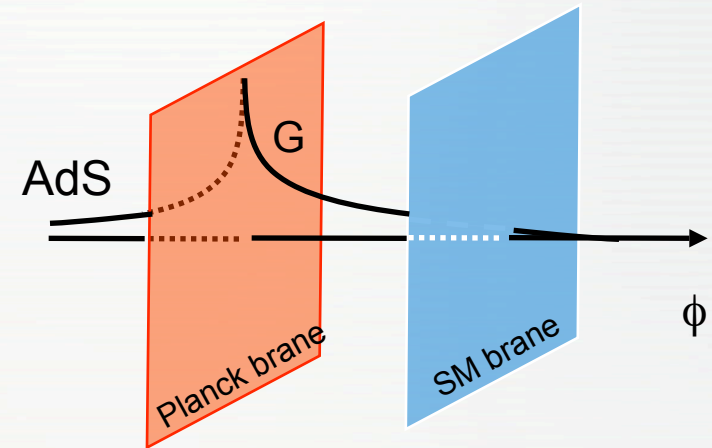
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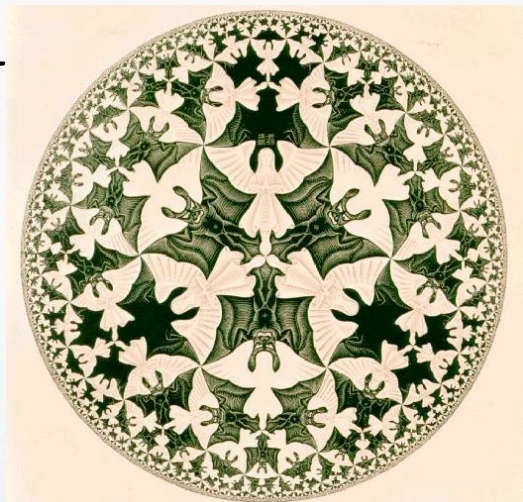
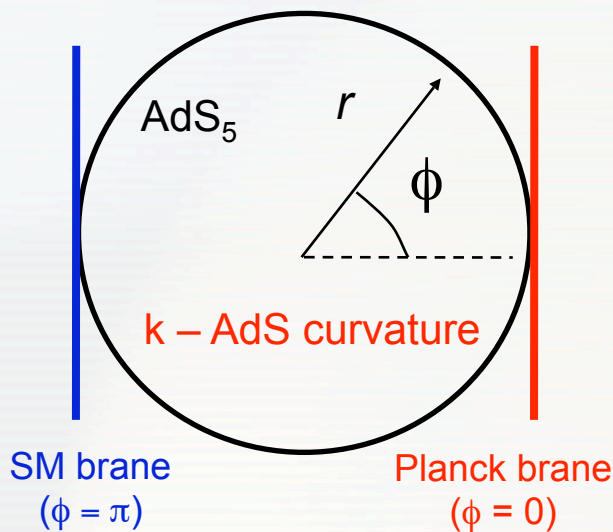
Anti-deSitter space-time metric:

$$ds^2 = e^{-2kr|\phi|} \eta_{\mu\nu} dx^\mu dx^\nu - r^2 d\phi^2$$

$$\Lambda_\pi = \overline{M}_{Pl} e^{-kr\pi}$$

Reduced Planck mass:

$$\overline{M}_{Pl} \equiv M_{Pl} / \sqrt{8\pi}$$

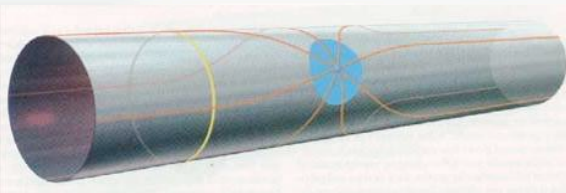




# Extra Dimensions: a Brief Recap

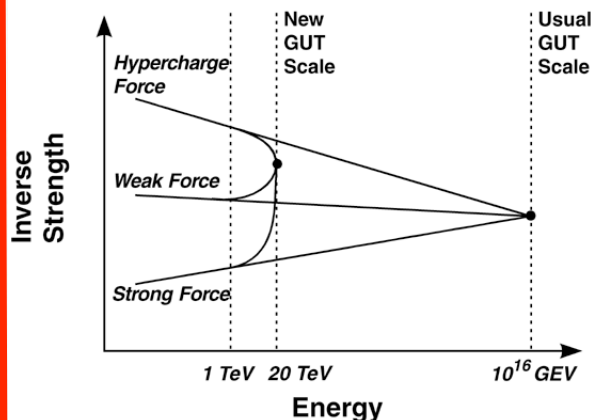
## ADD Paradigm:

- Pro: “Eliminates” the hierarchy problem by stating that physics ends at a TeV scale
- Only gravity lives in the “bulk” space
- Size of ED’s ( $n=2-7$ ) between  $\sim 100 \mu\text{m}$  and  $\sim 1 \text{ fm}$
- Black holes at the LHC and in the UHE cosmic rays
- Con: Doesn’t explain why ED are so large



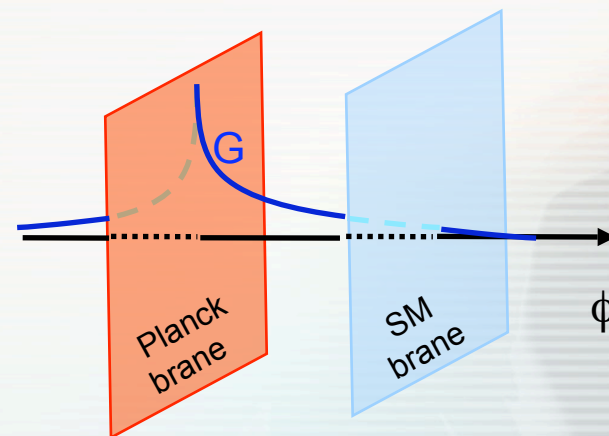
## TeV<sup>-1</sup> Scenario:

- Pro: Lowers GUT scale by changing the running of couplings
- Only gauge bosons ( $g/\gamma/W/Z$ ) “live” in ED’s
- Size of ED’s  $\sim 1 \text{ TeV}^{-1}$  or  $\sim 10^{-19} \text{ m}$  – i.e., natural EWSB size
- Con: Gravity is not in the picture



## RS Model:

- Pro: A rigorous solution to the hierarchy problem via localization of gravity
- Gravitons (and possibly other particles) propagate in a single ED, with special metric
- Black holes at the LHC and in UHE cosmic rays
- Con: Somewhat disfavored by precision EW fits

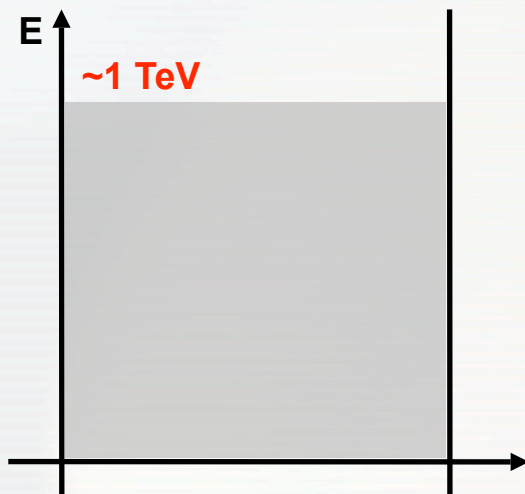




# ED: Kaluza-Klein Spectrum

## ADD Paradigm:

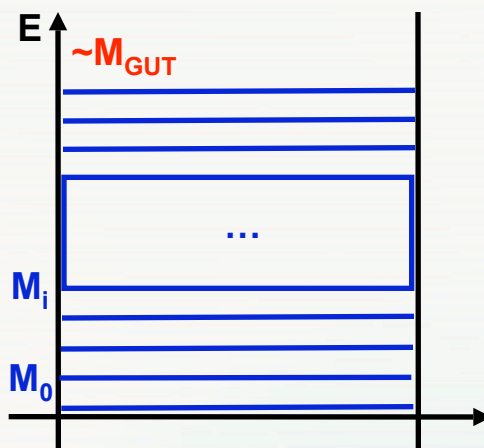
- Winding modes with energy spacing  $\sim 1/r$ , i.e. 1 meV – 100 MeV
- Experimentally can't resolve these modes – they appear as continuous spectrum
- Coupling:  $G_N$  per mode; compensated by large number of modes



## TeV<sup>-1</sup> Scenario:

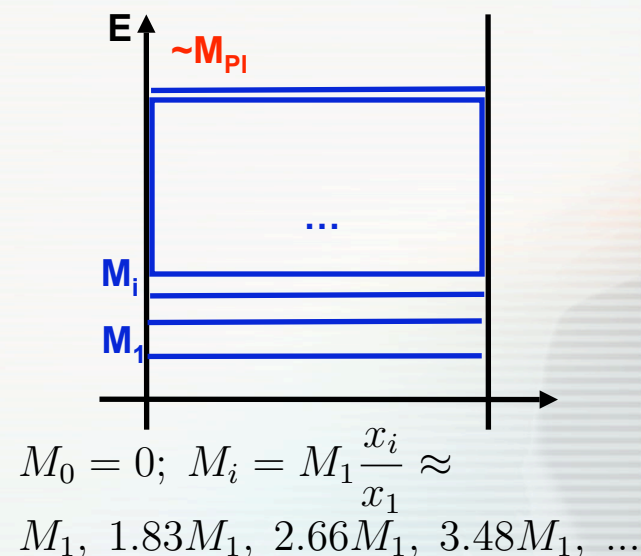
- Winding modes with nearly equal energy spacing  $\sim 1/r$ , i.e.  $\sim 1$  TeV
- Can excite individual modes at colliders or look for indirect effects
- Coupling:  $\sim g_w$  per mode

$$M_i = \sqrt{M_0^2 + i^2/r^2}$$



## RS Model:

- “Particle in a box” with special AdS metric
- Energy eigenvalues are given by the zeroes of Bessel function  $J_1$
- Light modes might be accessible at colliders
- Coupling:  $G_N$  for the zero mode;  $1/\Lambda_\pi^2$  for the others

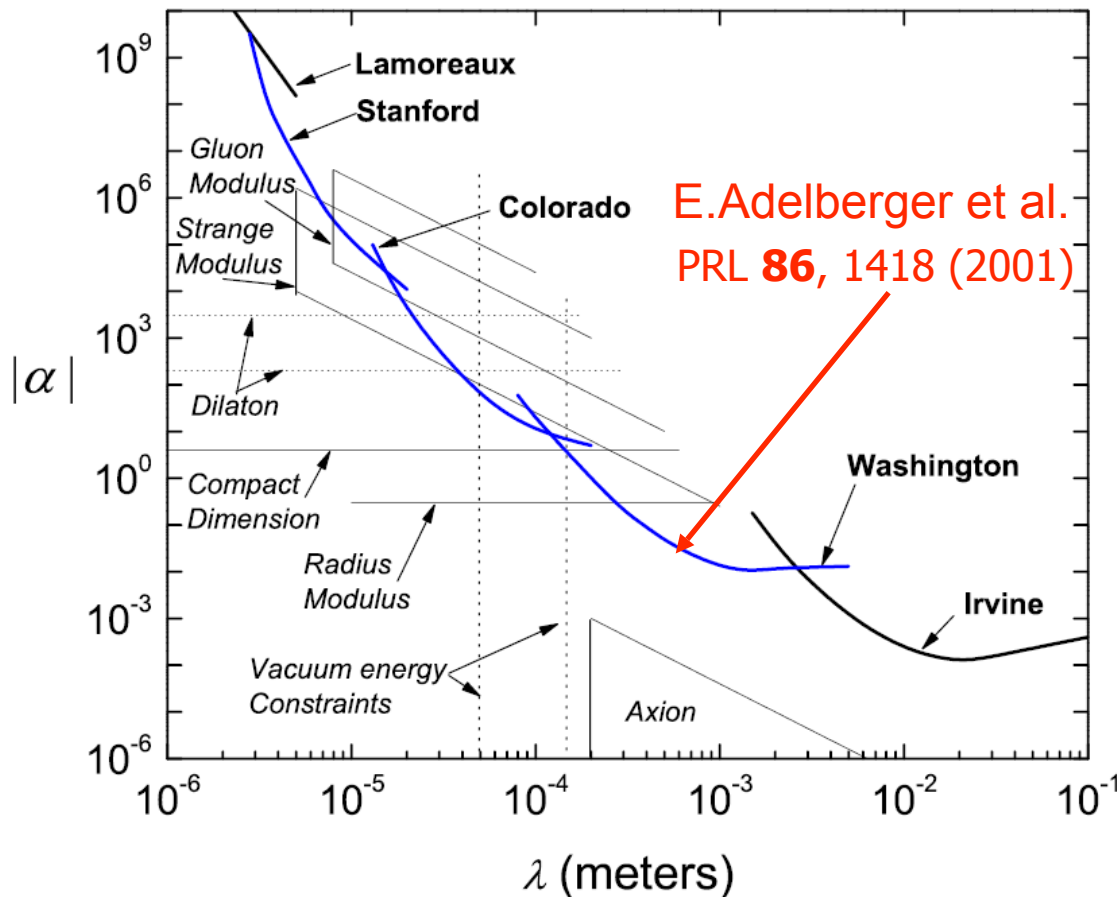






# Large ED: Gravity at Short Distances

[J. Long, J. Price, hep-ph/0303057]

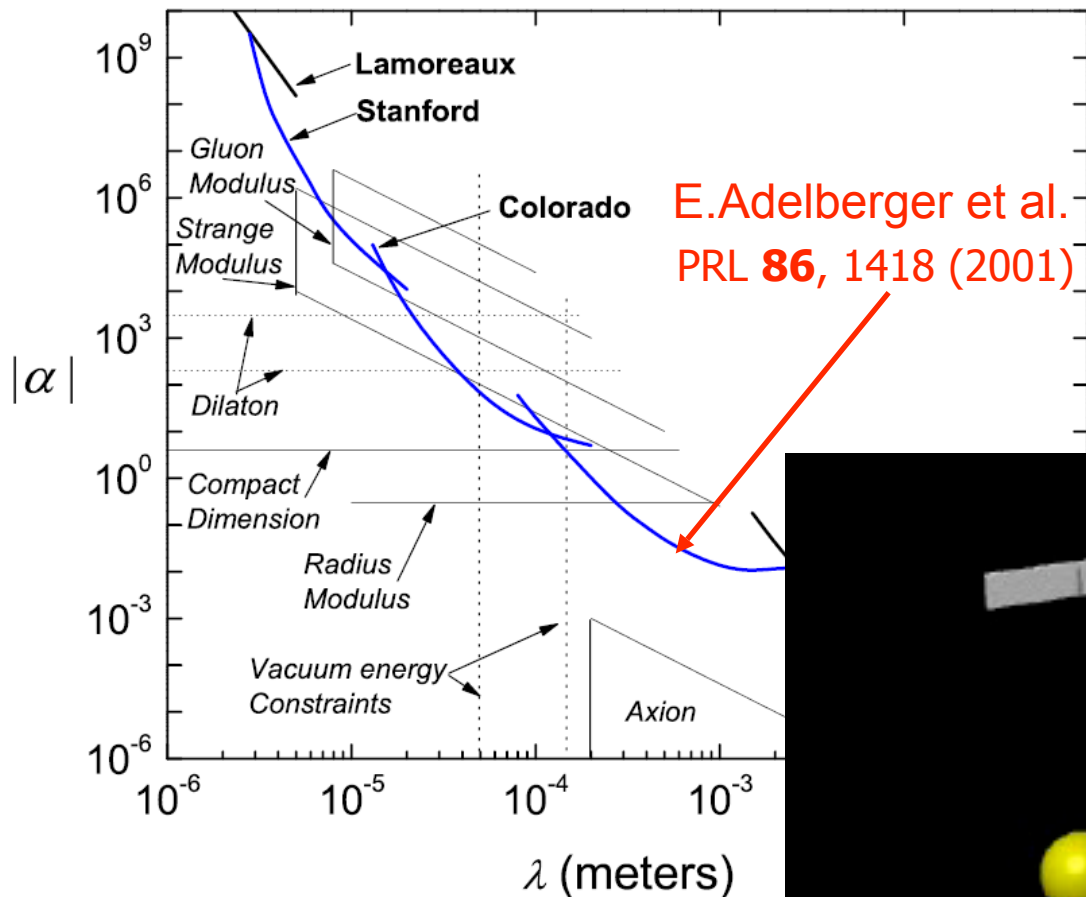


- Sub-millimeter gravity measurements could probe *only*  $n=2$  case *only* within the ADD model
  - The best sensitivity so far have been achieved in the U of Washington torsion balance experiment – a high-tech “remake” of the 1798 Cavendish experiment
    - $R \lesssim 0.16$  mm ( $M_D \gtrsim 1.7$  TeV)
- Sensitivity vanishes quickly with the distance – can’t push limits further down
  - Started restricting ADD with 2 extra dimensions; can’t probe any higher number
  - Ultimately push the sensitivity by a factor of two in terms of the distance
- No sensitivity to the  $\text{TeV}^{-1}$  and RS models



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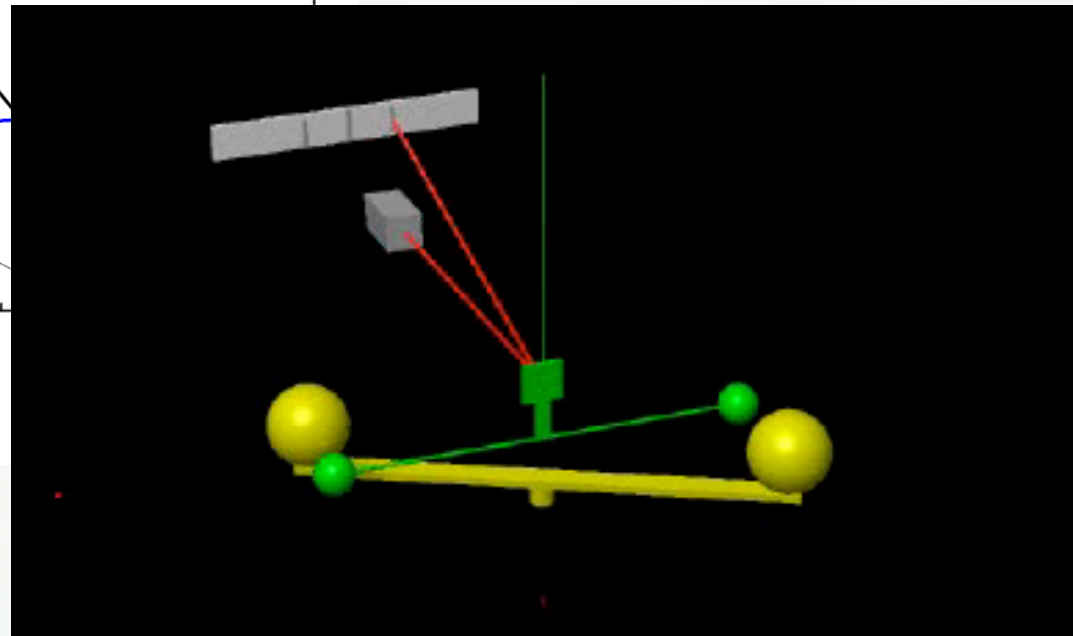
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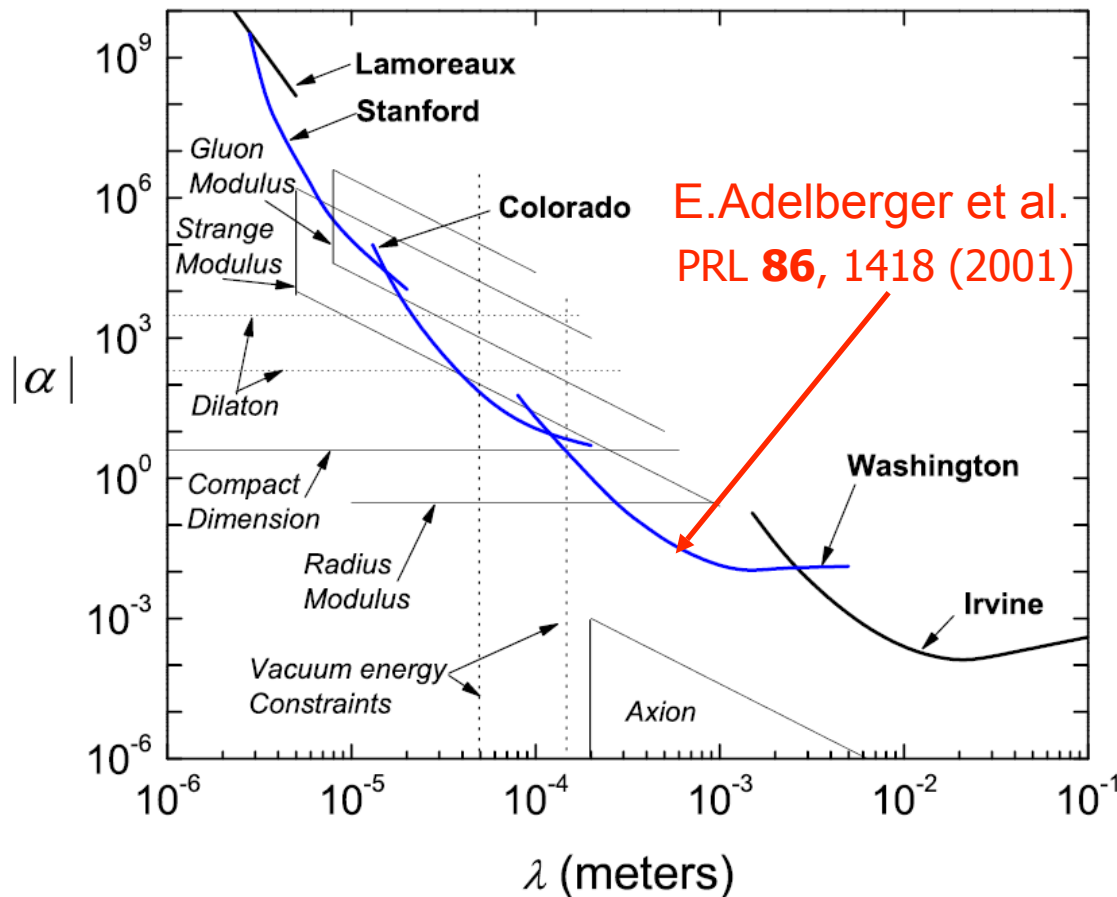
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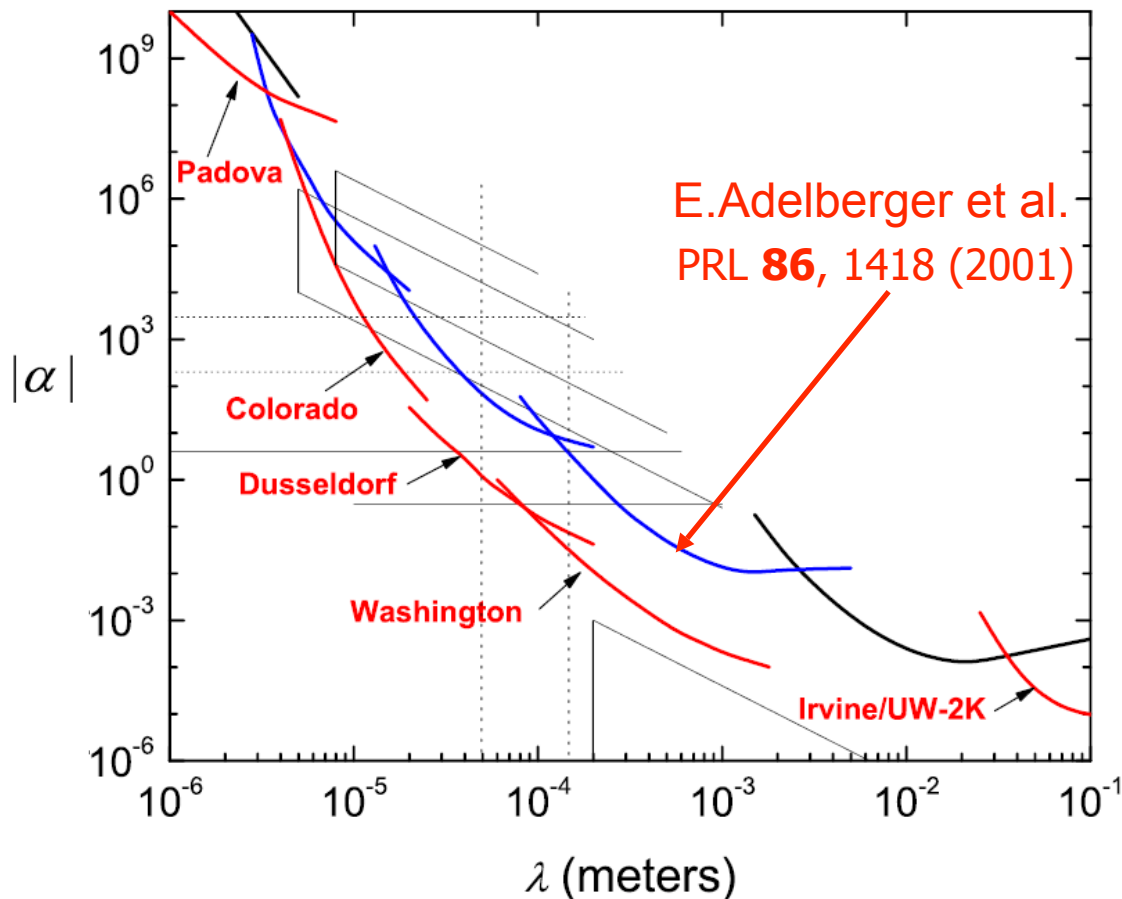


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# Large ED: Astro & Cosmo Constraints

- Supernova cooling due to graviton emission – an alternative cooling mechanism that would decrease the dominant one via neutrino emission
  - Tightest limits on any additional cooling sources come from the measurement of the SN1987A neutrino flux by Kamiokande and IMB
  - Application to the ADD scenario: Cullen and Perelstein [PRL **83**, 268 (1999)]; Hanhart, Phillips, Reddy, and Savage [Nucl. Phys. **B595**, 335 (2001)]:
    - $M_D > 25\text{-}30 \text{ TeV}$  ( $n=2$ )
    - $M_D > 2\text{-}4 \text{ TeV}$  ( $n=3$ )
- Distortion of the cosmic diffuse gamma radiation (CDG) spectrum due to the  $G_{KK} \rightarrow \gamma\gamma$  decays: Hall and Smith [PRD **60**, 085008 (1999)]:
  - $M_D > 100 \text{ TeV}$  ( $n=2$ )
  - $M_D > 5 \text{ TeV}$  ( $n=3$ )
- Overclosure of the universe, matter dominance in the early universe, Fairbairn [Phys. Lett. **B508**, 335 (2001)]; Fairbairn, Griffiths [JHEP 0202, **024** (2002)]:
  - $M_D > 86 \text{ TeV}$  ( $n=2$ )
  - $M_D > 7.4 \text{ TeV}$  ( $n=3$ )
- Neutron star  $\gamma$ -emission from radiative decays of the gravitons trapped during the supernova collapse, Hannestad and Raffelt [PRL **88**, 071301 (2002)]:
  - $M_D > 1700 \text{ TeV}$  ( $n=2$ )
  - $M_D > 60 \text{ TeV}$  ( $n=3$ )
- Caveat: there are many known (and unknown!) uncertainties, so the cosmological bounds are reliable only as an order of magnitude estimate
- Still,  $n=2$  is largely disfavored



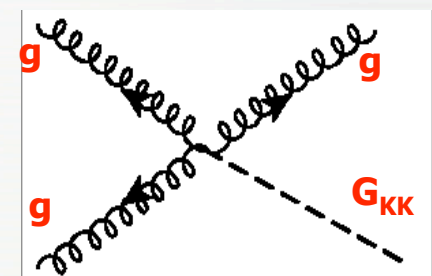
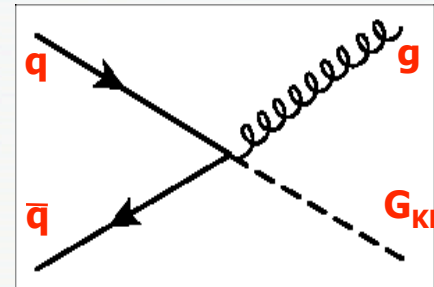


# Collider Signatures for Large ED

- Kaluza-Klein gravitons couple to the energy-momentum tensor, and therefore contribute to most of the SM processes
- For Feynman rules for  $G_{KK}$  see:
  - Han, Lykken, Zhang [PRD **59**, 105006 (1999)]
  - Giudice, Rattazzi, Wells [NP **B544**, 3 (1999)]
- Graviton emission: direct sensitivity to the fundamental Planck scale  $M_D$
- Virtual effects: sensitive to the ultraviolet cutoff  $M_S$ , expected to be  $\sim M_D$  (and likely  $< M_D$ )
- The two processes are complementary

## Real Graviton Emission

Monojets at hadron colliders



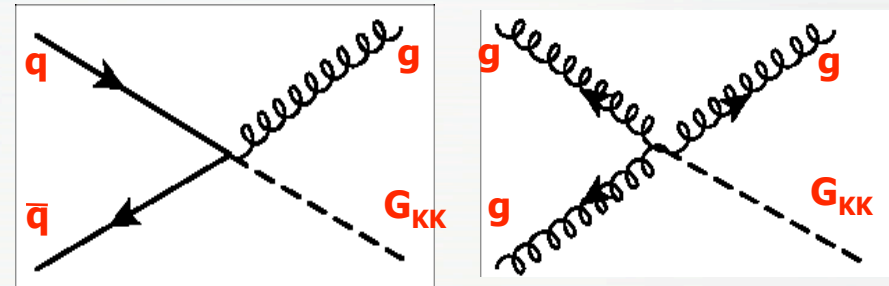


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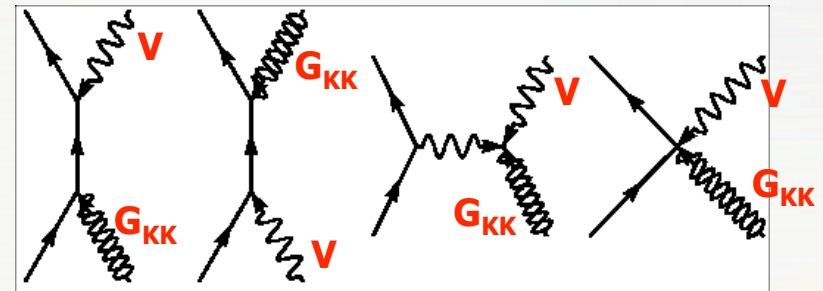
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Single VB at hadron or  $e^+e^-$  colliders



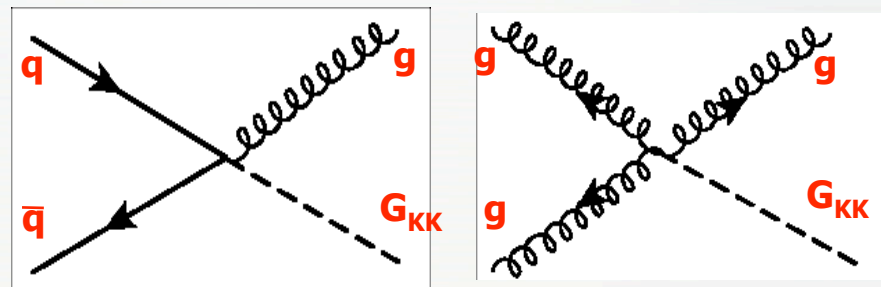


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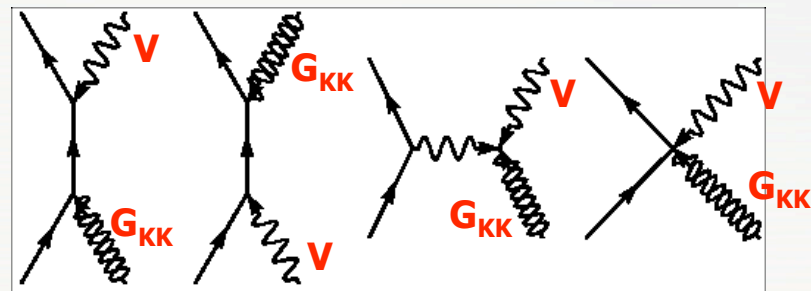
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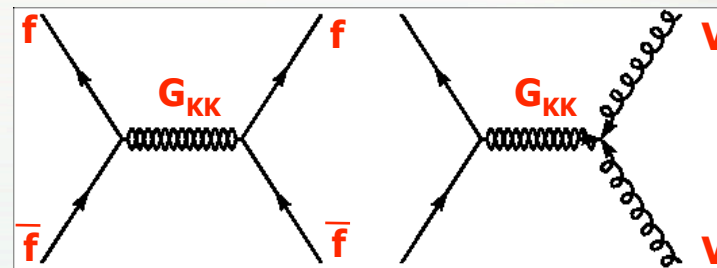


Single VB at hadron or  $e^+e^-$  colliders



## Virtual Graviton Effects

Fermion or VB pairs at hadron or  $e^+e^-$  colliders





# L'EPilogue (Large ED)

## Direct Graviton Emission

	$e^+e^- \rightarrow \gamma G$					$e^+e^- \rightarrow ZG$					
Experiment	n=2	n=3	n=4	n=5	n=6	n=2	n=3	n=4	n=5	n=6	Color coding
ALEPH	1.28	0.97	0.78	0.66	0.57	0.35	0.22	0.17	0.14	0.12	$\leq 184$ GeV
DELPHI	1.38	<b>1.02</b>	0.84	<b>0.68</b>	0.58						$\leq 189$ GeV
L3	1.02	0.81	0.67	0.58	0.51	0.60	0.38	0.29	<b>0.24</b>	<b>0.21</b>	$> 200$ GeV
OPAL	1.09	0.86	0.71	0.61	0.53						$\lambda = -1$ $\lambda = +1$ GL

All limits are in TeV

## Virtual Graviton Exchange

Experiment	$e^+e^-$	$\mu^+\mu^-$	$\tau^+\tau^-$	$qq$	$ff$	$\gamma\gamma$	$WW$	$ZZ$	Combined
ALEPH	1.04 0.81	0.65 0.67	0.60 0.62	0.53/0.57 0.46/0.46 (bb)	1.05 0.84	0.81 0.82			0.75/1.00 (<189)
DELPHI		0.59 0.73	0.56 0.65		0.60 0.76	0.83 0.91			0.60/0.76 (ff) (<202)
L3	0.98 1.06	0.56 0.69	0.58 0.54	0.49 0.49	0.84 1.00	0.99 0.84	0.68 0.79		1.0/1.1 (<202)
OPAL	1.15 1.00	0.62 0.66			0.62 0.66	0.89 0.83		0.63 0.74	1.17/1.03 (<209)

LEP Combined: 1.2/1.1 TeV



# Monojets: Tainted History

EXPERIMENTAL OBSERVATION OF EVENTS WITH LARGE MISSING TRANSVERSE ENERGY  
ACCOMPANIED BY A JET OR A PHOTON(S) IN  $p\bar{p}$  COLLISIONS

AT  $\sqrt{s} = 540$  GeV

**[PL, 139B, 115 (1984)]**

UA1 Collaboration, CERN, Geneva, Switzerland

## Abstract

We report the observation of five events in which a missing transverse energy larger than 40 GeV is associated with a narrow hadronic jet and of two similar events with a neutral electromagnetic cluster (either one or more closely spaced photons). We cannot find an explanation for such events in terms of backgrounds or within the expectations of the Standard Model.







# Monojets: Tainted History

EXPERIMENTAL OBSERVATION OF EVENTS WITH LARGE MISSING TRANSVERSE ENERGY  
ACCOMPANIED BY A JET OR A PHOTON(S) IN  $p\bar{p}$  COLLISIONS

AT  $\sqrt{s} = 540$  GeV

**[PL, 139B, 115 (1984)]**

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VOLUME 54, NUMBER 6

PHYSICAL REVIEW LETTERS

11 FEBRUARY 1985

## **Monojets from Z Decay without Extra Neutrinos or Higgs Particles**

Stephen F. King

*Lyman Laboratory of Physics, Harvard University, Cambridge, Massachusetts 02138*

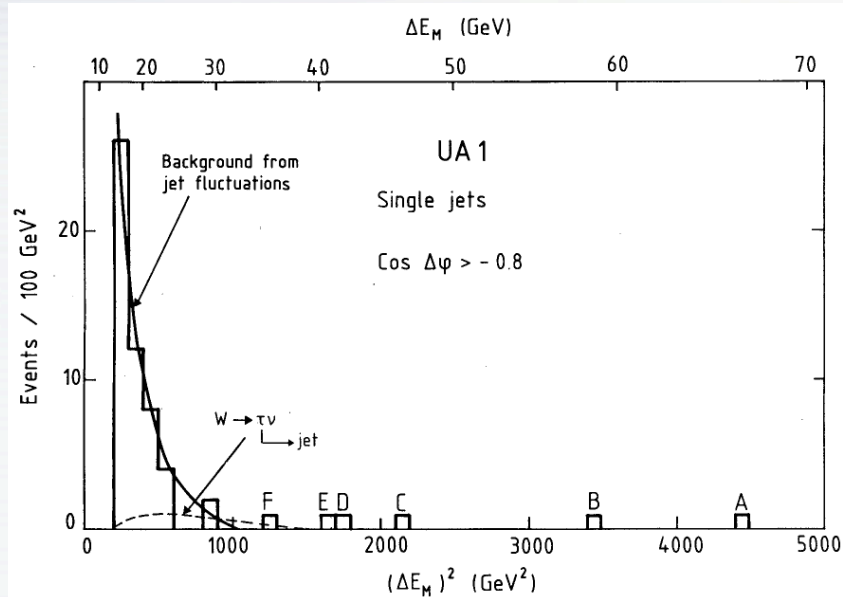
(Received 26 November 1984)

The recent discovery of monojets by Arnison *et al.*<sup>1</sup> at the CERN  $p\bar{p}$  collider has caused ripples of excitement throughout the particle physics world, since they cannot be explained by the minimal standard model.<sup>2</sup>

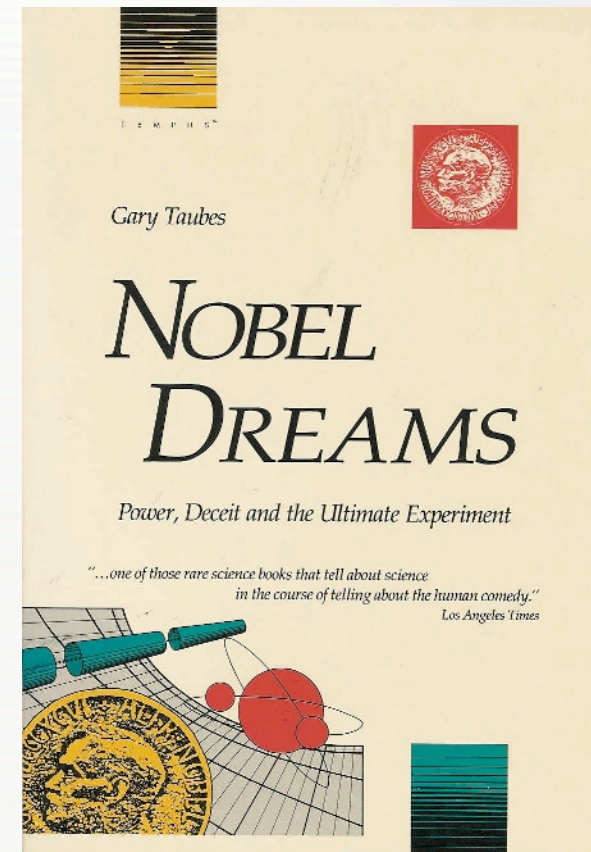




# Monojets: Tainted History



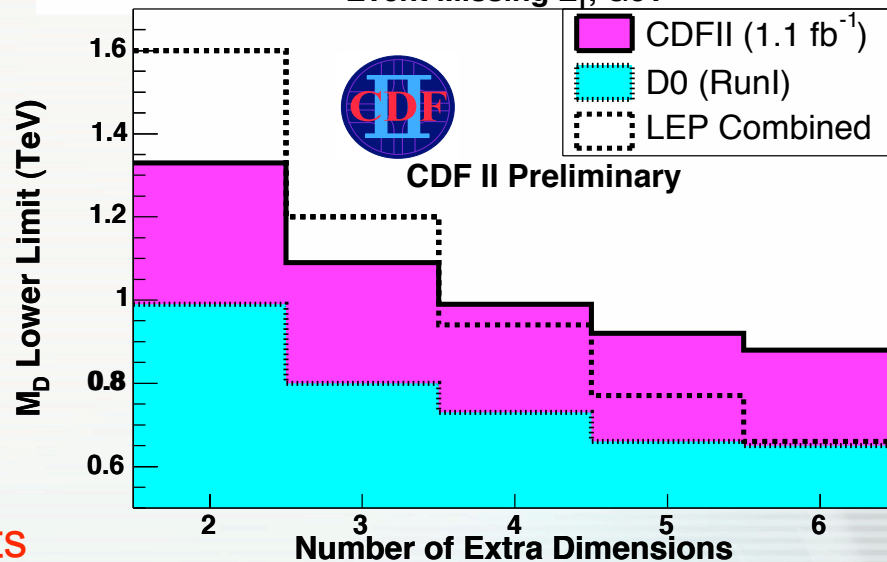
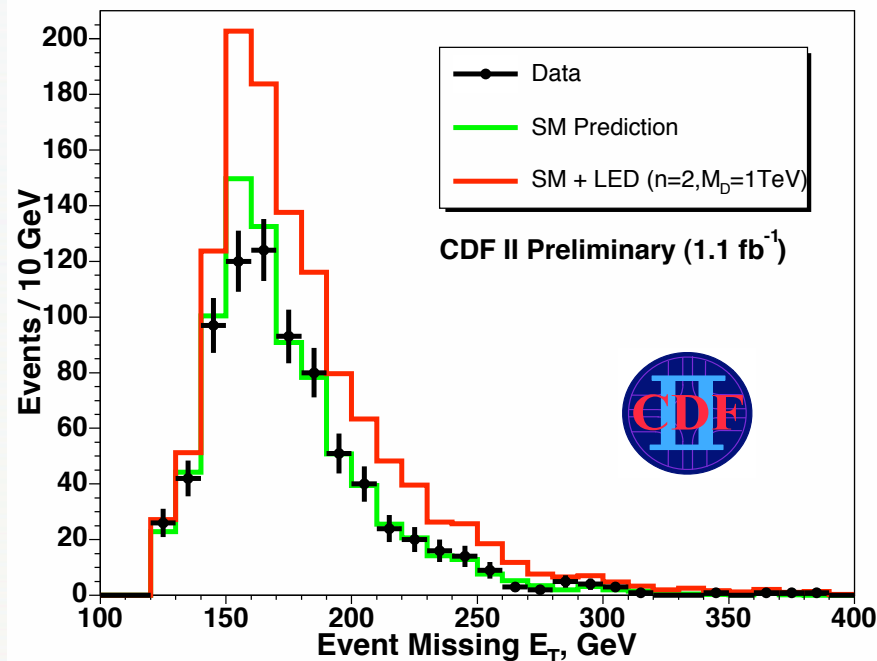
- These **monojets** turned out to be due to **unaccounted background**
- The **signature** was deemed **doomed** and nearly forgotten
- It **took many years** for **successful monojet analyses** at a hadron collider to be completed (CDF/DØ)





# Tevatron: Large ED Search via Monojets

- jets +  $ME_T$  final state
- $Z(\nu\nu)+\text{jets}$  is irreducible background
  - Challenging signature due to large instrumental backgrounds from jet mismeasurement, cosmics, etc.
- DØ pioneered this search and set limits [PRL, **90** 251802 (2003)]  
 $M_P > 1.0-0.6$  TeV for  $n=2\dots 7$
- CDF analysis based on  $1.1 \text{ fb}^{-1}$ 
  - Central jet w/  $E_T > 150 \text{ GeV}$
  - $ME_T > 120 \text{ GeV}$
  - No other jets w/  $E_T > 60 \text{ GeV}$
  - 779 events observed with  $819 \pm 71$  expected (half comes from  $Z(\nu\nu)+j$ )
  - Set limits on the fundamental Planck scale between 0.88 and 1.33 TeV
  - Similar results with looser  $ME_T$ ,  $E_T^j$  cuts

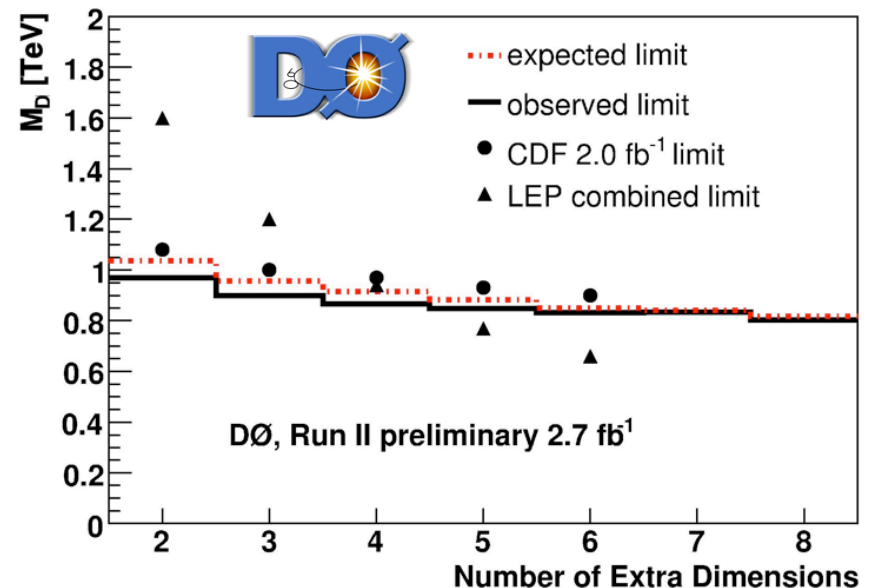
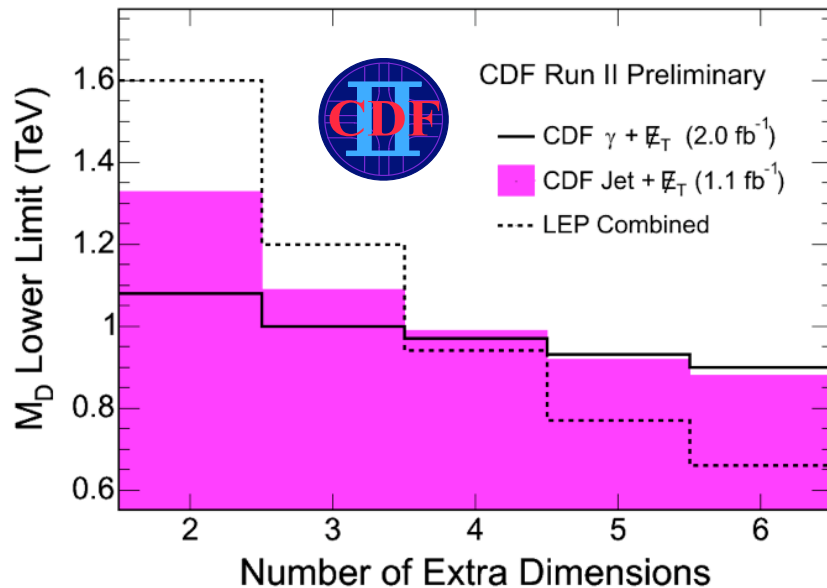






# Tevatron Searches for ED in Monophotons

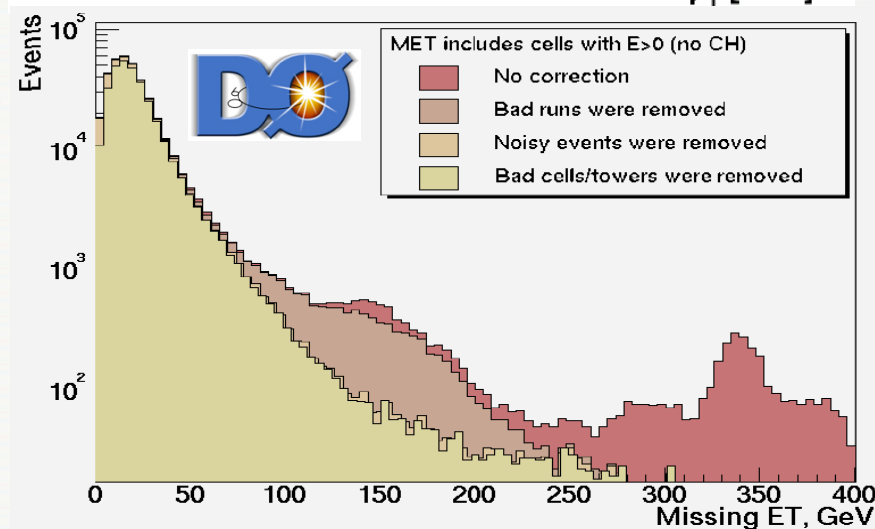
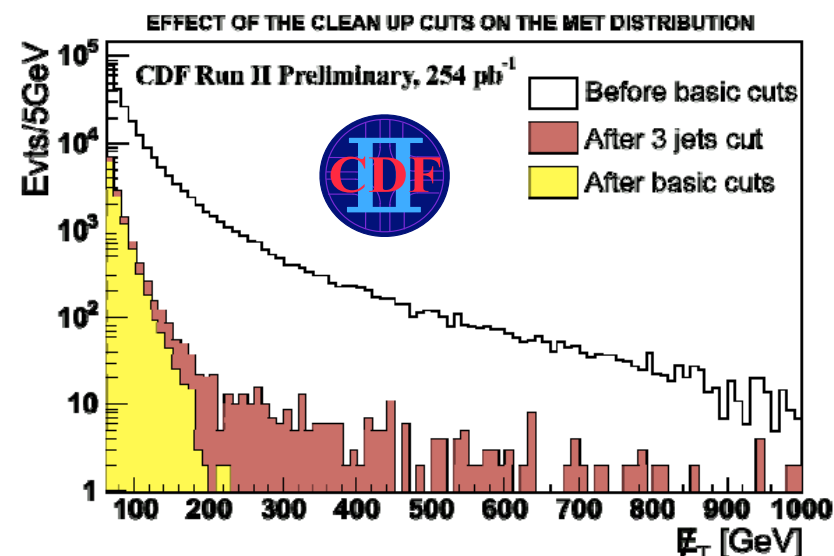
- Both CDF and DØ completed monophoton searches
- While easier than the monojet one, the sensitivity is typically not as good, especially for low number of ED
  - CDF monophoton limits approach monojet ones at large  $n$ , but require twice the luminosity





# Why $ME_T$ is Tough?

- Fake  $ME_T$  appears naturally in multijet events, which have enormous rate at the LHC
- Jets tend to fluctuate wildly:
  - Large shower fluctuation
  - Fluctuations in the e/h energy ratio
  - Non-linear calorimeter response
  - Non-compensation (i.e.,  $e/h \neq 1$ )
- Instrumental effects:
  - Dead or “hot” calorimeter cells
  - Cosmic ray bremsstrahlung
  - Poorly instrumented area of the detector
- Consequently, it will be a challenge to use in early LHC running
- Nevertheless,  $ME_T$  is one of the most prominent signatures for new physics and thus must be pursued

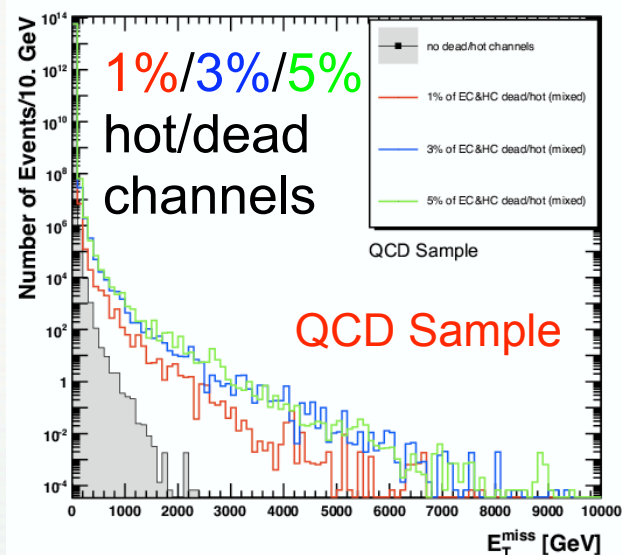
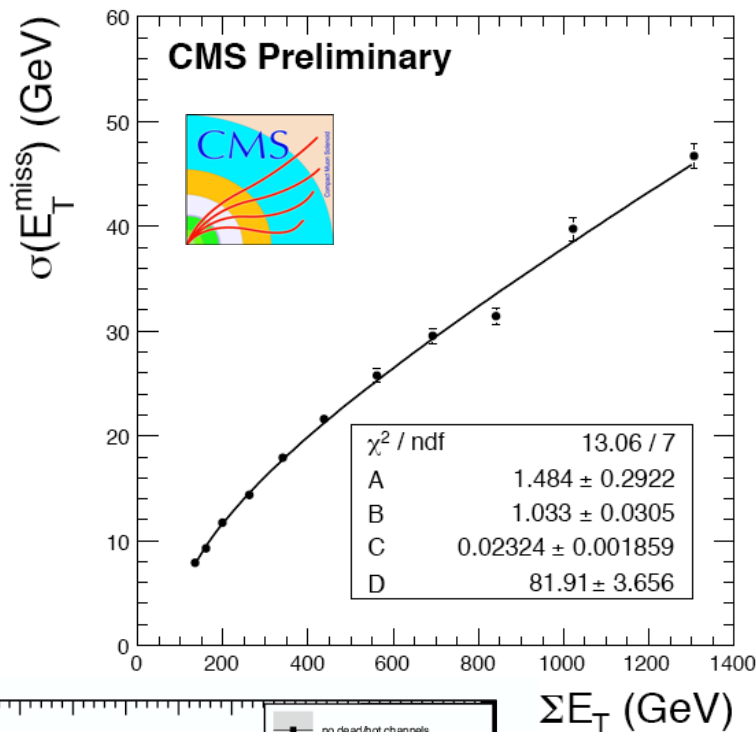


- Raw  $ME_T$  spectrum at the Tevatron and that after thorough clean-up



# ME<sub>T</sub> in CMS

- Parameters:
  - $A = 1.48$  GeV (noise term)
  - $B = 1.03$  GeV<sup>1/2</sup> (sampling term)
  - $C = 0.023$  (constant term; dominates at large  $S_T$ )
  - $D = 82$  GeV ( $\Sigma E_T$  with no beam)
- Apart from the resolution an important characteristic is the non-Gaussian tails
- Better performance at the low end is expected from particle flow
- Very hard to simulate; will have to wait for real data to see how large the effect is
  - A few special cases have been looked at already, e.g. the effect of hot/dead channels

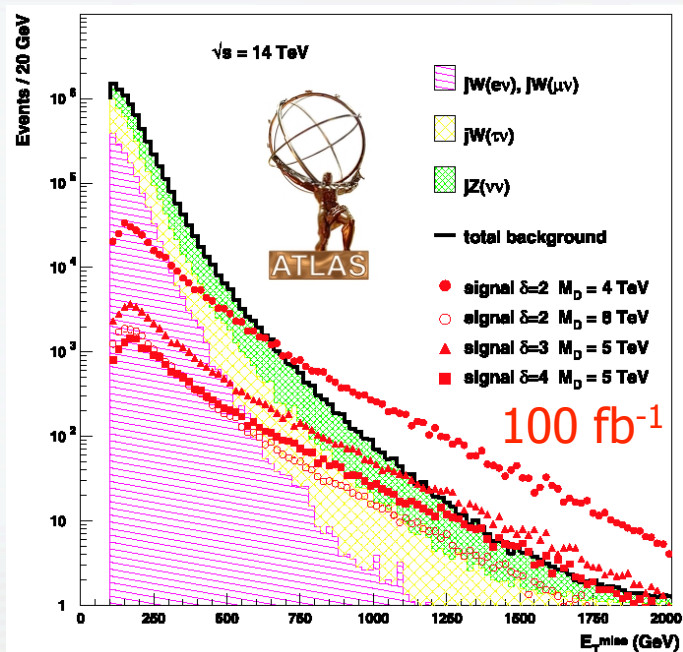




# Expectations at the LHC

## • Monojets:

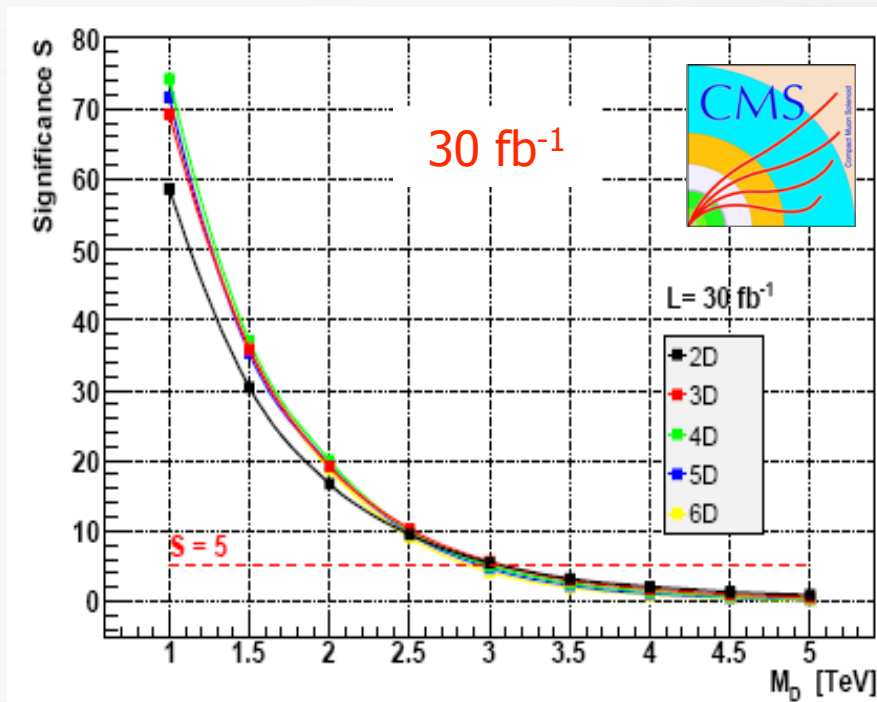
- ATLAS fast simulation for 30 and 100 fb<sup>-1</sup> (caveat: no instrumental bckg. included)



$\delta$	$M_D^{max}$ (TeV) LL, 30 fb <sup>-1</sup>	$M_D^{max}$ (TeV) HL, 100 fb <sup>-1</sup>	$M_D^{min}$ (TeV)
2	7.7	9.1	~ 4
3	6.2	7.0	~ 4.5
4	5.2	6.0	~ 5

## • Monophotons:

- ATLAS and CMS simulations for 100 fb<sup>-1</sup> and 30 fb<sup>-1</sup>, respectively



$\delta$	$M_D^{max}$ (TeV) HL, 100 fb <sup>-1</sup>	$M_D^{min}$ (TeV)
2	4	~ 3.5



# New CMS Monojet Analysis

- Jet  $E_T > 350$  GeV ( $|\eta| < 1.7$ ) and  $ME_T > 400$  GeV
- Second jet veto ( $E_T < 40$  GeV,  $|\eta| < 3$ )
- Dominated by the irreducible background (determined from W+jets)

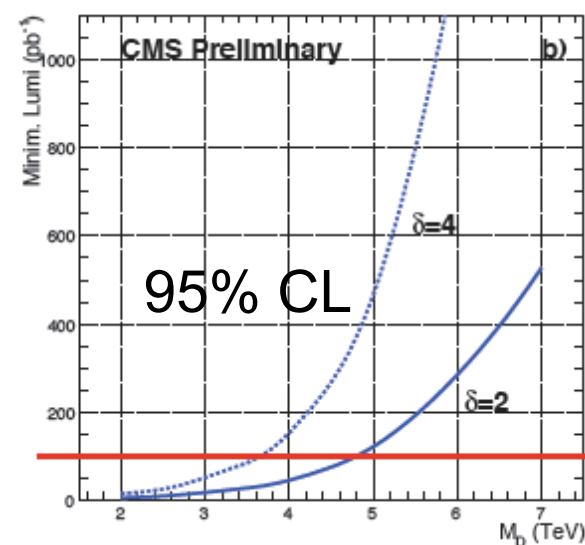
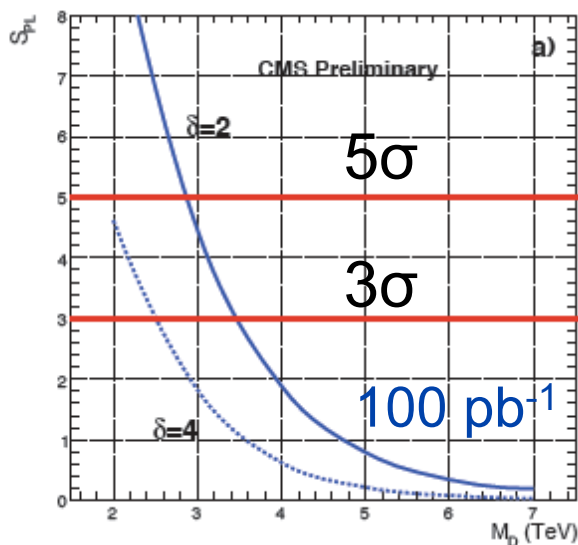
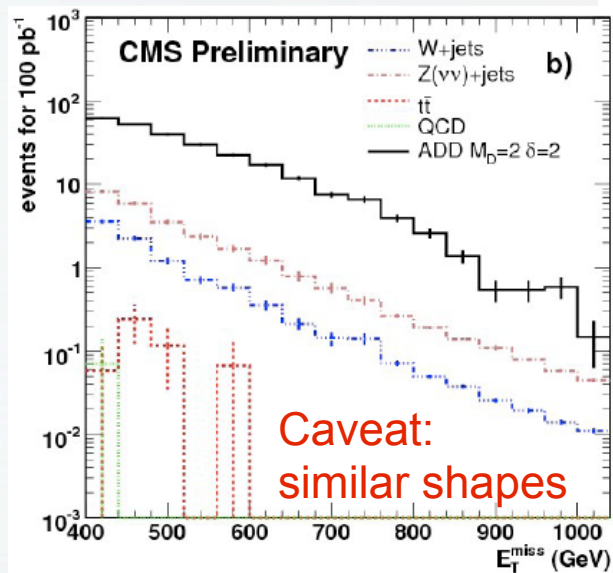


Table 1: Most recent 95% CL lower limits on the fundamental Planck scale  $M_D$  (in TeV).

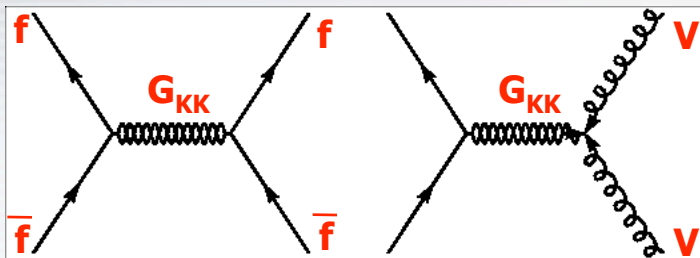
Experiment and channel	$n = 2$	$n = 3$	$n = 4$	$n = 5$	$n = 6$
LEP Combined [12]	1.60	1.20	0.94	0.77	0.66
CDF monophotons, 2.0 fb <sup>-1</sup> [18]	1.08	1.00	0.97	0.93	0.90
DØ monophotons, 2.7 fb <sup>-1</sup> [19]	0.97	0.90	0.87	0.85	0.83
CDF monojets, 1.1 fb <sup>-1</sup> [20]	1.31	1.08	0.98	0.91	0.88
CDF combined [18]	1.42	1.16	1.06	0.99	0.95

CMS reach w/ 100 pb<sup>-1</sup>  
 95% CL exclusion:  
 n=2 - 4.8 TeV  
 n=4 - 3.6 TeV  
 5σ observation:  
 n=2 - 2.8 TeV  
 n=4 - 1.8 TeV



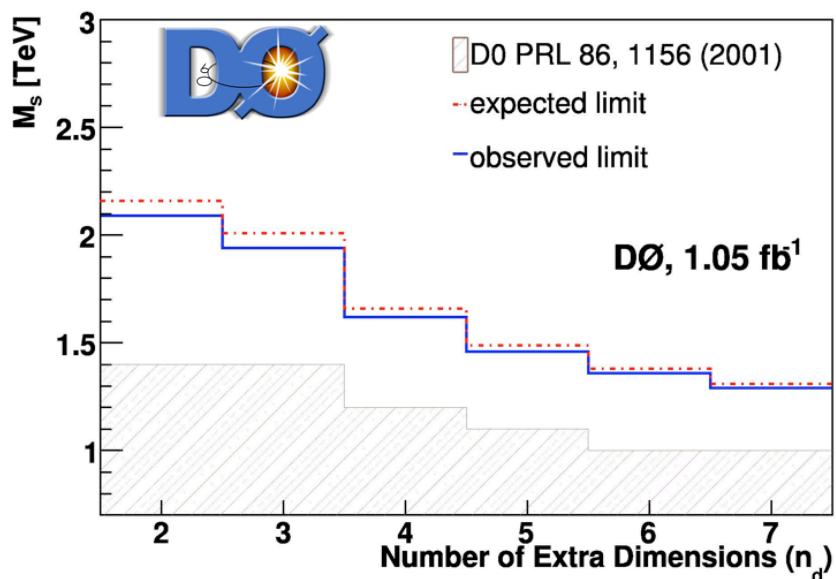


# Tevatron: Virtual Graviton Effects



- Expect an interference with the SM fermion or boson pair production

$$\frac{d^2\sigma}{d\cos\theta^*dM} = \frac{d^2\sigma_{\text{SM}}}{d\cos\theta^*dM} + \frac{a(n)}{M_S^4} f_1(\cos\theta^*, M) + \frac{b(n)}{M_S^8} f_2(\cos\theta^*, M)$$



- High-mass, low  $|\cos\theta^*|$  tail is a characteristic signature of LED  
Cheung, GL [PRD **62** 076003 (2000)]
- Best limits on the effective Planck scale come from 1 fb<sup>-1</sup> DØ Run II data:
  - $M_S > 1.3\text{--}2.1$  TeV ( $n=2\text{--}7$ ) - tightest to date
- Recent results from dijets yield similar sensitivity

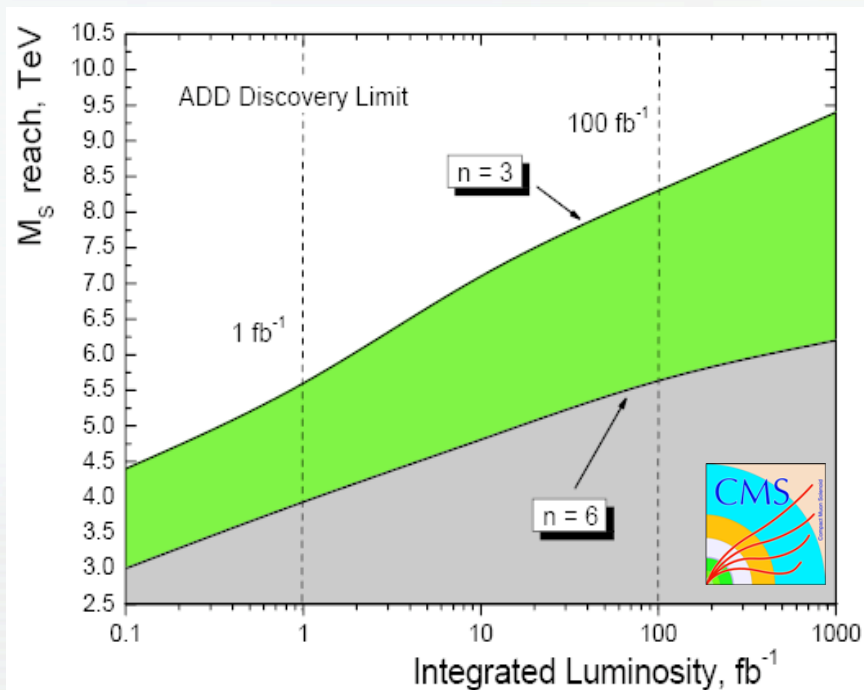
DØ Signature	GRW [2]		HLZ [11]				
		$n=2$	$n=3$	$n=4$	$n=5$	$n=6$	$n=7$
$ee + \gamma\gamma$ , 1.1 fb <sup>-1</sup> [21]	1.62	2.09	1.94	1.62	1.46	1.36	1.29
Dijets, 0.7 fb <sup>-1</sup> [22]	1.56		1.85	1.56	1.41	1.31	1.24



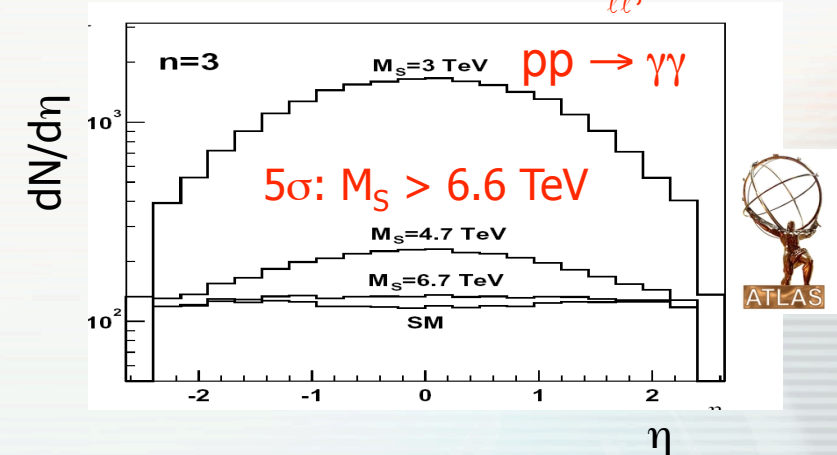
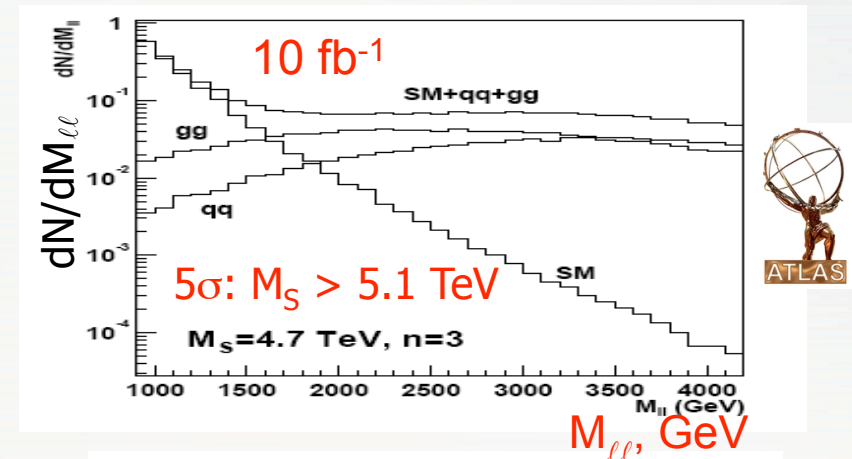


# Virtual Graviton Effects at the LHC

- Clean signature, with a huge potential of a quick discovery in dimuon, dielectron, and diphoton channels:
  - Factor of  $\sim 3$  gain over the Tevatron/Cosmic Ray limits in just  $100 \text{ pb}^{-1}$
  - Will also probe generic compositeness models with similar increase in sensitivity compared to the existing limits



CMS reach for large ED in the dimuon channel

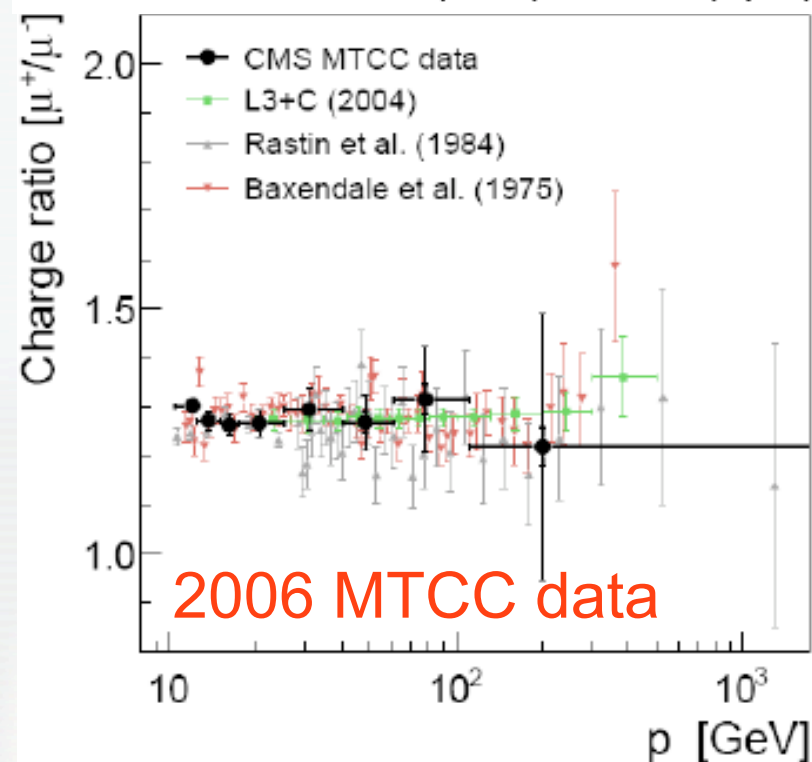




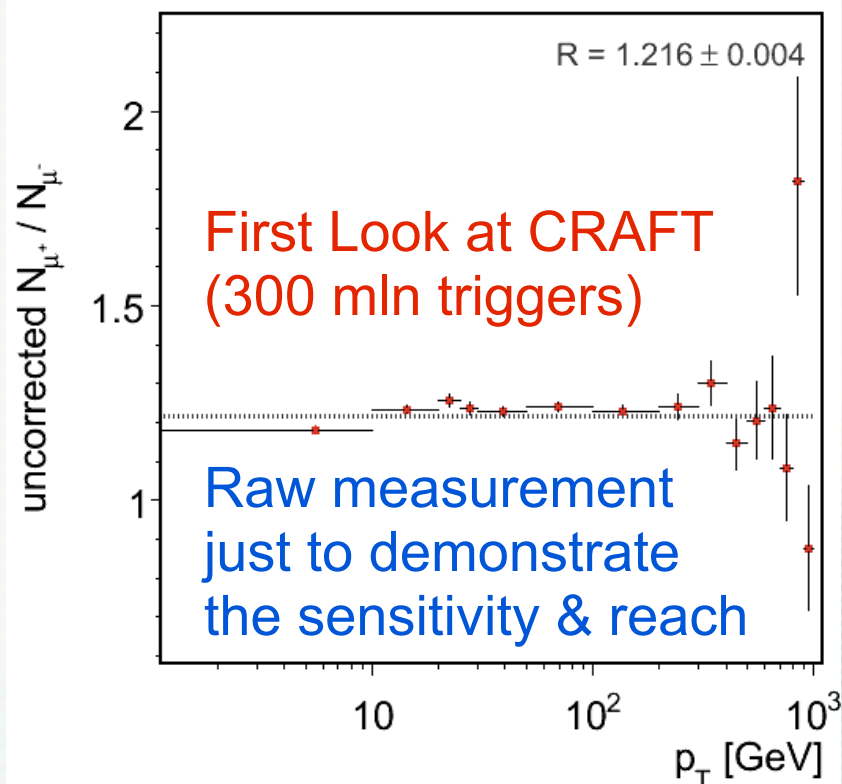
# First High- $p_T$ Muons in CMS

- While waiting for the collision data, CMS is already looking for high- $p_T$  muons from cosmic rays
- Charge ratio for atmospheric muons agrees with other measurements and approaches their precision

$$\langle R \rangle = 1.282 \pm 0.004(\text{stat}) \pm 0.007(\text{syst})$$



CRAFT - GLB Muons 1 Leg Barrel Only\_cmb

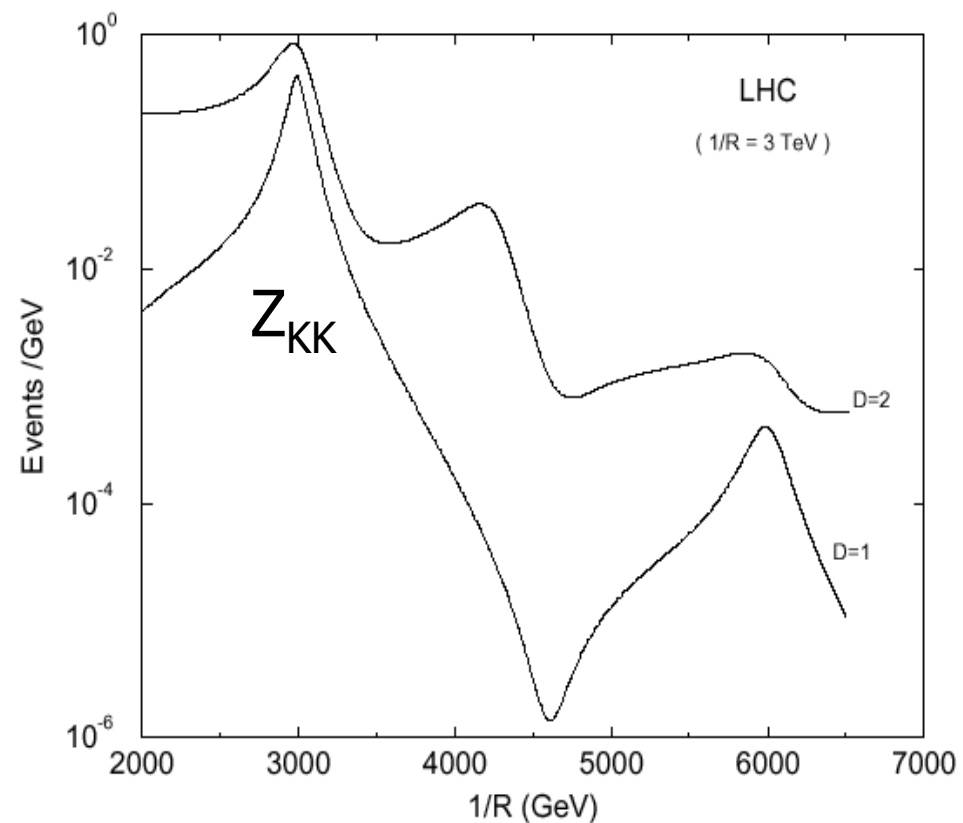




# TeV<sup>-1</sup> Extra Dimensions

- Intermediate-size **extra dimensions** with  $\sim \text{TeV}^{-1}$  radius
- Introduced by **Antoniadis** [PL **B246**, 377 (1990)] in the string theory context
- Used by **Dienes, Dudas, and Gherghetta** [PL **B436**, 55 (1998)] to allow for low-energy unification
  - Expect  $Z_{KK}$ ,  $W_{KK}$ ,  $g_{KK}$  resonances at the LHC energies
  - At lower energies, can study effects of virtual exchange of the Kaluza-Klein modes of vector bosons
- Current indirect constraints come from **precision EW** measurements:
  - $1/r \sim 6 \text{ TeV}$

Antoniadis, Benaklis, and Quiros  
[PL **B460**, 176 (1999)]





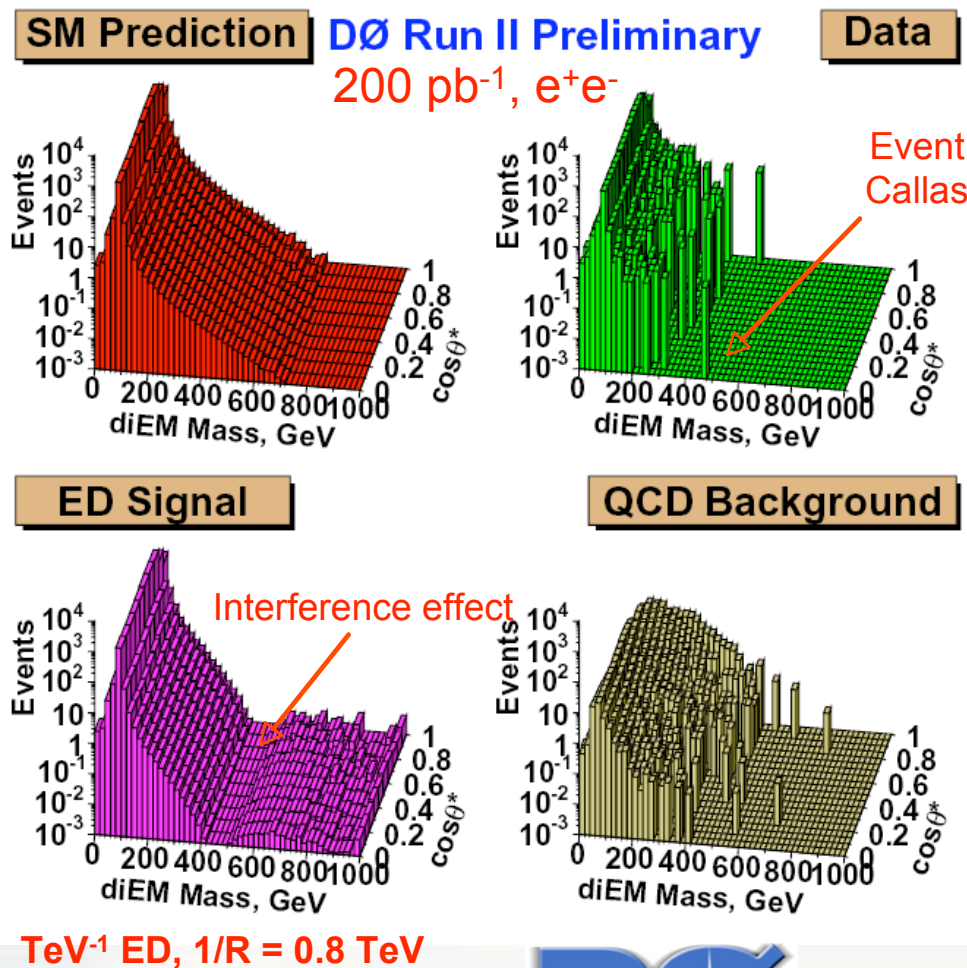
# Current Limits on $\text{TeV}^{-1}$ ED

From Cheung & GL [PRD **65**, 076003 (2002)]

	$\eta \text{ (TeV}^{-2}\text{)}$	$\eta_{95} \text{ (TeV}^{-2}\text{)}$	$M_C^{95} \text{ (TeV)}$
LEP 2:			
hadronic cross section, ang. dist., $R_{b,c}$	$-0.33^{+0.13}_{-0.13}$	0.12	5.3
$\mu, \tau$ cross section & ang. dist.	$0.09^{+0.18}_{-0.18}$	0.42	2.8
$ee$ cross section & ang. dist.	$-0.62^{+0.20}_{-0.20}$	0.16	4.5
LEP combined	$-0.28^{+0.092}_{-0.092}$	0.076	6.6
HERA:			
NC	$-2.74^{+1.49}_{-1.51}$	1.59	1.4
CC	$-0.057^{+1.28}_{-1.31}$	2.45	1.2
HERA combined	$-1.23^{+0.98}_{-0.99}$	1.25	1.6
TEVATRON:			
Drell-yan	$-0.87^{+1.12}_{-1.03}$	1.96	1.3
Tevatron dijet	$0.46^{+0.37}_{-0.58}$	1.0	1.8
Tevatron top production	$-0.53^{+0.51}_{-0.49}$	9.2	0.60
Tevatron combined	$-0.38^{+0.52}_{-0.48}$	0.65	2.3
All combined	$-0.29^{+0.090}_{-0.090}$	0.071	6.8



# First Dedicated Search for $\text{TeV}^{-1}$ ED

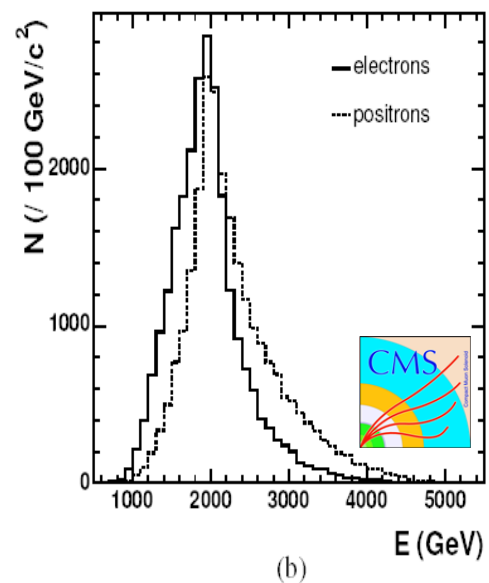
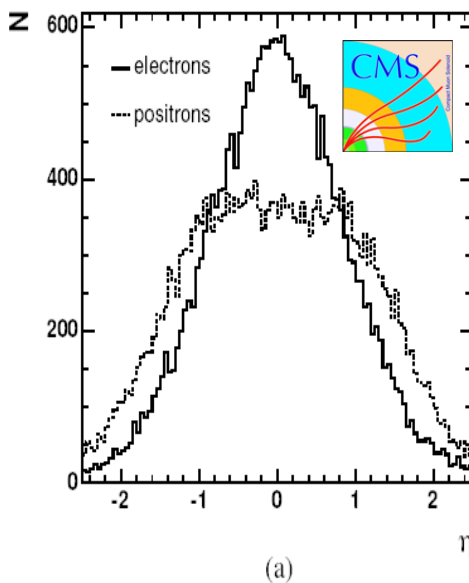
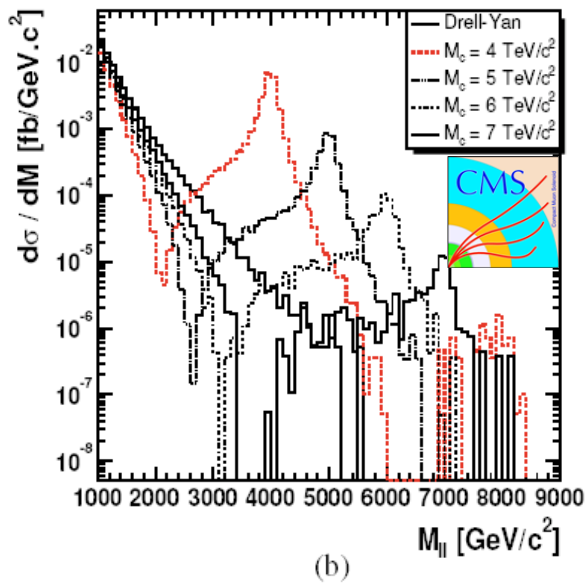
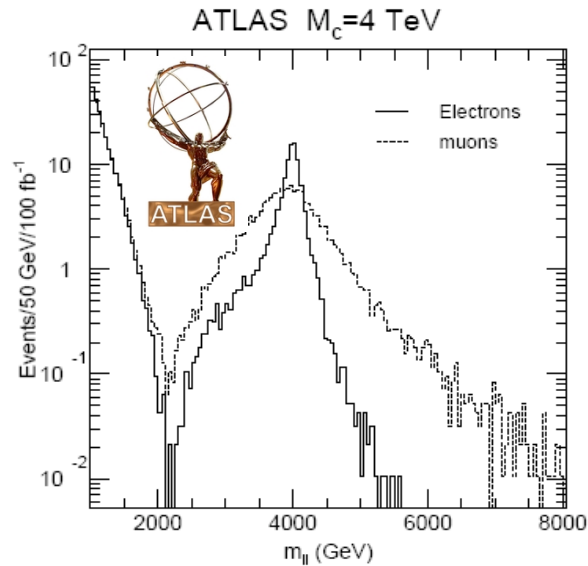
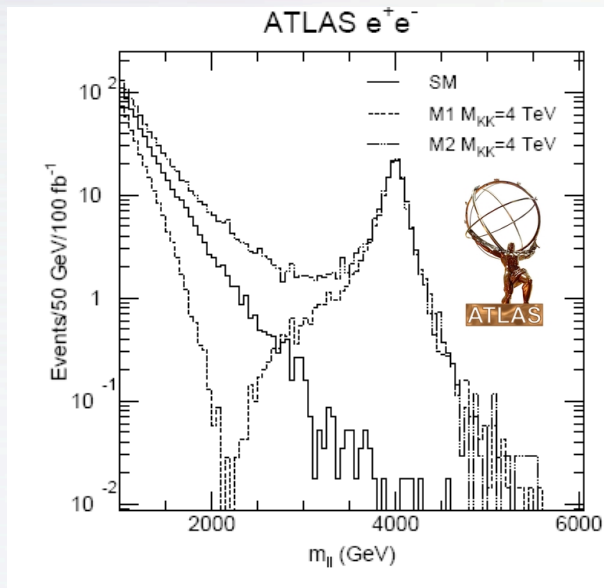


- While the Tevatron sensitivity is inferior to indirect limits, it explores the effects of virtual KK modes at higher energies, i.e. complementary to those in the EW data
- DØ has performed the first dedicated search of this kind in the dielectron channel based on 200 pb<sup>-1</sup> of Run II data ( $Z_{KK}, \gamma_{KK} \rightarrow e^+e^-$ )
- The 2D-technique similar to the search for ADD effects in the virtual exchange yields the best sensitivity Cheung, GL[PRD **65**, 076003 (2002)]
- Data agree with the SM predictions, which resulted in the following limit:
  - $1/r > 1.12 \text{ TeV @ 95\% CL}$
  - $r < 1.75 \times 10^{-19} \text{ m}$
- From dijets (700 pb<sup>-1</sup>) the limit is:
  - $1/r > 1.4 \text{ TeV @ 95\% CL}$





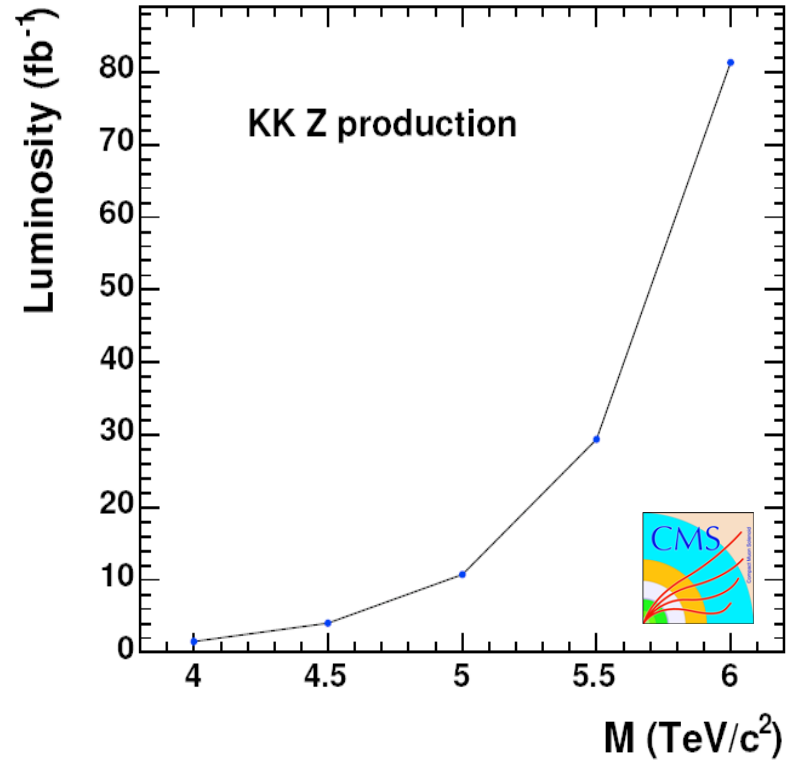
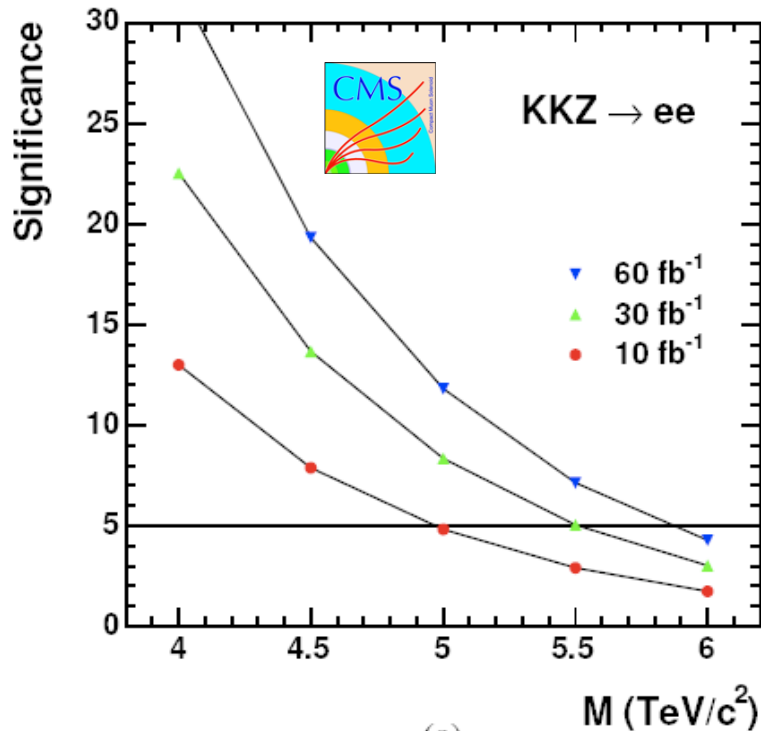
# LHC: KK Excitations of the Z Boson





# KK Resonance Reach at the LHC

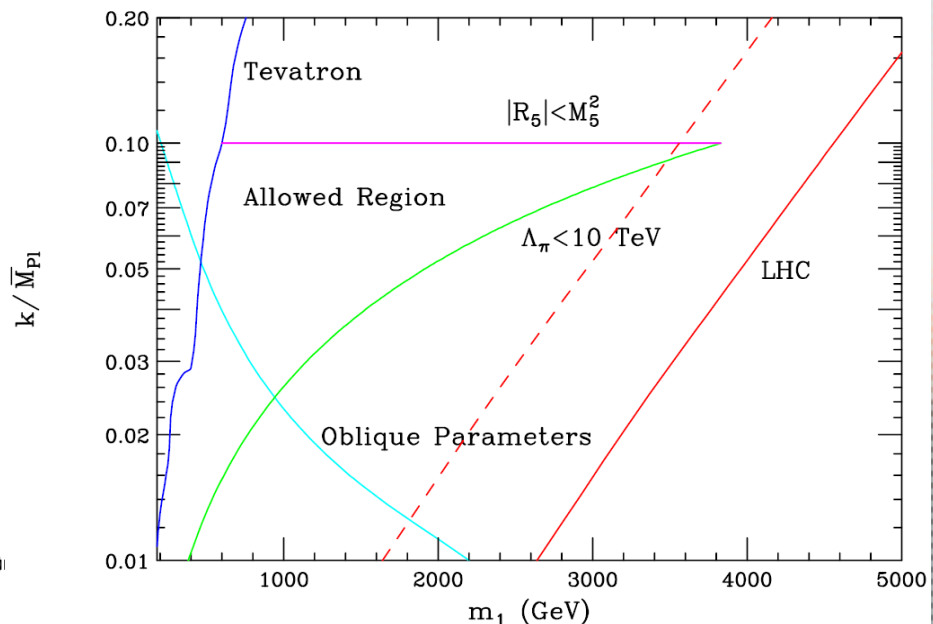
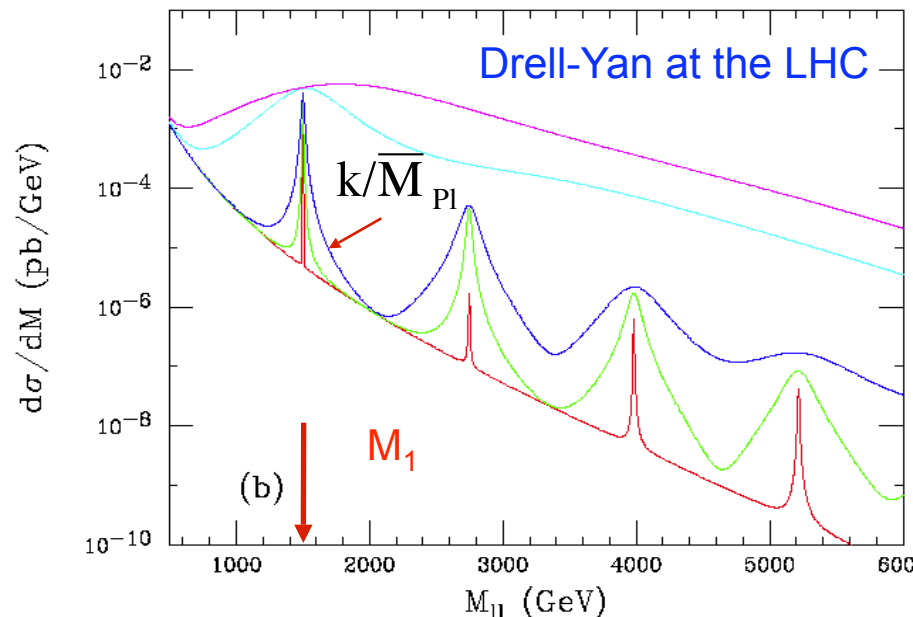
- Dramatic reach even with  $\sim 1 \text{ fb}^{-1}$





# Randall-Sundrum Model Observables

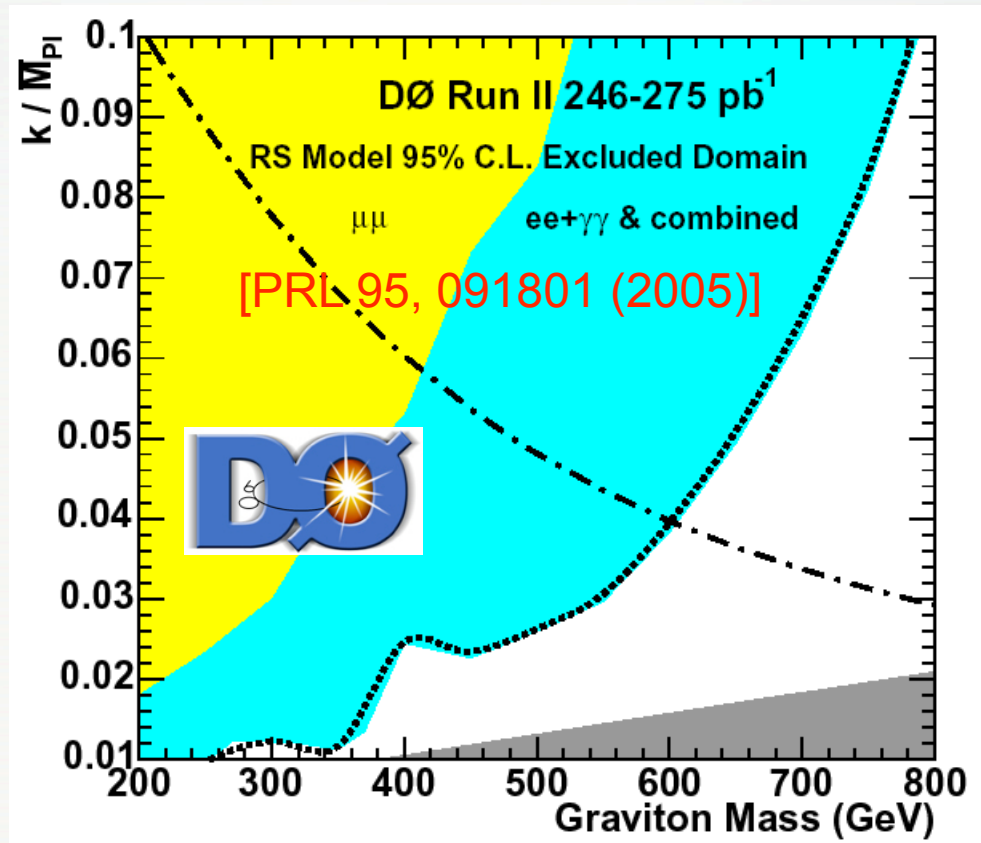
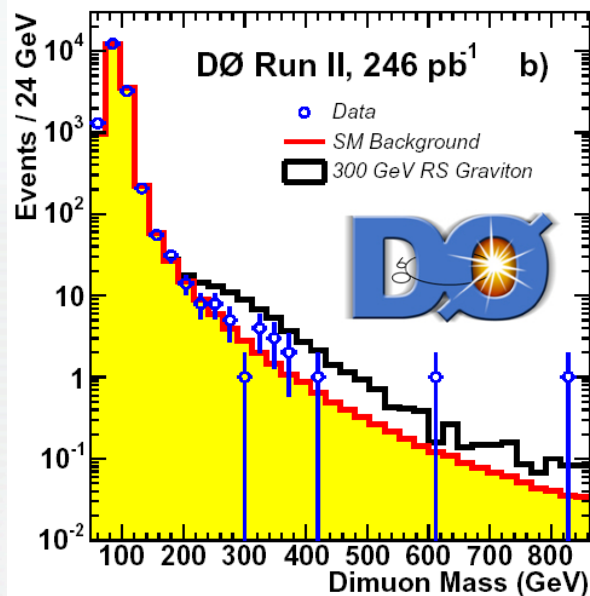
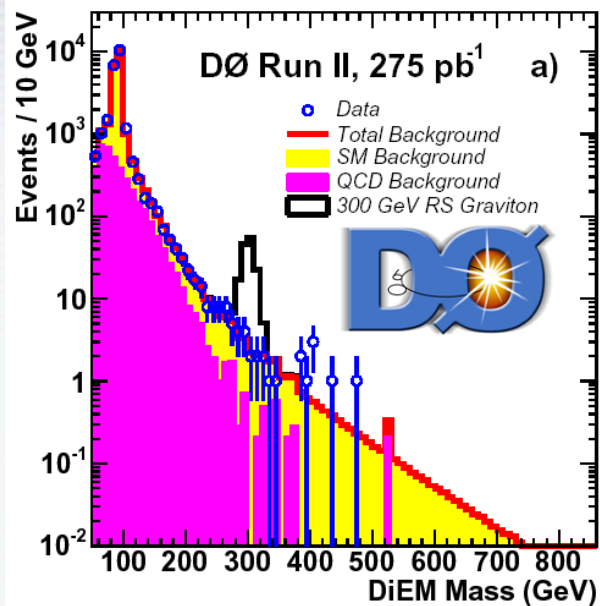
- Need only **two parameters** to define the model:  **$k$**  and  **$r$**
- **Equivalent set of parameters:**
  - The mass of the first KK mode,  $M_1$
  - Dimensionless coupling  $k/\overline{M}_{\text{Pl}}$ , which determines the graviton width
- To avoid fine-tuning and non-perturbative regime, **coupling can't be too large or too small**
- $0.01 \leq k/\overline{M}_{\text{Pl}} \leq 0.10$  is the expected range
- Gravitons are narrow
- Similar observables for  $Z_{\text{KK}}/g_{\text{KK}}$  in  $\text{TeV}^{-1}$  models



Davoudiasl, Hewett, Rizzo [PRD **63**, 075004 (2001)]



# First Search for RS Gravitons



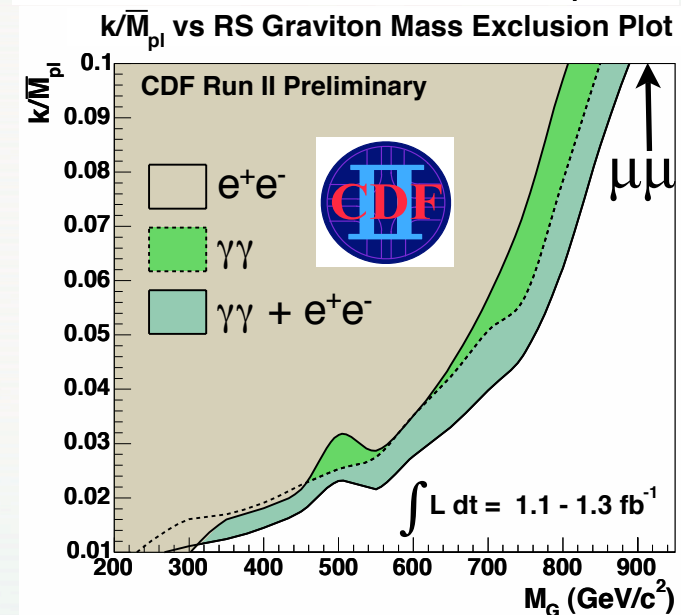
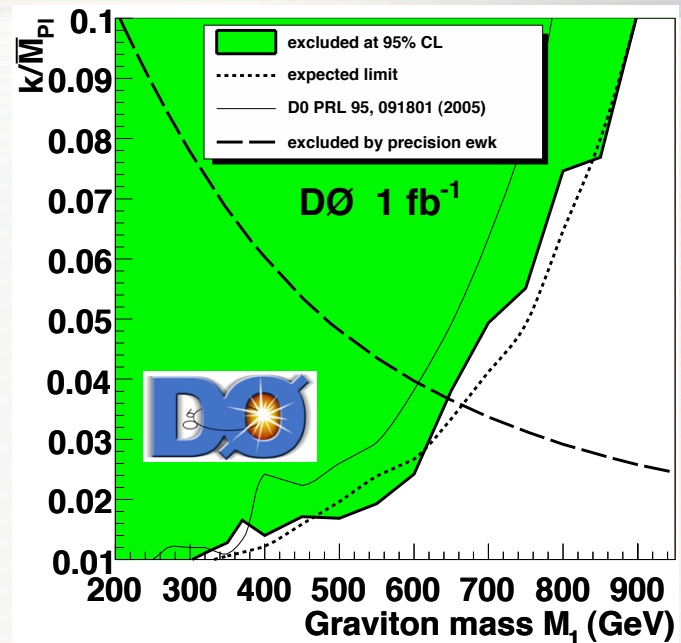
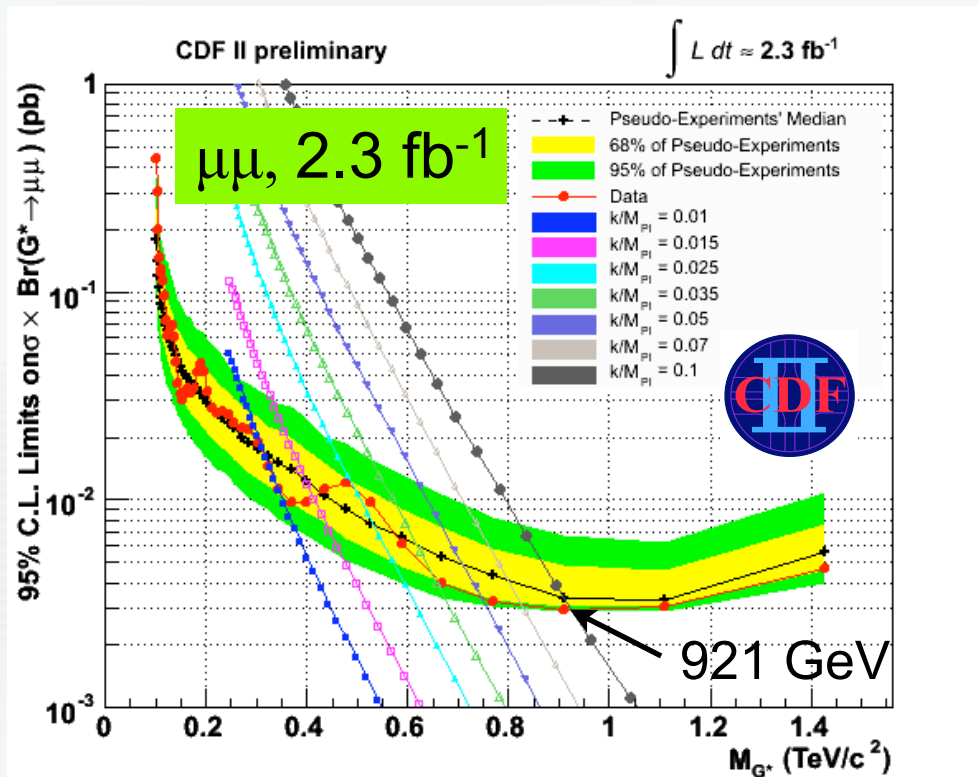
Assume a fixed K-factor of 1.3 for the signal





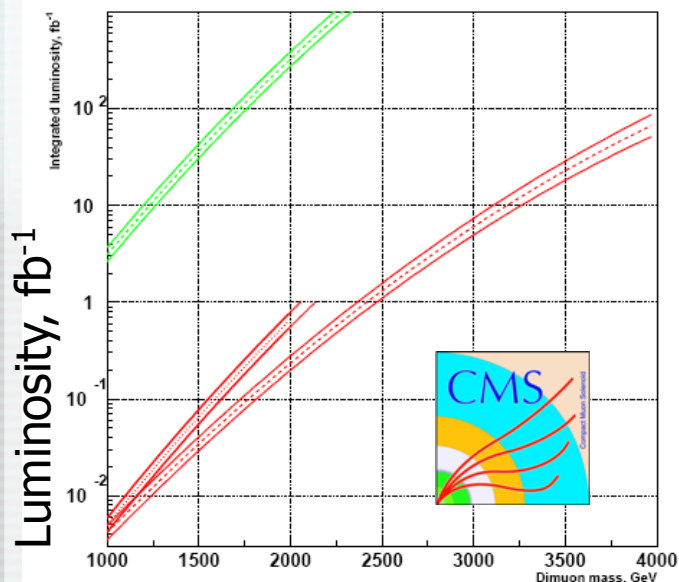
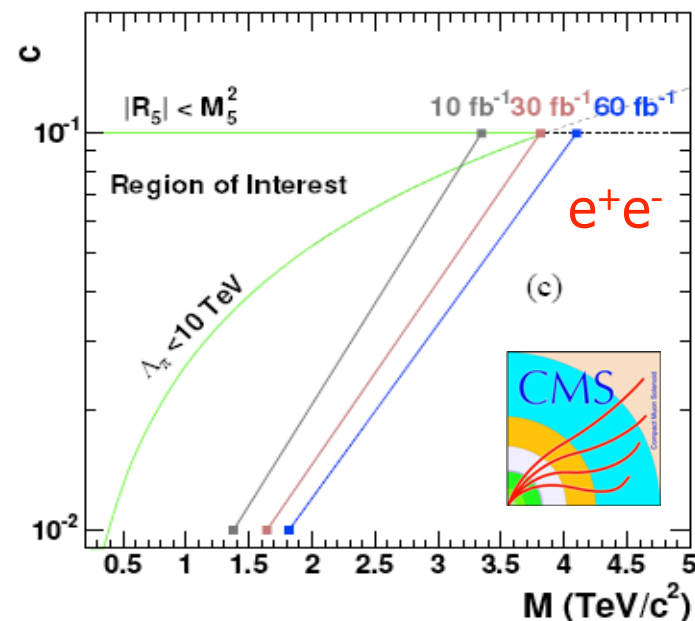
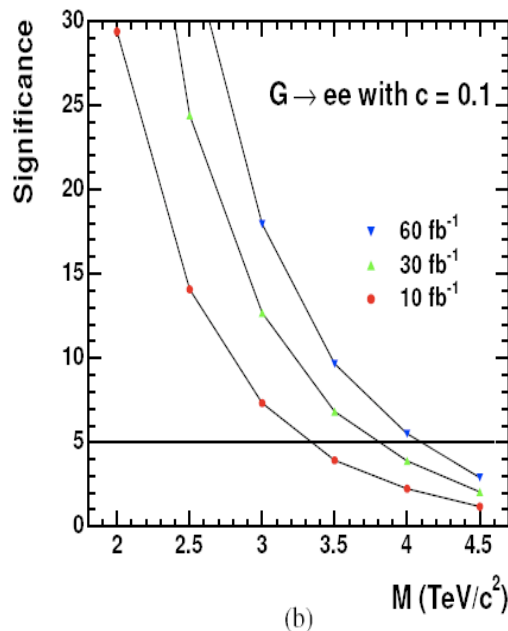
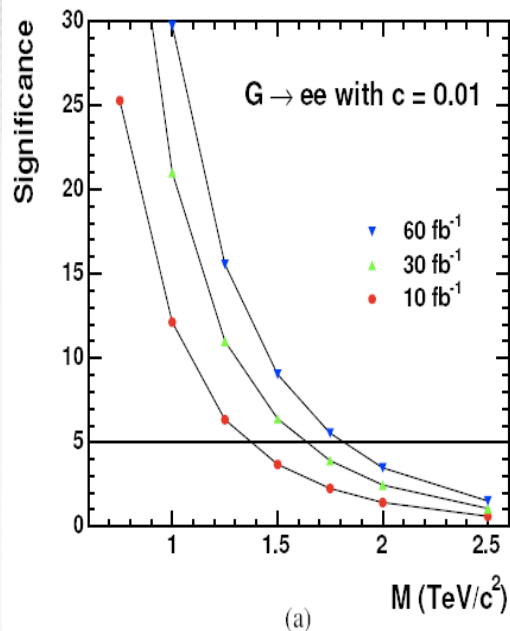
# Most Recent Limits

- Latest limits are 10% higher than the original ones despite 4x statistics
  - Tevatron sensitivity has really maxed out - need higher energies!

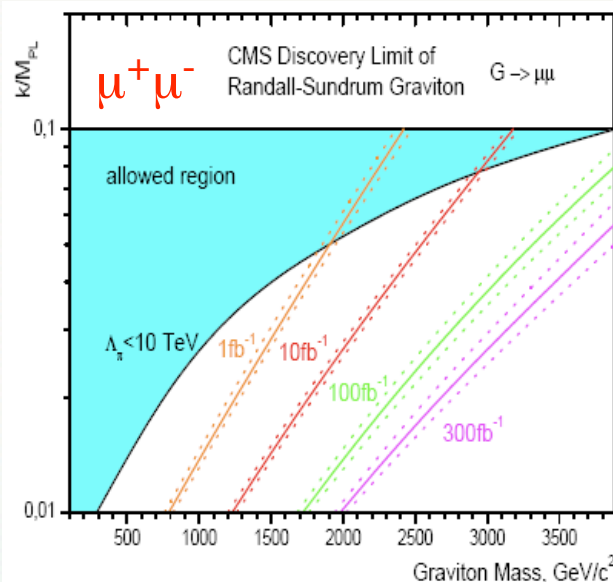




# LHC: Randall-Sundrum Graviton Reach



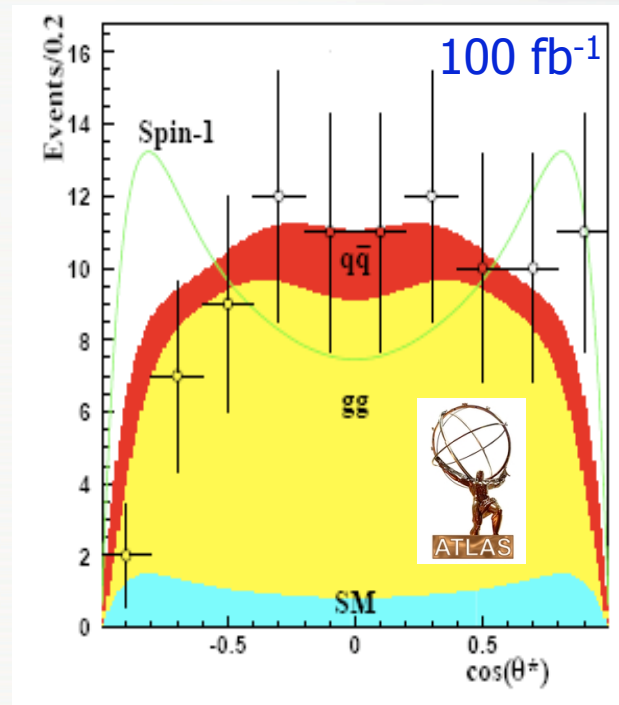
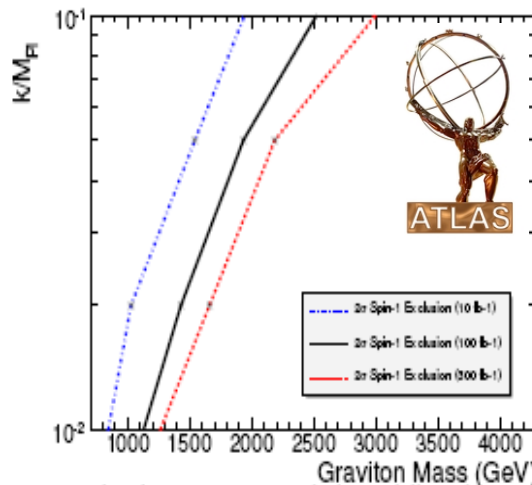
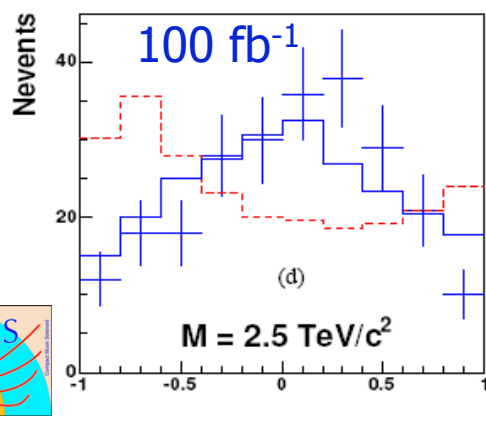
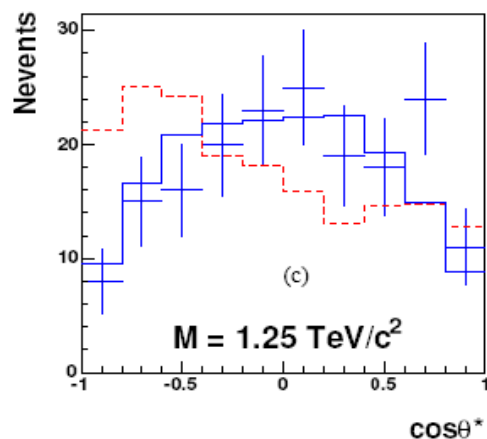
Coupling constant $c$	Estimator	1 $\text{fb}^{-1}$	10 $\text{fb}^{-1}$
0.01	$S_{cP}$	0.75	1.20
	$S_{cL}$	0.77	1.21
	$S_L$	0.78	1.23
0.02	$S_{cP}$	1.21	1.72
	$S_{cL}$	1.22	1.72
	$S_L$	1.22	1.74
0.05	$S_{cP}$	1.83	2.48
	$S_{cL}$	1.85	2.49
	$S_L$	1.85	2.51
0.1	$S_{cP}$	2.34	3.11
	$S_{cL}$	2.36	3.13
	$S_L$	2.36	3.16





# LHC: Graviton Spin?

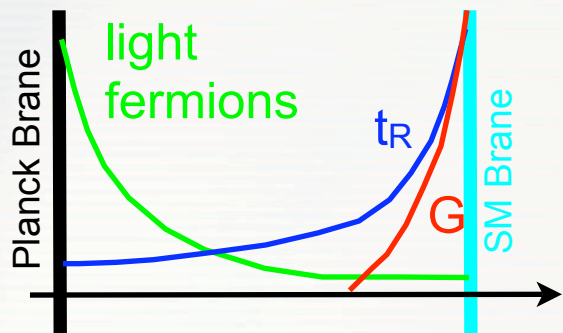
- Not in the early running!
  - “One event – discovery; two events – cross section measurement; three events – angular distributions”



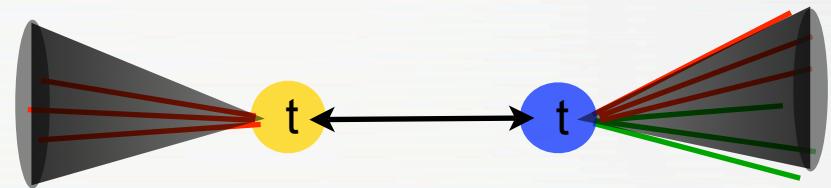


# But: Life May be More Complicated!

- **Simple RS model** has many potential **problems**: FCNC, CP-violation
  - Those can be solved by putting fermions in the bulk
- **Top quark is localized near the SM brane**; light fermions are near the Planck brane
- **Graviton mainly couples to the top quark**, and thus the dominant decay mode is a pair of top quarks



- For graviton masses  $\sim 2\text{--}3\text{ TeV}$ , **top quarks emerge highly boosted**, which makes it challenging to reconstruct them



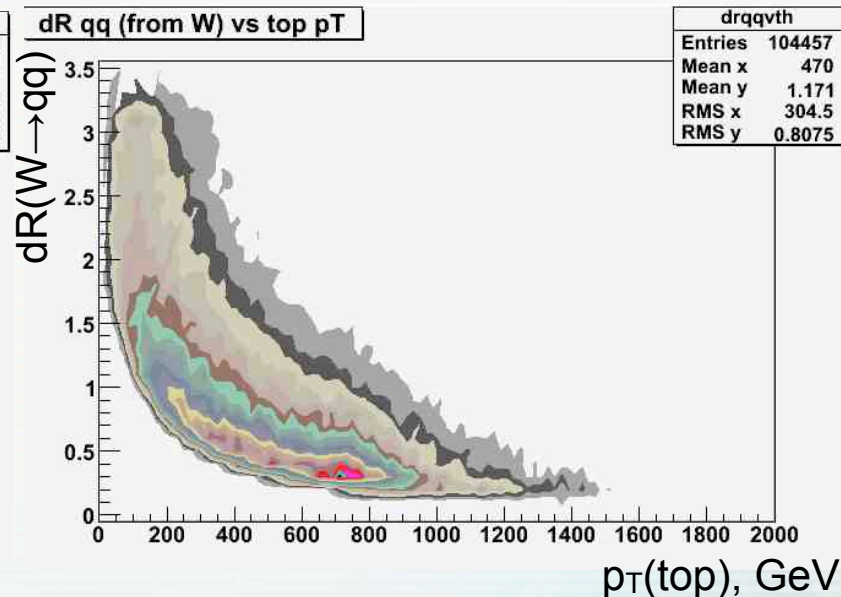
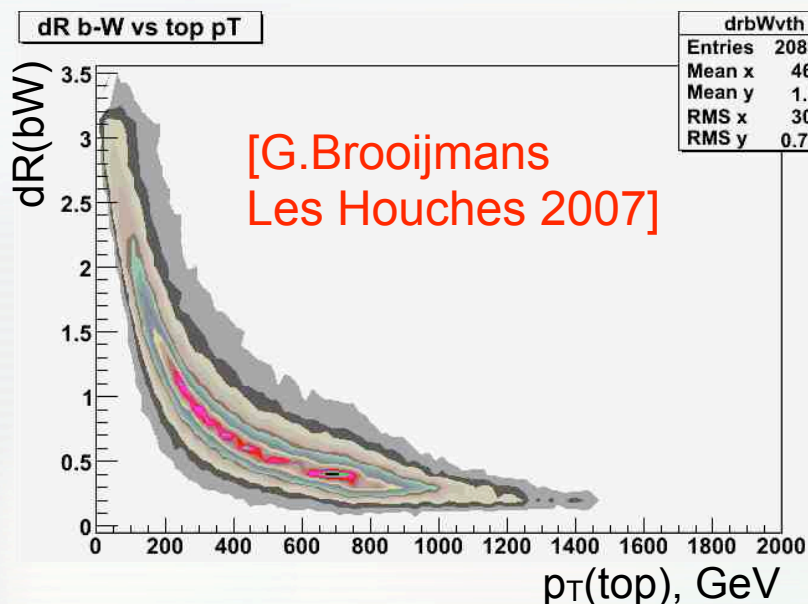
- Several challenges:
  - for 3-jet top decays jets are often merged in a single “fat” jet
  - b-tagging efficiency drops dramatically, as the opening angle between the tracks becomes small.





# Possible Remedies

- Several have been suggested:
  - Use of top-jet mass: poorly defined; depends on the jet algorithm and pile-up/underlying event; depends on the calorimeter granularity
  - 3D b-tagging
  - $k_T$  algorithm with small  $D \sim 0.1$ -0.2 within the “fat” jet
- Requires serious experimental studies with realistic detector effects
  - dedicated groups in both ATLAS and CMS are exploring this topic





# More Challenges: Universal ED

- The most “democratic” ED model: *all* the SM fields are free to propagate in extra dimension(s) with the size  $r = 1/M_c \sim 1 \text{ TeV}^{-1}$  Appelquist, Cheng, Dobrescu [PRD **64**, 035002 (2001)]
  - Instead of chiral doublets and singlets, model contains vector-like quarks and leptons
  - Gravitational force is not included in this model
- The number of universal extra dimensions is not fixed:
  - it’s feasible that there is just one (MUED)
  - the case of two extra dimensions is theoretically attractive, as it breaks down to the chiral Standard Model and has additional nice features, such as guaranteed proton stability, etc.
- Every particle acquires KK modes with the masses  $M_n^2 = M_0^2 + M_c^2$ ,  $n = 0, 1, 2, \dots$
- Kaluza-Klein number ( $n$ ) is conserved at tree level, i.e.  $n_1 \pm n_2 \pm n_3 \pm \dots = 0$ ; consequently, the lightest KK mode could be stable (and is an excellent dark matter candidate Cheng, Feng, Matchev [PRL **89**, 211301 (2002)])
- Hence, first level KK-excitations are produced in pairs, similar to SUSY particles
- Consequently, current limits (dominated by precision electroweak measurements, particularly T-parameter) are sufficiently low ( $M_c \sim 300 \text{ GeV}$  for one ED and of the same order, albeit more model-dependent for  $>1$  ED)
- CDF unpublished Run I analysis based on  $90 \text{ pb}^{-1}$  set  $M_c > 280 \text{ GeV}$

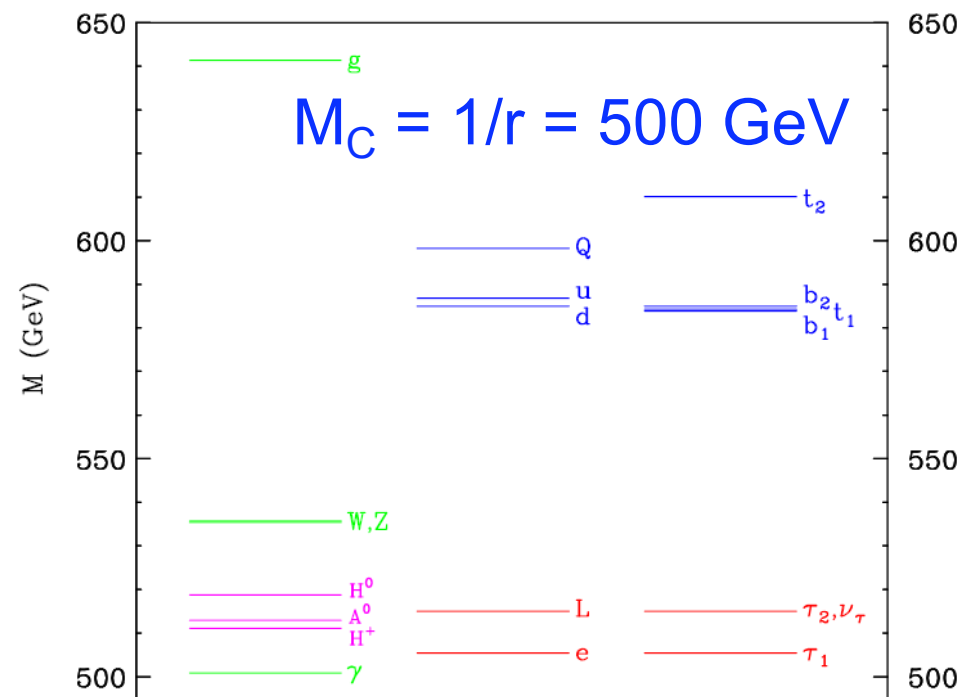


# UED Phenomenology

- Naively, one would expect large clusters of **nearly degenerate states** with masses around  $1/r, 2/r, \dots$
- Cheng, Feng, Matchev, Schmaltz: **not true, as radiative corrections tend to be large** (up to 30%); thus the KK excitation mass spectrum resembles that of SUSY!
- **Minimal UED model with a single extra dimension**, compactified on an  $S_1/Z_2$  orbifold
  - **Odd fields** do not have 0 modes, so we identify them w/ “**wrong**” chiralities, so that they **vanish** in the SM

- $Q, L (q, l)$  are  $SU(2)$  doublets (**singlets**) and contain both chiralities

Cheng, Matchev, Schmaltz  
[PRD **66**, 056006 (2002)]



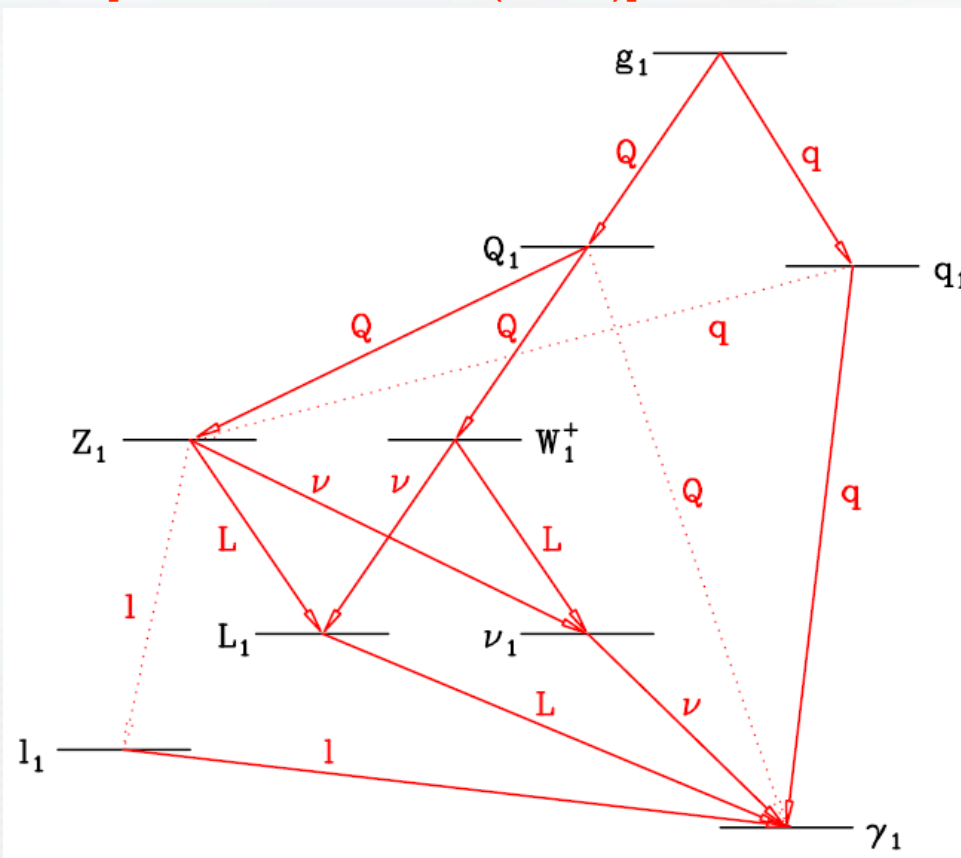


# Mass Spectrum and Decays

- First level KK-states spectroscopy

Cheng, Matchev, Schmaltz

[PRD **66**, 056006 (2002)]



Decay:

$$B(g_1 \rightarrow Q_1 Q) \sim 50\%$$

$$B(g_1 \rightarrow q_1 q) \sim 50\%$$

$$B(q_1 \rightarrow q \gamma_1) \sim 100\%$$

$$B(t_1 \rightarrow W_1 b, H_1^+ b) \sim 100\%$$

$$B(Q_1 \rightarrow Q Z_1 : W_1 : \gamma_1) \sim 33\% : 65\% : 2\%$$

$$B(W_1 \rightarrow \nu L_1 : \nu_1 L) = 1/6 : 1/6 \text{ (per flavor)}$$

$$B(Z_1 \rightarrow \nu \nu_1 : L L_1) \sim 1/6 : 1/6 \text{ (per flavor)}$$

$$B(L_1 \rightarrow \gamma_1 L) \sim 100\%$$

$$B(\nu_1 \rightarrow \gamma_1 \nu) \sim 100\%$$

$$B(H_1^\pm \rightarrow \gamma \gamma_1, H^\pm \gamma_1) \sim 100\%$$

Production:

$$q_1 q_1 + X \rightarrow ME_T + \text{jets} (\sim \sigma_{\text{had}}/4); \text{ but:}$$

low  $ME_T$

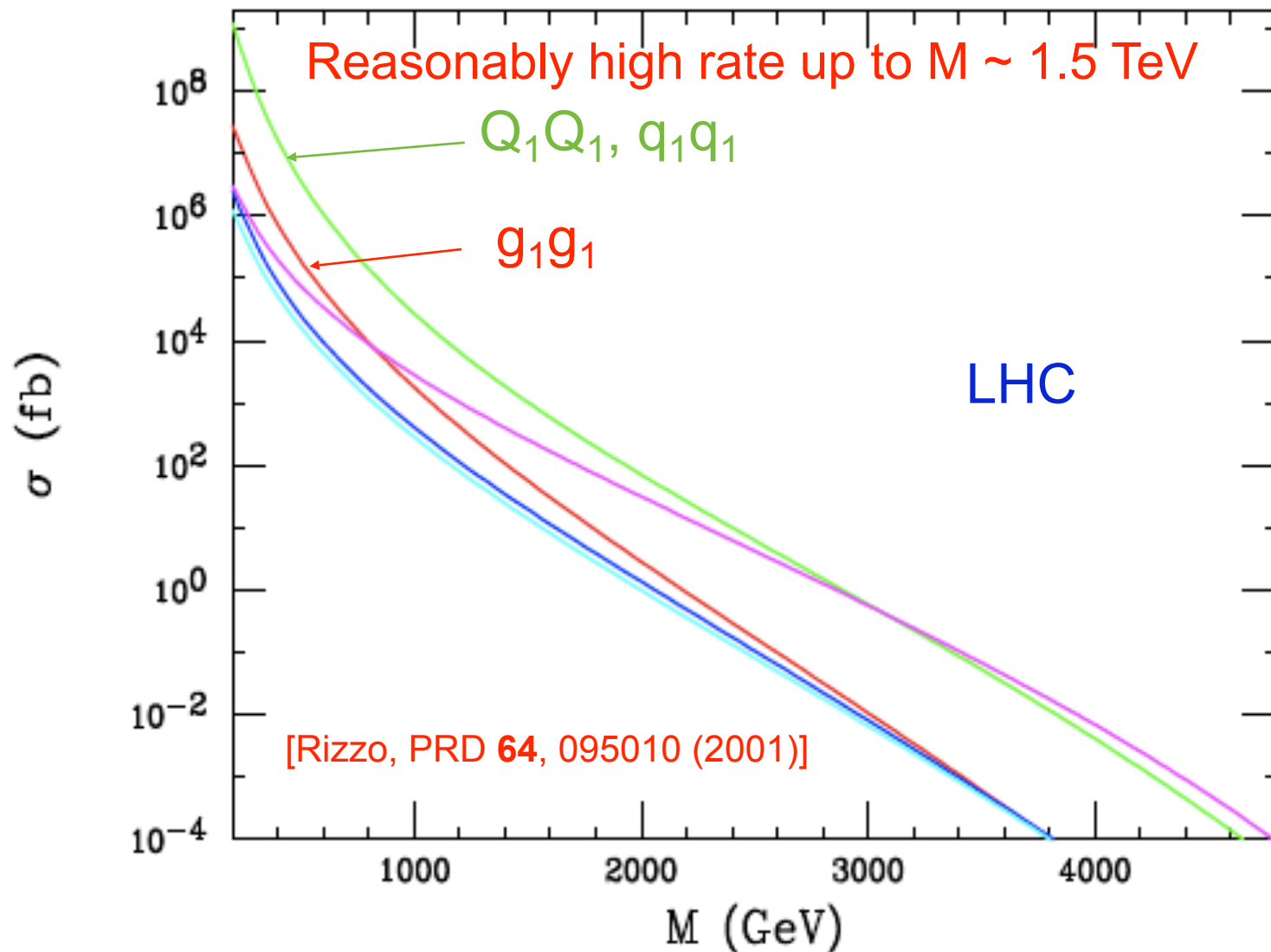
$$Q_1 Q_1 + X \rightarrow V_1 V'_1 + \text{jets} \rightarrow 2\text{-}4 \ell + ME_T$$

$$(\sim \sigma_{\text{had}}/4)$$





# Production Cross Section

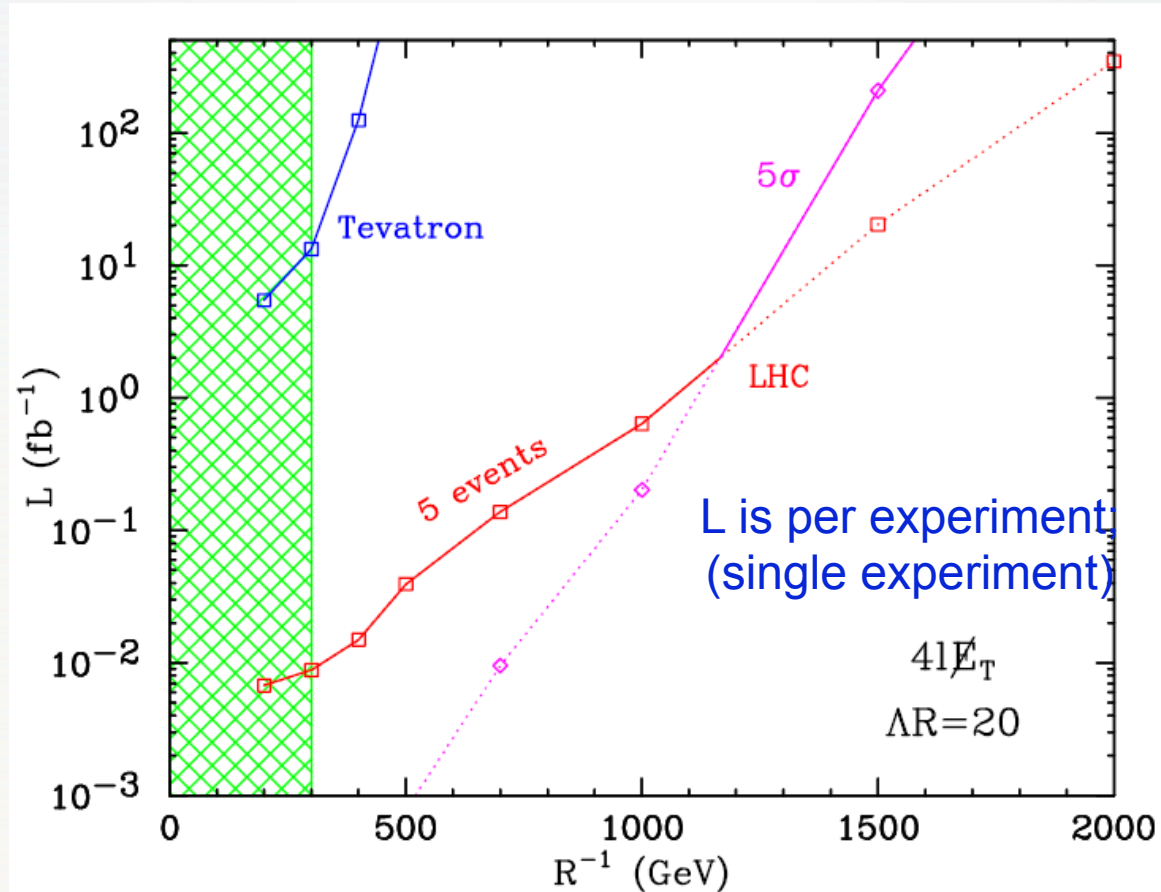




# Sensitivity in the Four-Lepton Mode

- Only the gold-plated 4-leptons +  $ME_T$  mode has been considered in the original paper and the subsequent studies
- Other promising channels:
  - dileptons + jets +  $ME_T$  + X (x9 cross section)
  - trileptons + jets +  $ME_T$  + X (x5 cross section)
  - Single production of the second KK excitation (via one loop)
- Detailed simulations are required: CompHEP and PYTHIA implementations now exist

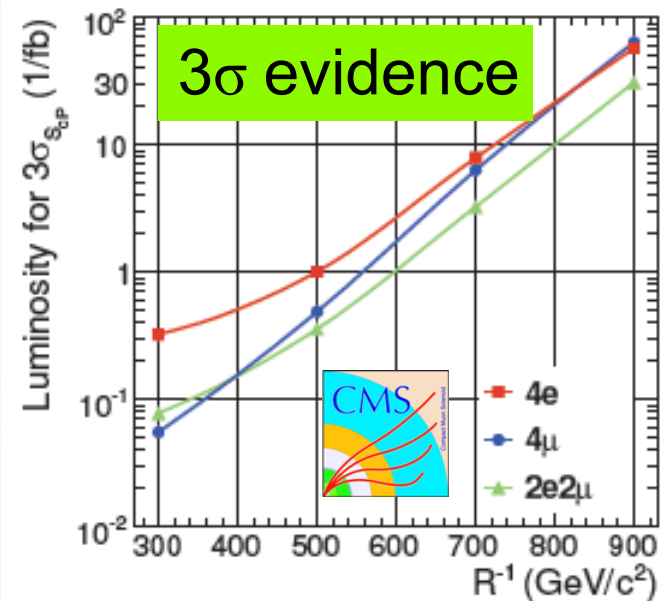
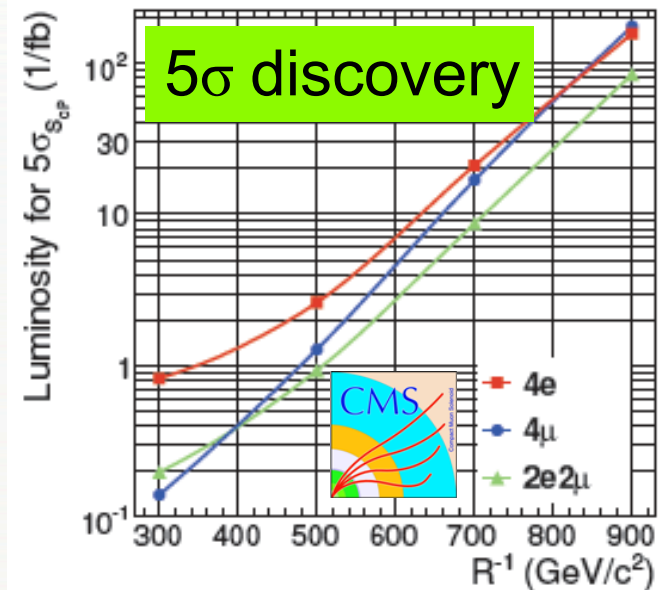
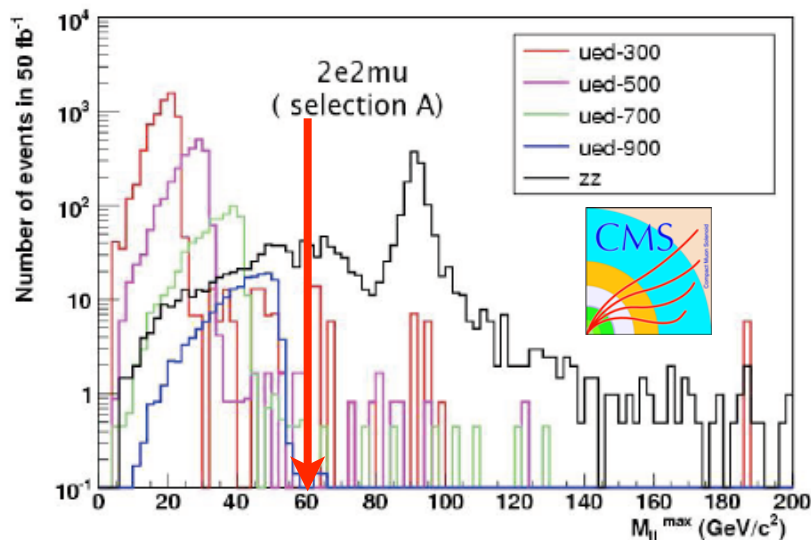
Cheng, Matchev, Schmaltz [PRD **66**, 056006 (2002)]





# Early UED Searches in CMS

- Very recently approved analysis by the LIP group
- Consider  $4e$ ,  $4\mu$ ,  $2e2\mu$  channels
- Tight selection for low  $1/R$  and looser selection for high  $1/R$
- Signal is found at low dilepton invariant mass
- Background is dominated by the physics  $(Z/\gamma^*)(Z/\gamma^*)$  background







# Black Holes at the LHC?







# Black Holes on Demand

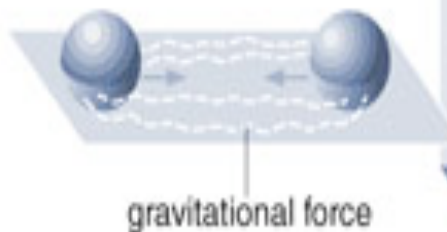
## Black Holes on Demand

NYT, 9/11/01

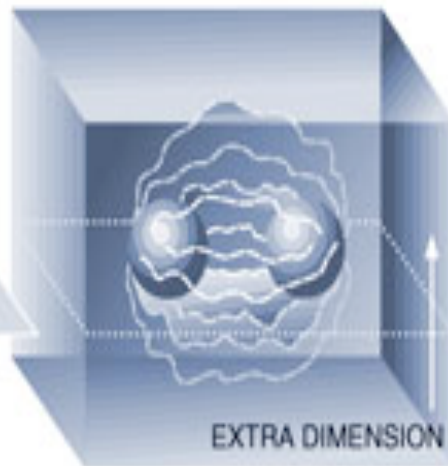
The New York Times  
ON THE WEB

Scientists are exploring the possibility of producing miniature black holes on demand by smashing particles together. Their plans hinge on the theory that the universe contains more than the three dimensions of everyday life. Here's the idea:

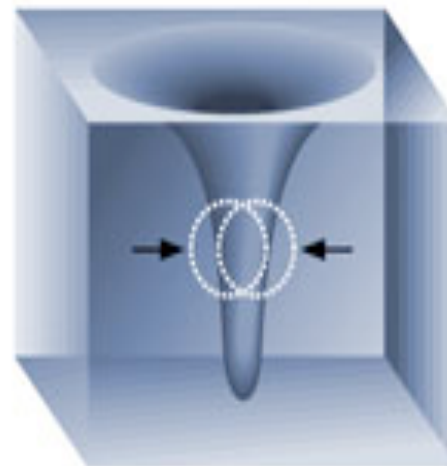
*Particles collide in three dimensional space, shown below as a flat plane.*



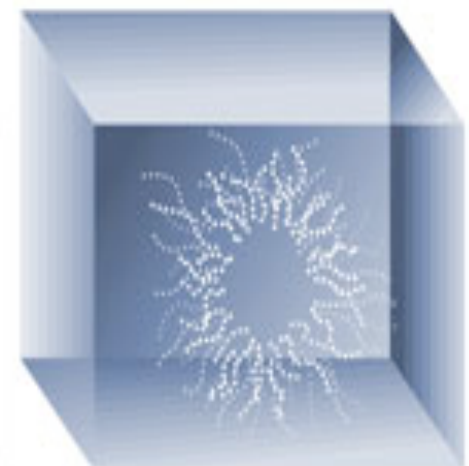
As the particles approach in a particle accelerator, their gravitational attraction increases steadily.



When the particles are extremely close, they may enter space with more dimensions, shown above as a cube.



The extra dimensions would allow gravity to increase more rapidly so a black hole can form.

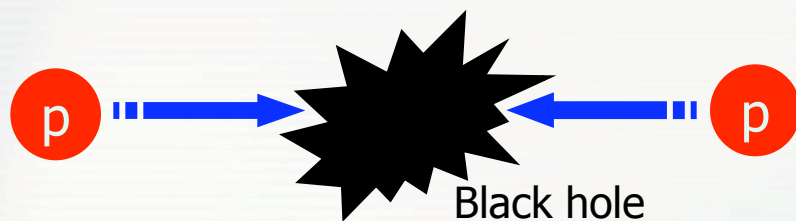


Such a black hole would immediately evaporate, sending out a unique pattern of radiation.

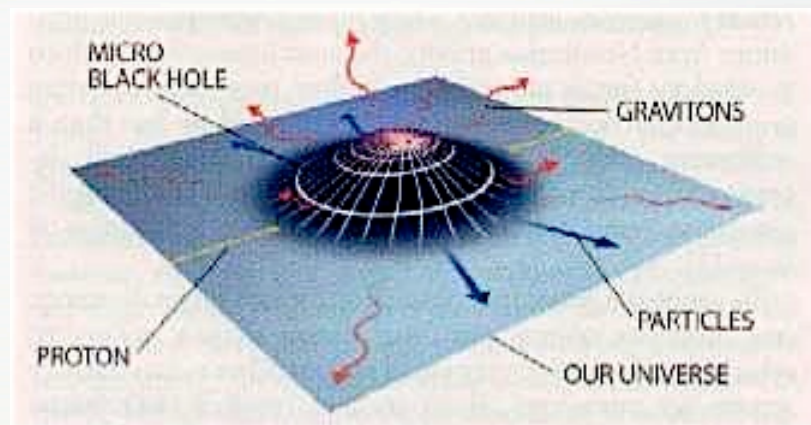


# BH at LHC: Theoretical Framework

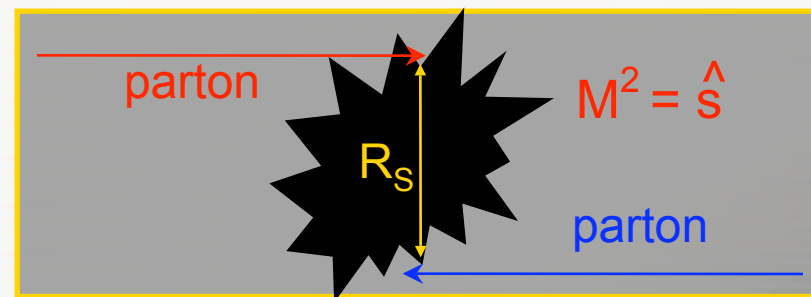
- Based on the work done with Dimopoulos a few years ago [PRL **87**, 161602 (2001)] and a related study by Giddings/Thomas [PRD **65**, 056010 (2002)]
- Extends previous, more theoretical studies by Argyres/Dimopoulos/March-Russell [PL **B441**, 96 (1998)], Banks/Fischler [JHEP, **9906**, 014 (1999)], Emparan/Horowitz/Myers [PRL **85**, 499 (2000)] to collider phenomenology
- Big surprise: BH production is not an exotic remote possibility, but the dominant effect!
- Main idea: when the c.o.m. energy reaches the fundamental Planck scale, a BH is formed!
- Also true in the RS models where  $\Lambda_\pi$  is the characteristic scale



Artist's view:



Cross section is given by a black disk approximation:



$\sigma \sim \pi R_s^2 \sim 1 \text{ TeV}^{-2} \sim 10^{-38} \text{ m}^2 \sim 100 \text{ pb}$   
 Comparable with that of the top-quark pair production!



# Assumptions and Approximations

- Fundamental limitation: our **lack of knowledge of quantum gravity effects** close to the Planck scale
- Consequently, **no attempts for partial improvement** of the results, e.g.:
  - Grey body factors
  - BH spin, charge, color hair
  - Relativistic effects and time-dependence
- The underlying assumptions rely on two simple qualitative properties:
  - The absence of small couplings;
  - The “democratic” nature of BH decays
- We **expect these features to survive for light BH**
- Use **semi-classical approach** strictly valid only for  $M_{\text{BH}} \gg M_{\text{P}}$ ; only consider  $M_{\text{BH}} > M_{\text{P}}$
- Clearly, these are **important limitations**, but there is **no way around them without the knowledge of QG**



# Black Hole Production

- Schwarzschild radius is given by Argyres et al. [hep-th/9808138], after Myers/Perry [Ann. Phys. **172**, 304(1986)]; it leads to:

$$\sigma(\hat{s} = M_{\text{BH}}^2) = \pi R_S^2 = \frac{1}{M_{\text{Pl}}^2} \left[ \frac{M_{\text{BH}}}{M_{\text{Pl}}} \frac{8\Gamma\left(\frac{n+3}{2}\right)}{n+2} \right]^{\frac{2}{n+1}}$$

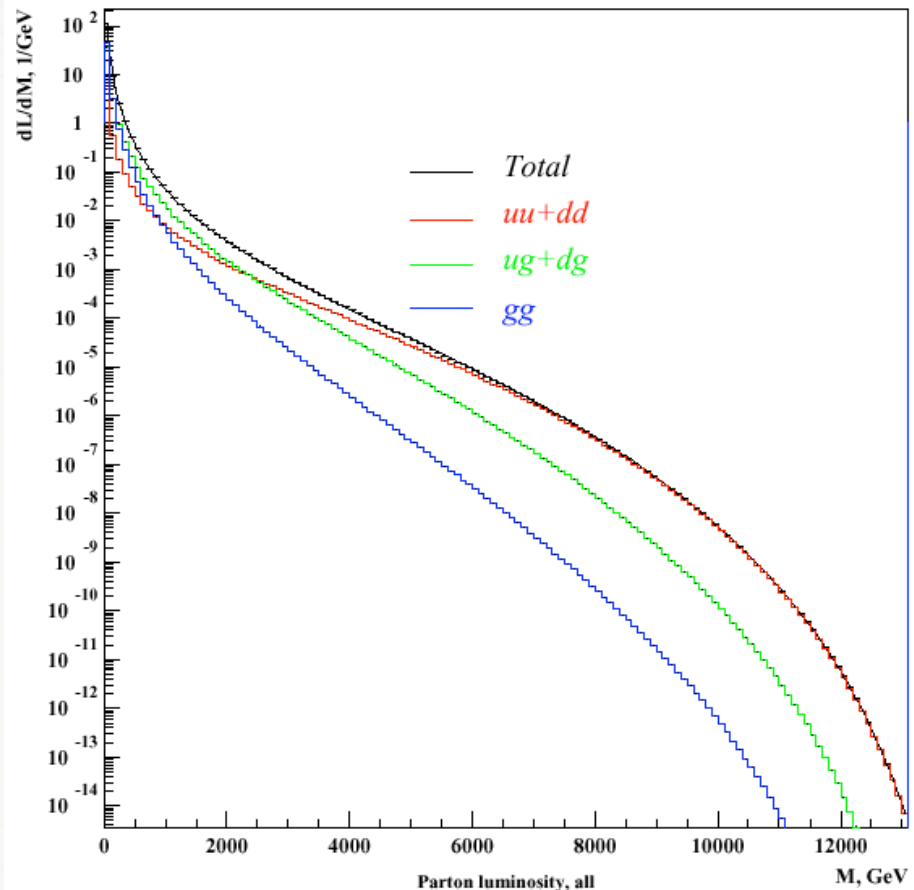
- Use parton luminosity approach with quark momentum distribution given by parton distribution functions

$$\frac{d\sigma(pp \rightarrow \text{BH} + X)}{dM_{\text{BH}}} = \frac{dL}{dM_{\text{BH}}} \hat{\sigma}(ab \rightarrow \text{BH})|_{\hat{s}=M_{\text{BH}}^2}$$

$$\frac{dL}{dM_{\text{BH}}} = \frac{2M_{\text{BH}}}{s} \sum_{a,b} \int_{M_{\text{BH}}^2/s}^1 \frac{dx_a}{x_a} f_a(x_a) f_b\left(\frac{M_{\text{BH}}^2}{sx_a}\right)$$

- Note: at c.o.m. energies  $\sim 1$  TeV the dominant contribution is from quark-quark interactions (BH w/ color,  $B \neq 0$ )

Dimopoulos, GL [PRL **87**, 161602 (2001)]







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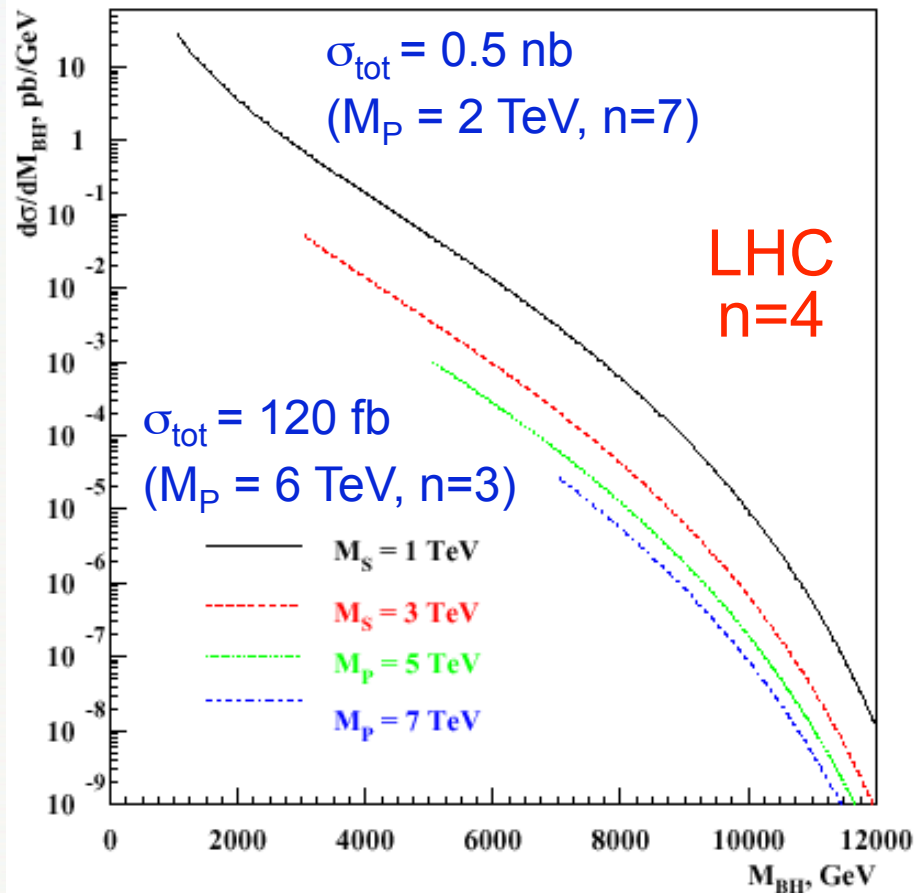
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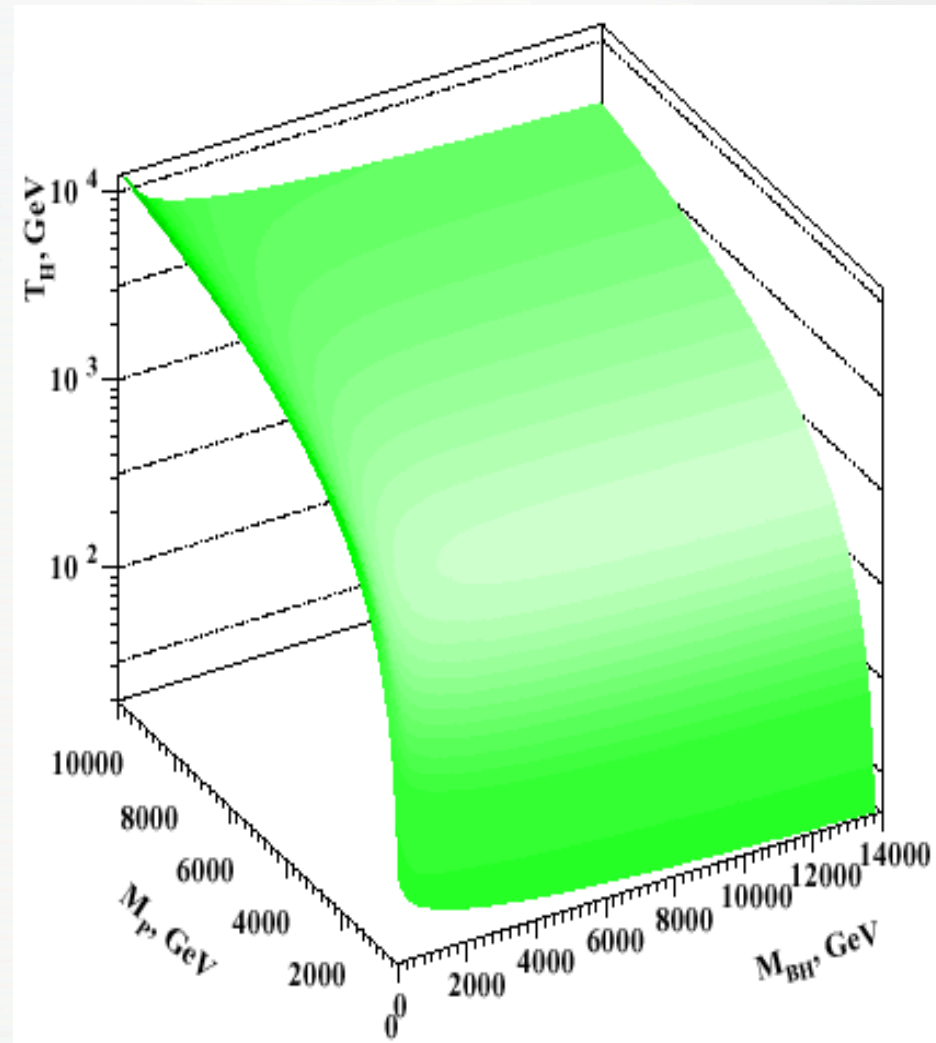




# Black Hole Decay

- **Hawking temperature:**  $R_S T_H = (n+1)/4\pi$   
(in natural units  $\hbar = c = k = 1$ )
- **BH radiates mainly in our 3D world:**  
**Emparan/Horowitz/Myers**  
**[PRL 85, 499 (2000)]**
  - $\lambda \sim 2\pi/T_H > R_S$ ; hence, the **BH is** a point radiator, **producing s-waves**, which depends only on the radial component
  - The **decay into a particle on the brane and in the bulk is thus the same**
  - Since there are **much more particles on the brane, than in the bulk**, decay into gravitons is largely suppressed
- **Democratic couplings to  $\sim 120$  SM d.o.f.** yield probability of Hawking evaporation into  $\gamma$ ,  $\ell^\pm$ , and  $\nu$   **$\sim 2\%$ ,  $10\%$ , and  $5\%$  respectively**
- Averaging over the BB spectrum gives **average multiplicity of decay products:**

$$\langle N \rangle \approx \frac{M_{\text{BH}}}{2T_H}$$



**Stefan's law:**  $\tau \sim 10^{-26} \text{ s}$

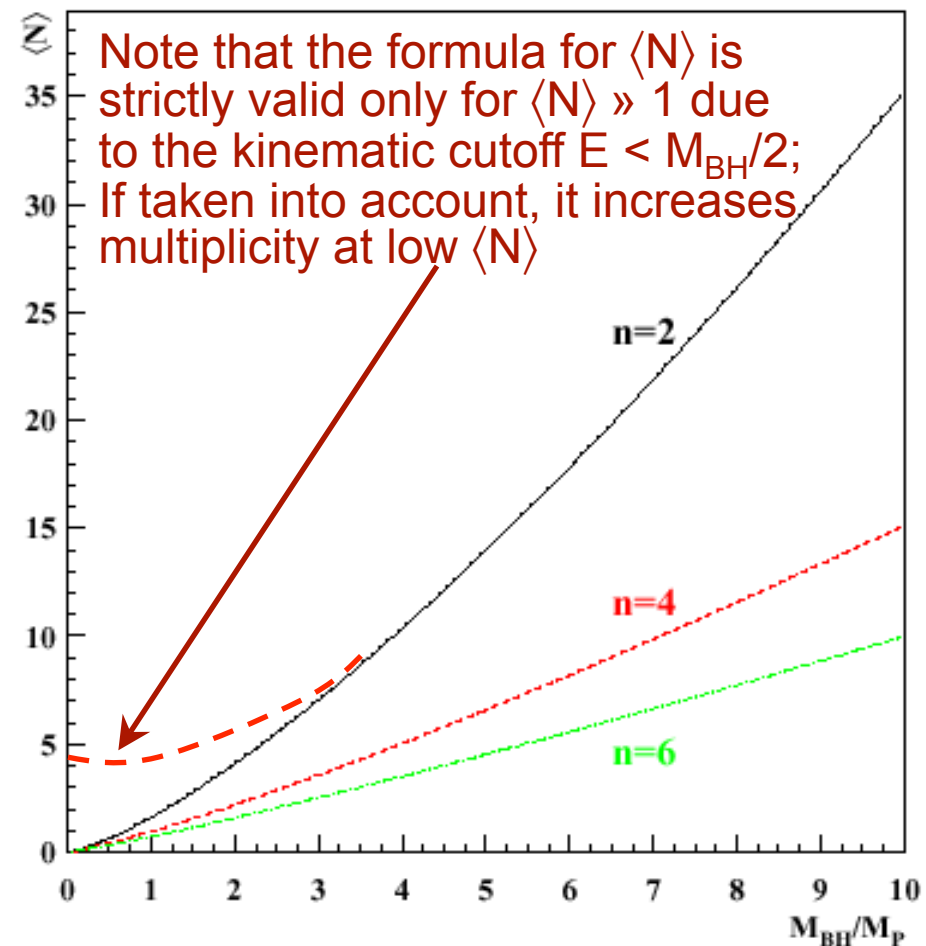


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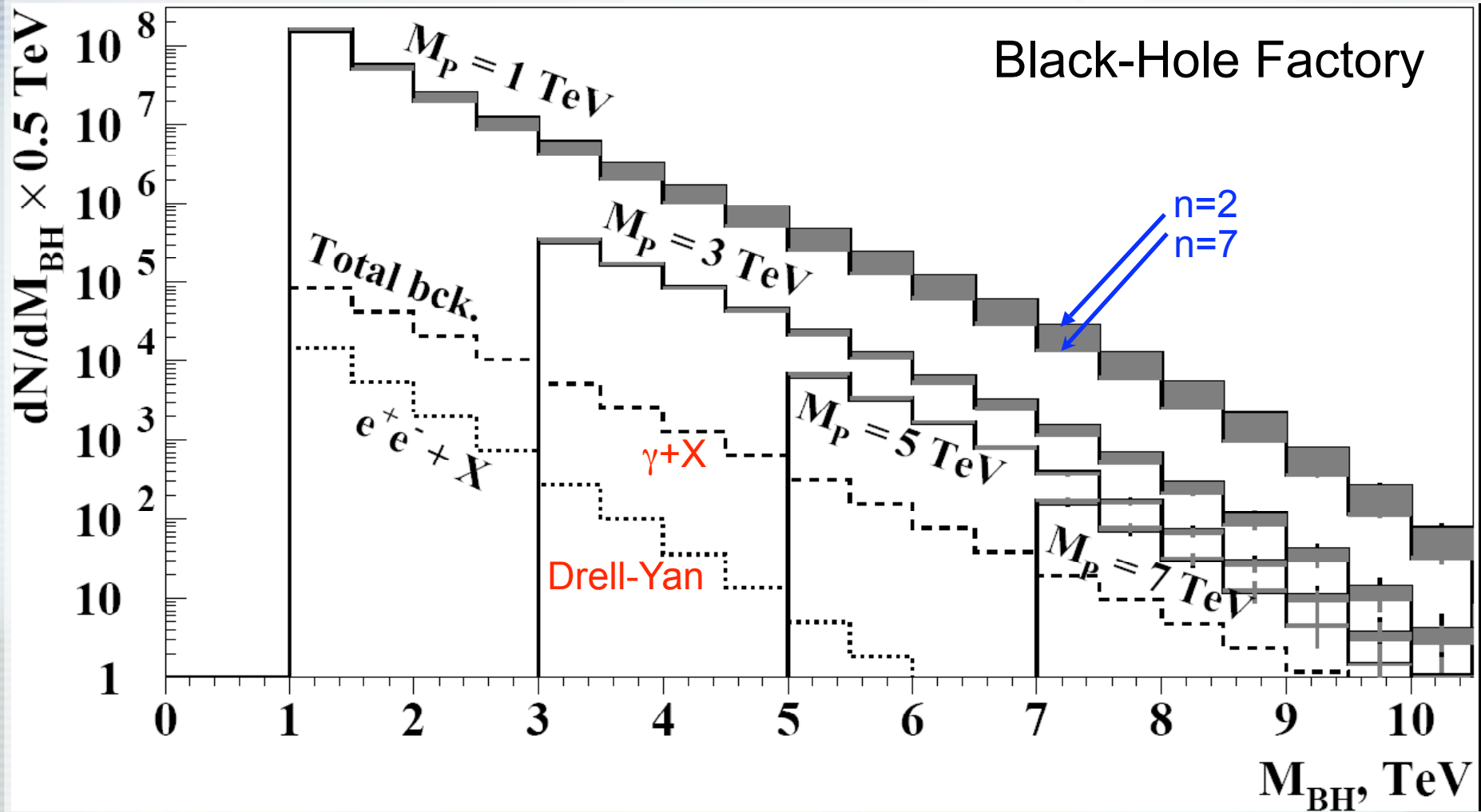


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# Black Hole Factory

Dimopoulos, GL [PRL **87**, 161602 (2001)]



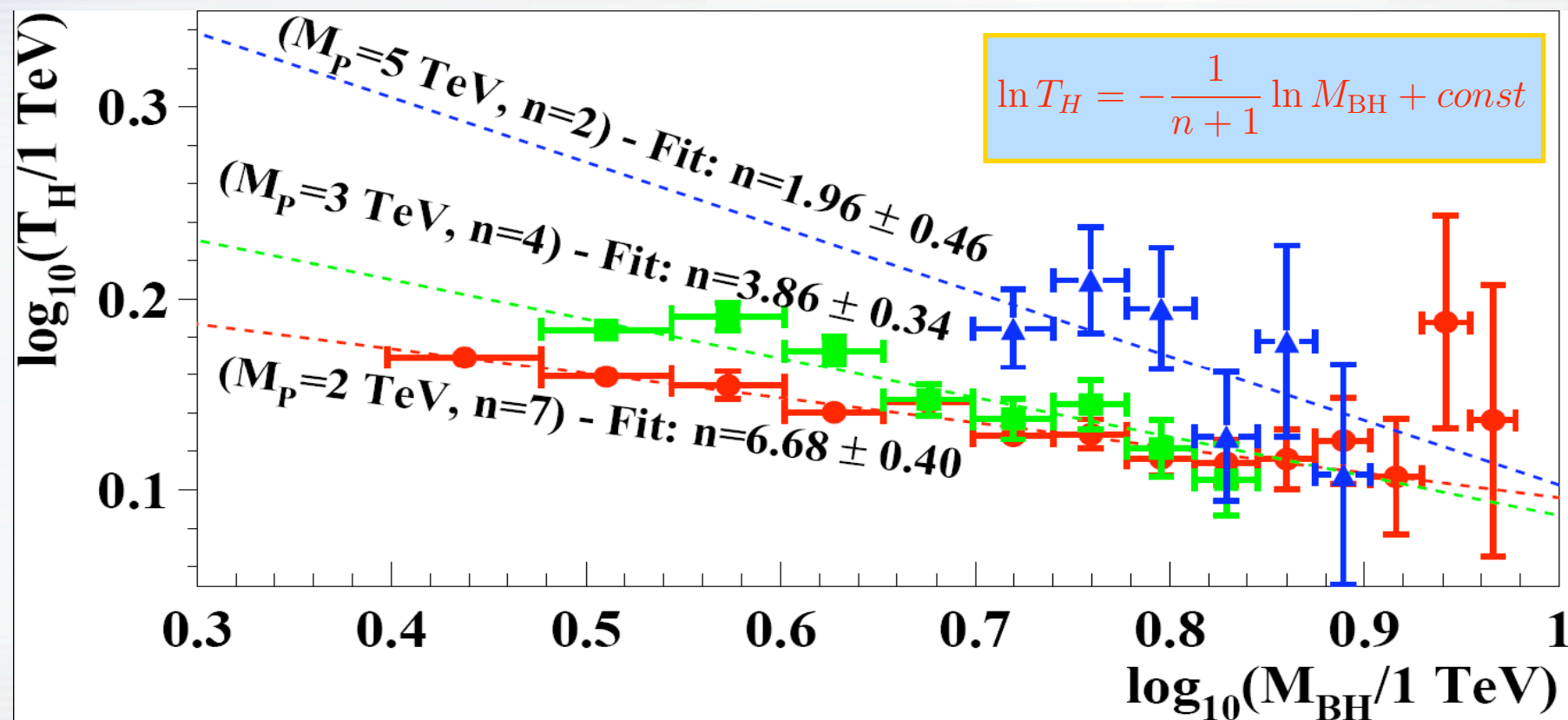
Spectrum of BH produced at the LHC with subsequent decay into final states tagged with an electron or a photon





# Shape of Gravity at the LHC

Dimopoulos, GL [PRL **87**, 161602 (2001)]



- Relationship between  $\log T_H$  and  $\log M_{BH}$  allows to find the number of ED
  - This result is independent of their shape!
  - This approach drastically differs from analyzing other collider signatures and would constitute a “smoking cannon” signature for a TeV Planck scale



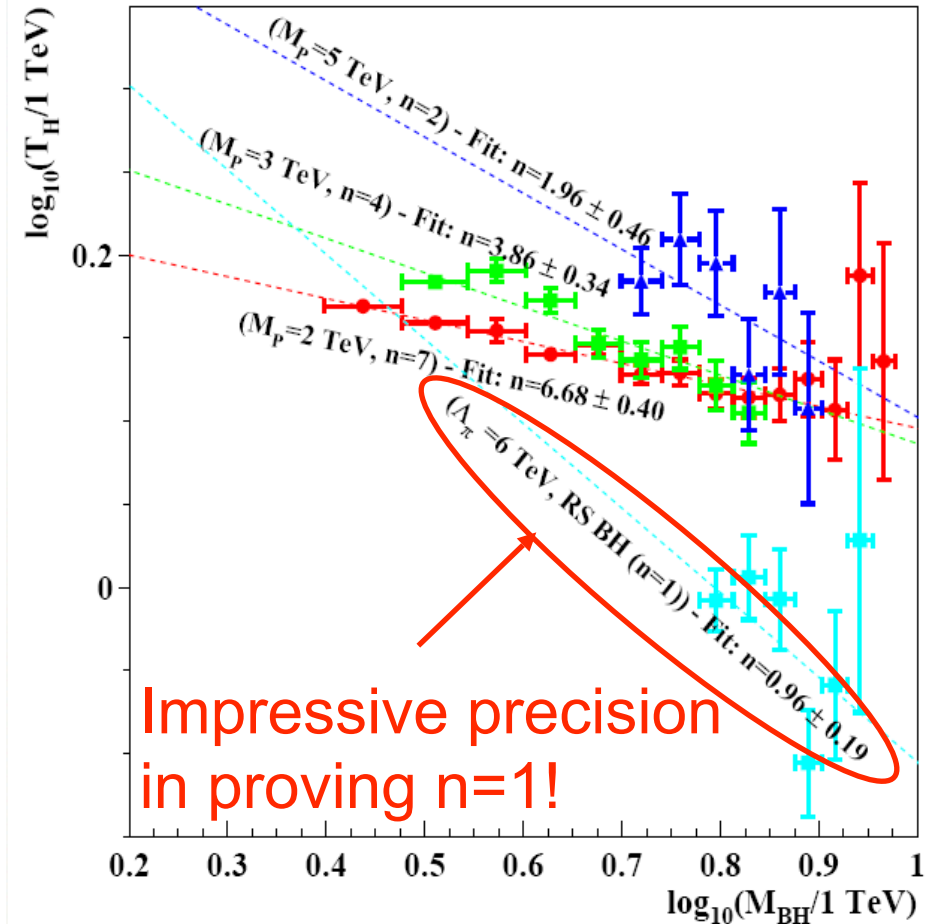
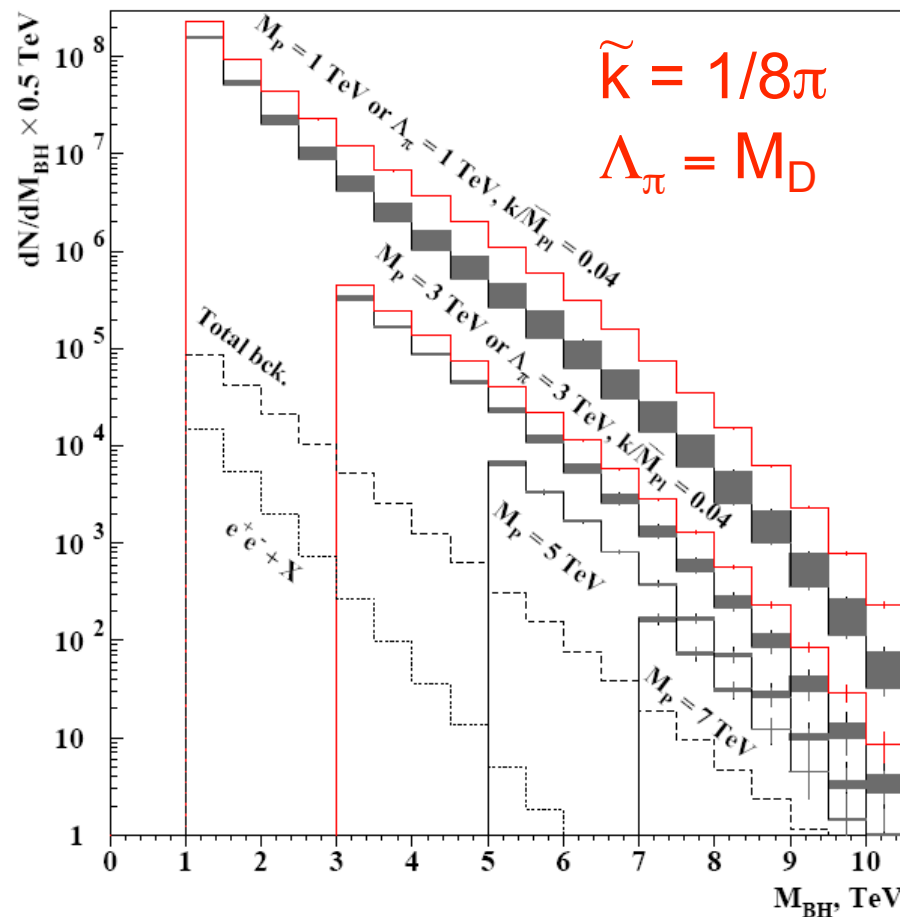
# Randall-Sundrum Black Holes

- Not nearly as studied as BH in large ED
  - Originally suggested in [Anchordoqui, Goldberg, Shapere \[PRD 66, 024033 \(2002\)\]](#)
  - A few authors extended work to various cases: [Rizzo \[JHEP 0501, 28 \(2005\); hep-ph/0510420; hep-ph/0603242\]](#); [Stojkovic \[PRL 94, 011603 \(2005\)\]](#)
  - The event horizon has a **pancake-like shape** (squashed in the 5<sup>th</sup> dimension **by  $e^{-k\pi r}$** )
- Nevertheless, the **comparison with the ADD BH is trivial**, [GL \[J. Phys. G32, R337 \(2006\)\]](#)
  - If  $R_s e^{-k\pi r} \ll \pi r$  the BH is still “small” and **can be treated as a 5D BH in flat space** (ignoring the AdS curvature at the SM brane  $\sim k^2 \ll 1$ )
  - For BH production,  $\Lambda_\pi$  in the RS model plays the same role as the **fundamental Planck scale  $M_D$  in the ADD model**
  - Recent paper by [Meade/Randall \[arXiv:0708.3017\]](#) used a different characteristic scale:  $\overline{M}_{Pl} e^{-k\pi r}$ , which resulted in a more conservative cross section estimate



# RS BH: Samples & Wien's Law

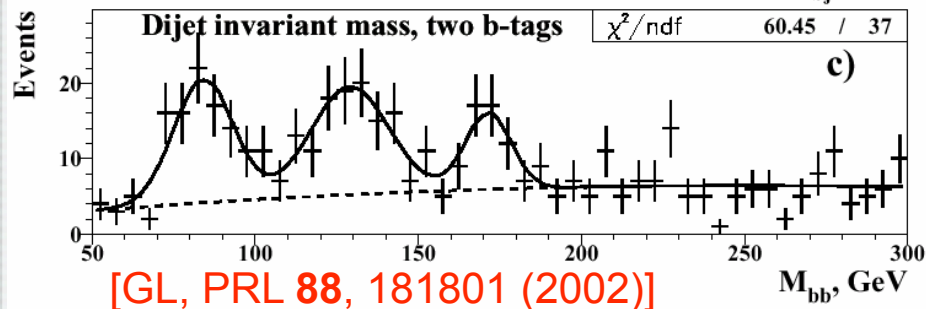
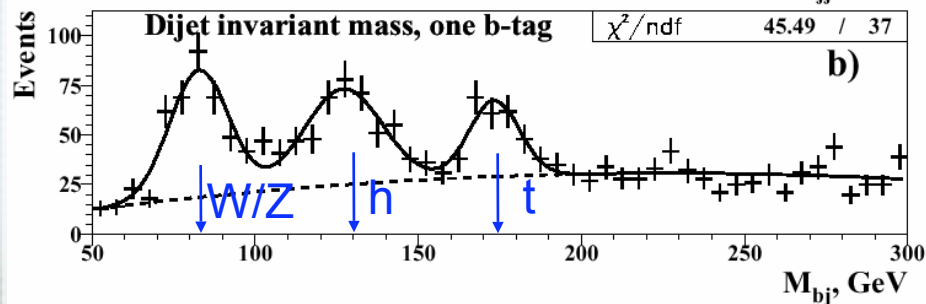
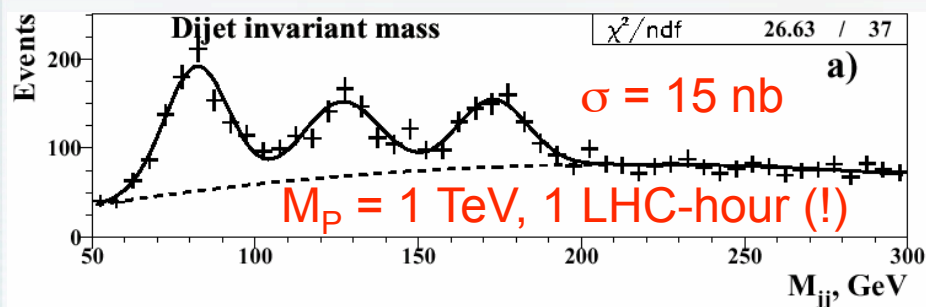
100 fb<sup>-1</sup> @ the LHC





# New Physics in BH Decays

- Example: Higgs with the mass of 130 GeV decays predominantly into  $b\bar{b}$ 
  - Tag BH events with leptons or photons, and look at the dijet invariant mass; does not even require b-tagging!
- Use typical LHC detector response to obtain realistic results

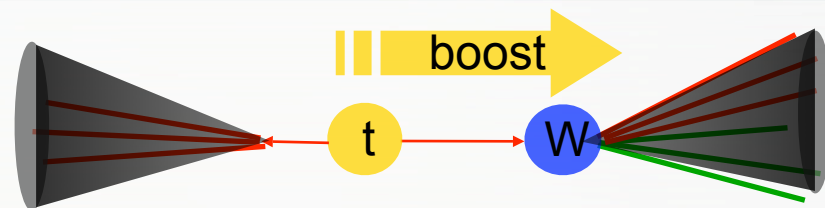
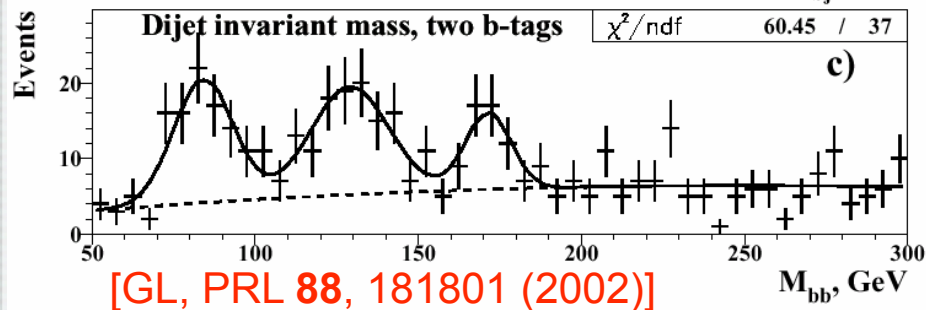
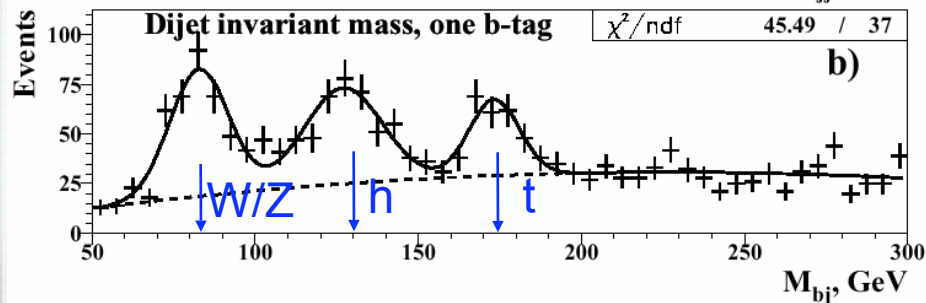
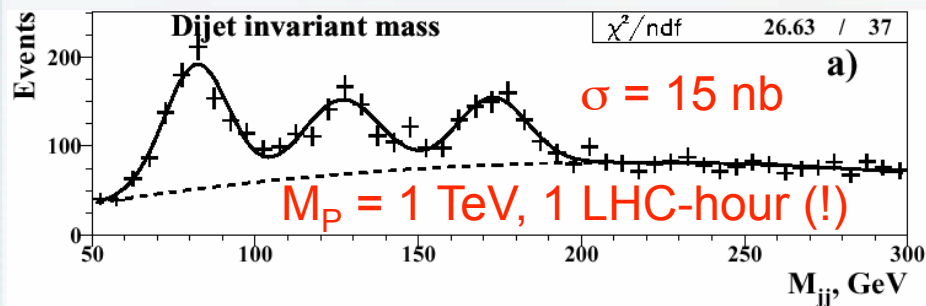






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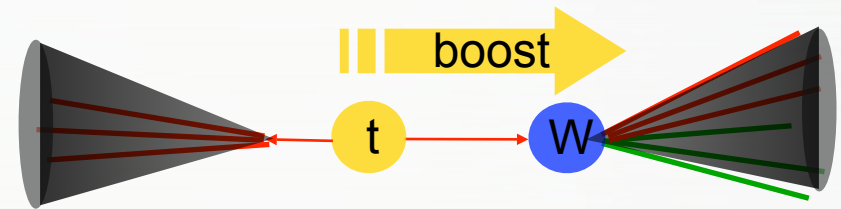
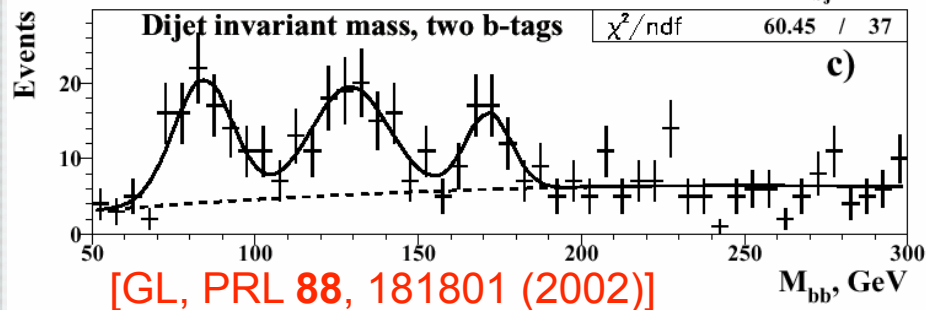
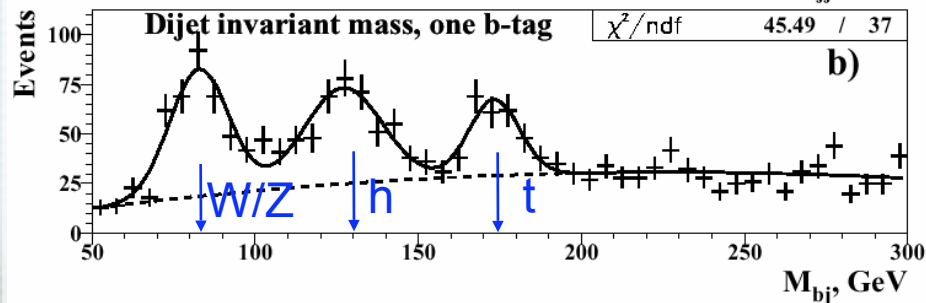
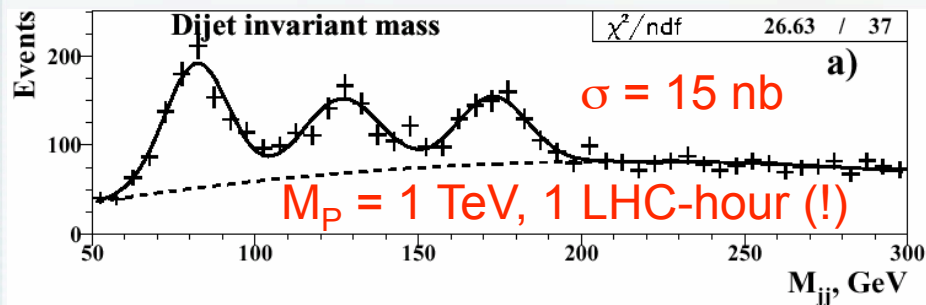
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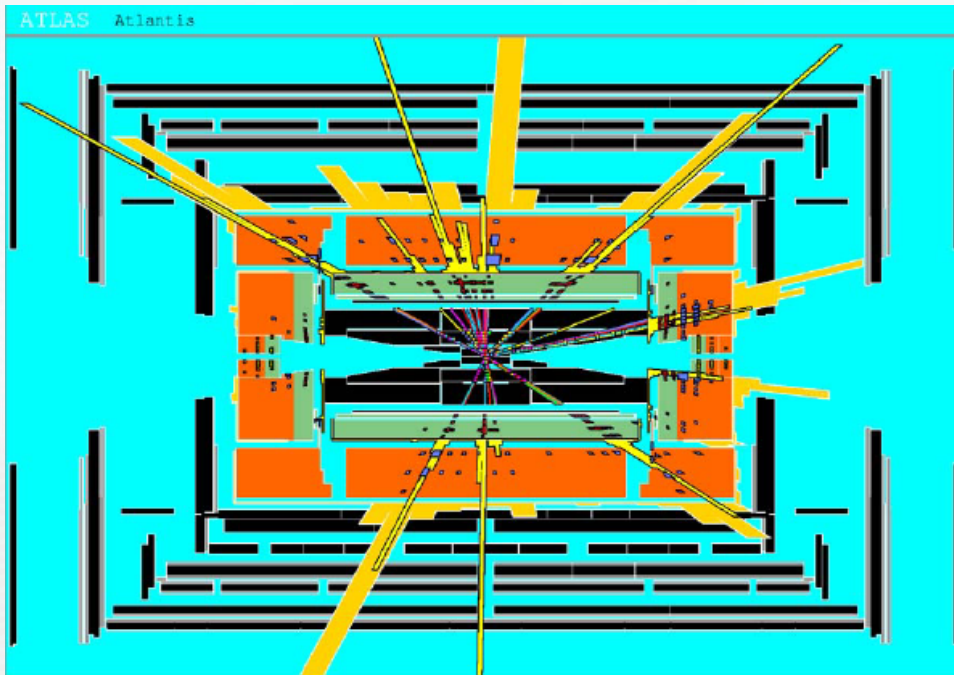


- Higgs observation in the black hole decays is possible at the LHC as early as in the first day of running even with the incomplete and poorly calibrated detectors!
- For  $M_P = 1, 2, 3,$  and  $4$  TeV one needs 1 day, 1 week, 1 month, or 1 year of running to find a  $5\sigma$  signal
- Higgs is just an example – this applies to most of the new particles with the mass  $\sim 100$  GeV

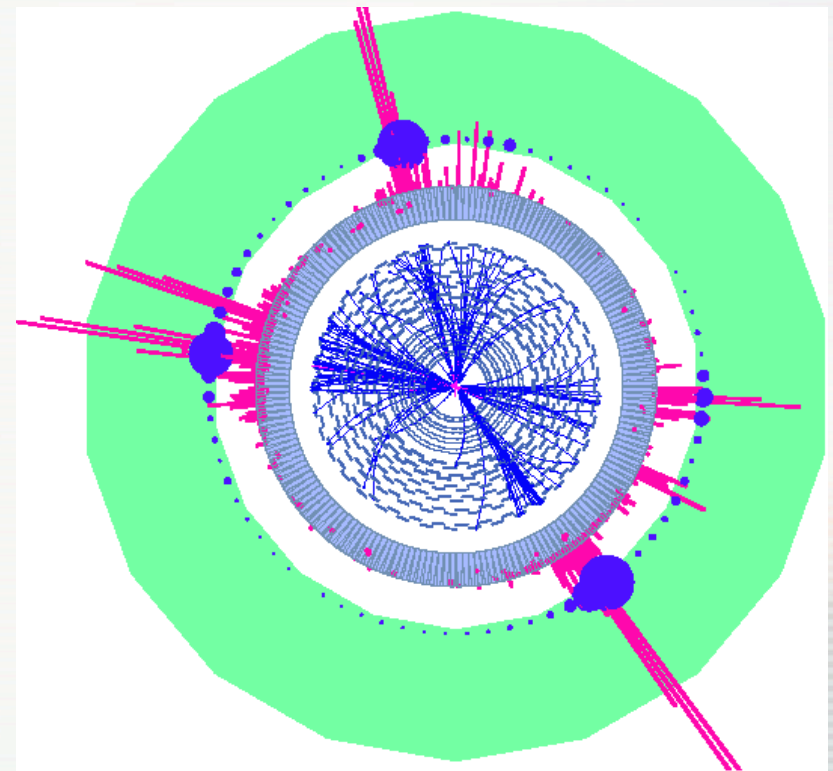


# Black Hole Events

- Detailed studies already started in ATLAS and CMS
  - ATLAS – CHARYBDIS (HERWIG-based generator with an elaborated decay model by Harris/Richardson/Webber)
  - CMS – TRUENOIR, GL/CHARYBDIS
- The hunt is going on!



Simulated black hole event in the ATLAS detector, from ATLAS-Japan Group

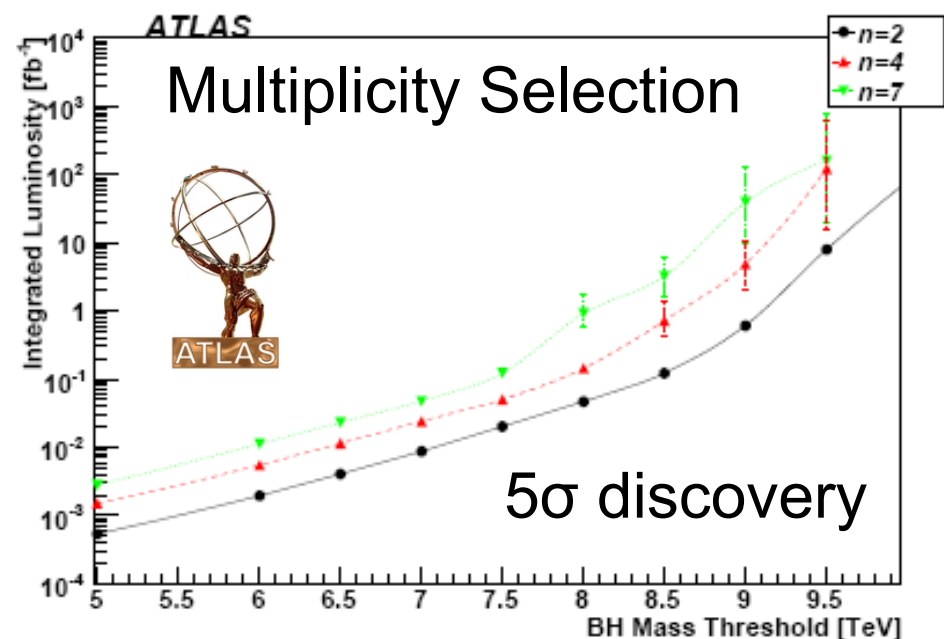
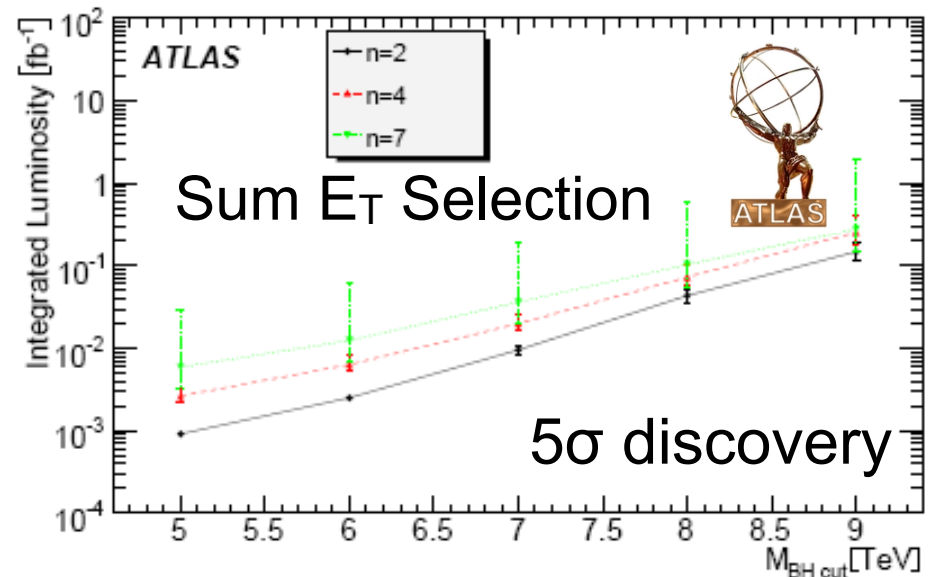


Simulated black hole event in the CMS detector, A. deRoeck & S. Wynhoff



# ATLAS Early Search for Black Holes

- Considered two selections:
  - Sum  $E_T$  selection:  $S_T > 2.5$  TeV,  $p_{T^l} > 50$  GeV
  - Multiplicity selection: require at least four energetic objects ( $p_T > 200$  GeV) in the final state; at least one is lepton
- For  $M_D$  of 1 TeV discovery of  $>5$  TeV BH's is possible with a fraction of  $\text{fb}^{-1}$







# Conclusions



# Conclusions

- Possibility of Extra Dimensions in space is a **bold theoretical idea**, which recently has acquired a new face:
  - Attempts to solve the hierarchy problem and other problems of the SM via an alternative framework



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- Such a possibility would fulfill our dreams for **Grand Unification of an ultimate kind**: that of particle physics, astrophysics, and cosmology!