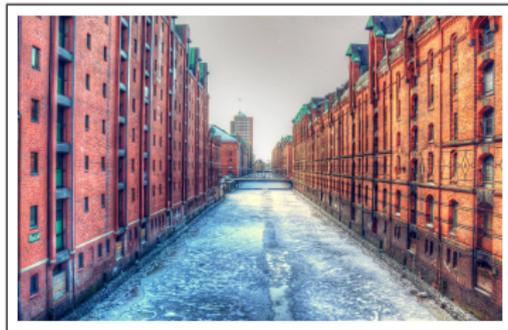


Spin structure of the proton at low x and low Q^2 in two-dimensional bins from COMPASS

Ana Sofia Nunes (LIP-Lisbon), on behalf of the COMPASS Collaboration



DIS2016, DESY, Hamburg, Germany – April 11th-15th, 2016



Aknowledgements:

FCT Fundação para a Ciência e a Tecnologia
MINISTÉRIO DA CIÊNCIA, TECNOLOGIA E ENSINO SUPERIOR

CERN/FIS-NUC/0017/2015

Motivation

- low $x \leftrightarrow$ high parton densities
- low x and low $Q^2 \leftrightarrow$ **transition from the regime of photoproduction to the regime of DIS (described by pQCD)**
- A_1^P and g_1^P as functions of x and ν showed, **for the first time, positive spin effects at very low x** (cf. $A_1^d \sim 0$ at low x , and SMC sample - 150x smaller)
- theoretical predictions for g_1^P as function of two kinematic variables:
 - ▶ Badełek et al., Eur.Phys.J. C26 (2002) 45
“Spin structure function $g_1(x, Q^2)$ and the DHGMY integral $I(Q^2)$ at low Q^2 : Predictions from the GVM model”
 - ▶ Ermolaev et al., Eur.Phys.J. C58 (2008) 29
“Comment on the recent COMPASS data on the spin structure function g_1 ”
 - ▶ Ermolaev et al., Riv.Nuovo Cim. 33 (2010) 57
“Overview of the spin structure function g_1 at arbitrary x and Q^2 ”
“one can parameterize g_1 by the set of variables x, Q^2 or, alternatively, $\omega[\equiv 2pq = 2M(E - E')]$, Q^2 , or ν, Q^2 ”
- COMPASS' $\sim 7 \times 10^8$ events allow a 2D extraction
- extraction, for the first time, in 4 2D grids: (x, Q^2) , (ν, Q^2) , (x, ν) , (Q^2, x)

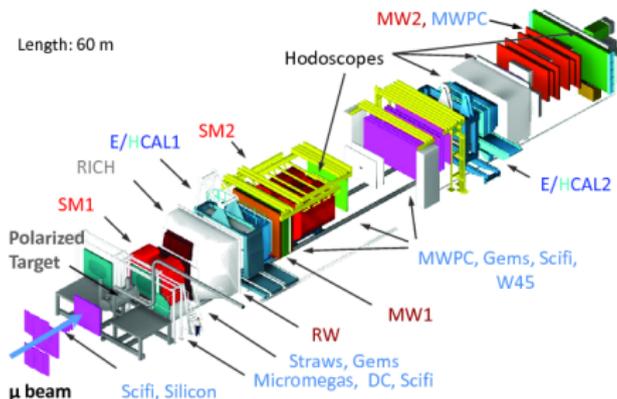
The COMPASS experiment at CERN

COMPASS @ CERN

COmmun **M**uon **P**roton
Apparatus for **S**tructure
and **S**pectroscopy

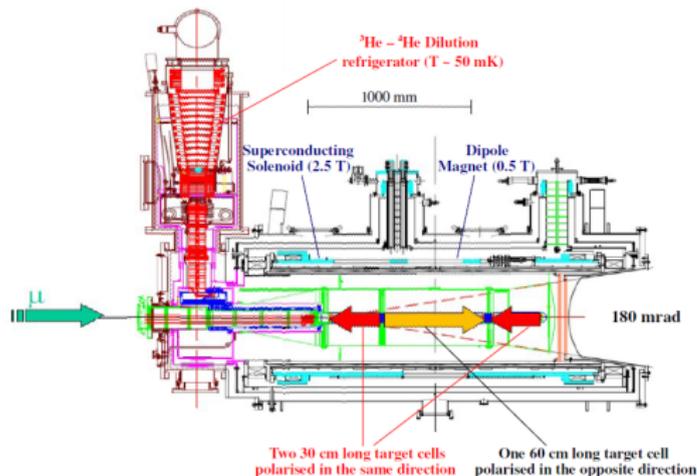


- **Fixed target experiment** at the SPS using a tertiary **muon beam**
- Collaboration of about 200 members from 11 countries and 23 institutions



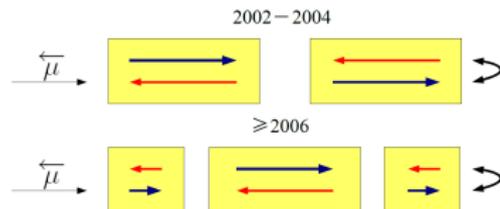
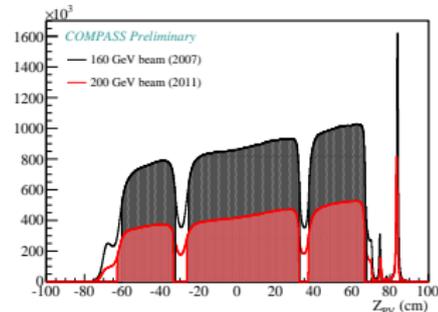
- 160/200 GeV μ^+ **polarised beam**, $P_b \sim -80\%$
- ${}^6\text{LiD}$ or NH_3 , 1.2 m long, **polarised target** @ 2.5 T and 60 mK, $P_{\text{target}} \sim 50/85\%$
- large acceptance, two staged spectrometer
- tracking, calorimetry, PID

Polarised target



Material	Dilution factor (f)	Polarisation (P_{target})
⁶ LiD	0.40	50%
NH ₃	0.16	85%

Vertex coordinate z_{PV}



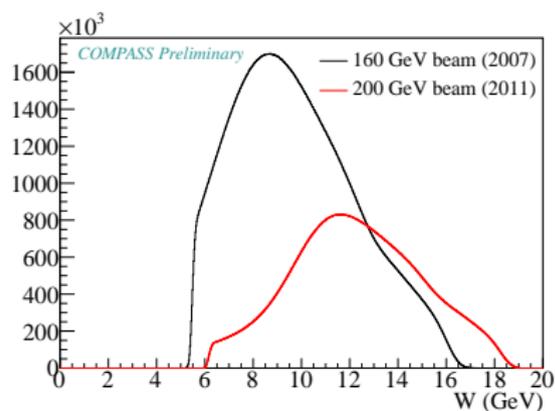
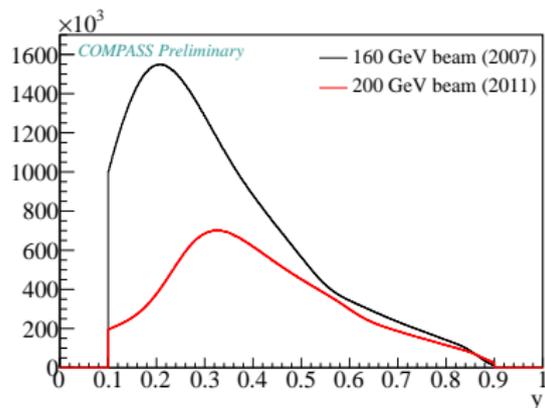
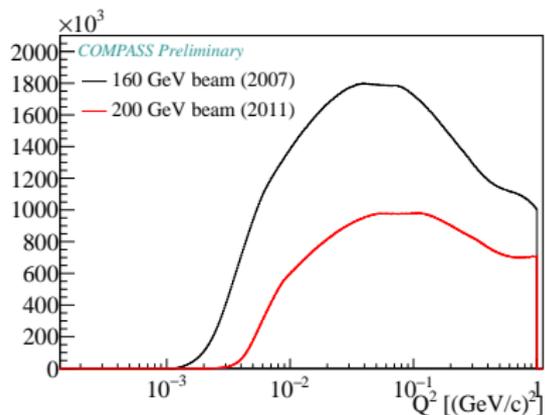
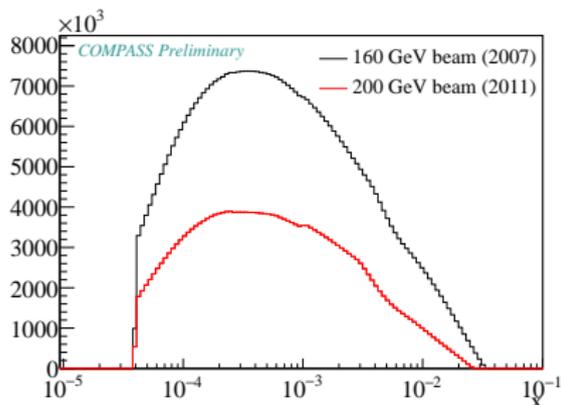
Data samples for the extraction of A_1^p and g_1^p

- Longitudinally polarised target (NH_3): **676×10^6 events**
(447×10^6 with 160 GeV beam in 2007, 229×10^6 with 200 GeV beam in 2011)
- Before, SMC low x , low Q^2 proton data: 4.5×10^6 events
 \Rightarrow The COMPASS data set has **150** \times more events than SMC

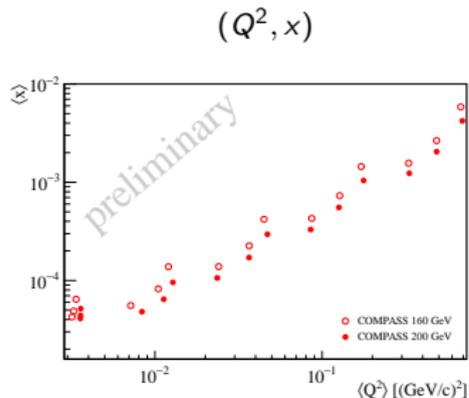
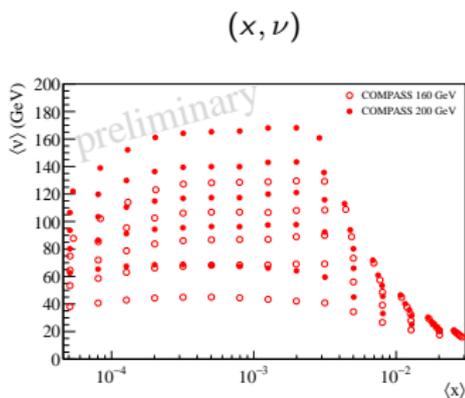
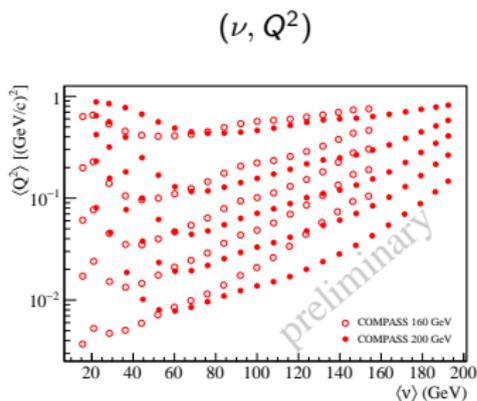
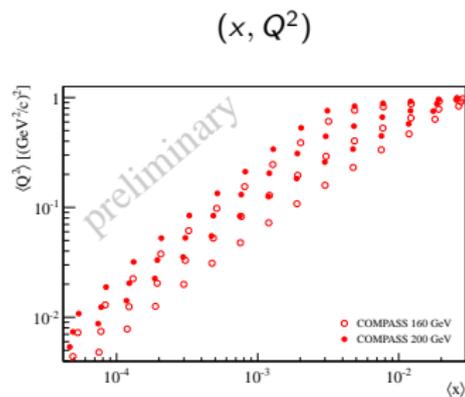
Main selection criteria:

- at least one additional track (besides the scattered muon) in the interaction point ("hadron method") - SMC proved there is no bias to the inclusive asymmetries at low x
- not a μe elastic scattering event
- $Q^2 < 1 \text{ (GeV}/c)^2$
- $x \geq 4 \times 10^{-5}$
- $0.1 < y < 0.9$

Characteristics of the final sample



Phase-space coverage of the 2D analysis



Double spin longitudinal asymmetry A_1^P

- $N^{\rightarrow, \leftarrow} = a \phi n \bar{\sigma} (1 \pm \mathbf{P}_{\text{beam}} \mathbf{P}_{\text{target}} \mathbf{f} \mathbf{D} \mathbf{A}_1^P)$ $\frac{N^{\rightarrow,1} \cdot N^{\leftarrow,2}}{N^{\leftarrow,1} \cdot N^{\rightarrow,2}} \rightarrow 2^{\text{nd}}$ ord. eq. on A_1^P
- Each event is given a weight $\omega = \mathbf{f} \mathbf{D} |\mathbf{P}_{\text{beam}}|$ to optimize the statistical errors
- Unpolarised radiative corrections (RC), included in the dilution factor, from TERAD
[A.A. Akhundov, *et al.*, Fortschr. Phys. 44 (1996) 373]
- Polarised radiative corrections ($A^{\text{RC}} \leq 0.25 \delta A_1^{\text{stat}}$) from POLRAD
[I. Akushevich *et al.*, Comput.Phys.Commun. 104 (1997) 201]
- Corrected for polarisable ^{14}N ($A^{^{14}\text{N}} \leq 0.01 \delta A_1^{\text{stat}}$)
- Thorough checks on possible sources of false asymmetries \Rightarrow systematic errors similar to the statistical errors

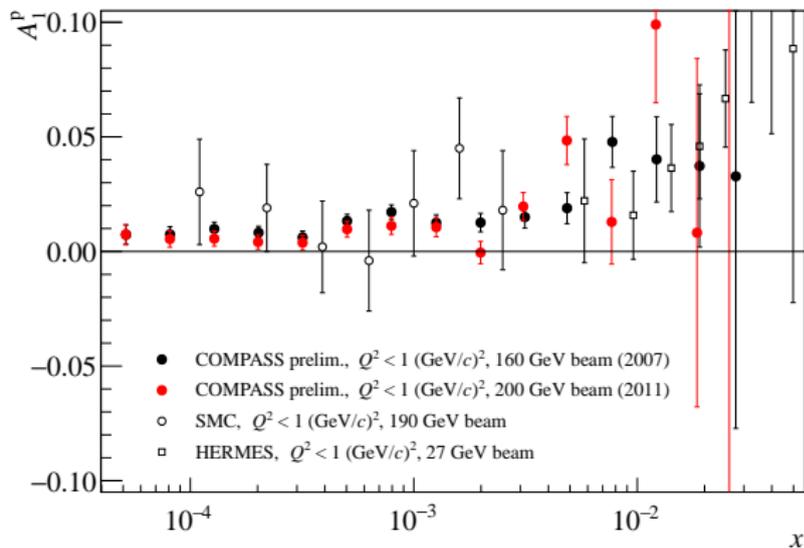
Spin dependent structure function g_1^p

- The structure function is obtained in bins of x or ν according to:

$$g_1^p(\langle x \rangle, \langle Q^2 \rangle) = \frac{F_2^p(\langle x \rangle, \langle Q^2 \rangle)}{2x[1 + R(\langle x \rangle, \langle Q^2 \rangle)]} A_1^p(\langle x \rangle, \langle Q^2 \rangle)$$

- $F_2^p(\langle x \rangle, \langle Q^2 \rangle)$ from the SMC fit on data or from a model (for low x and Q^2)
[SMC, Phys.Rev. D58 (1998), 112001; B. Badełek & J. Kwieciński, Phys.Lett. B295 (1992) 263]
- $R(\langle x \rangle, \langle Q^2 \rangle)$ based on SLAC parameterization, extended to low Q^2
[COMPASS, PLB 647 (2007) 330]

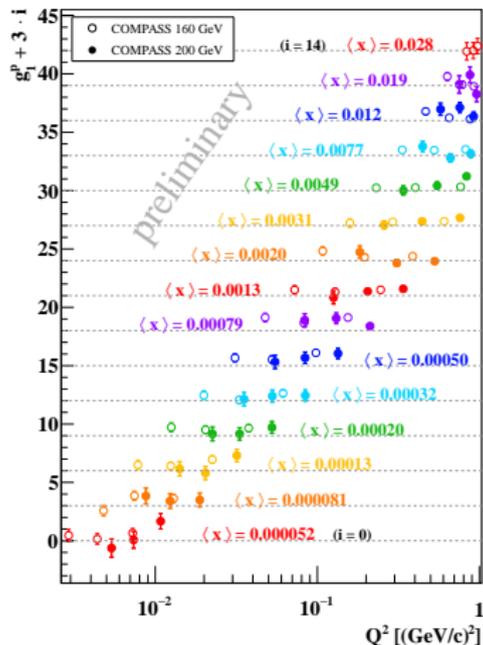
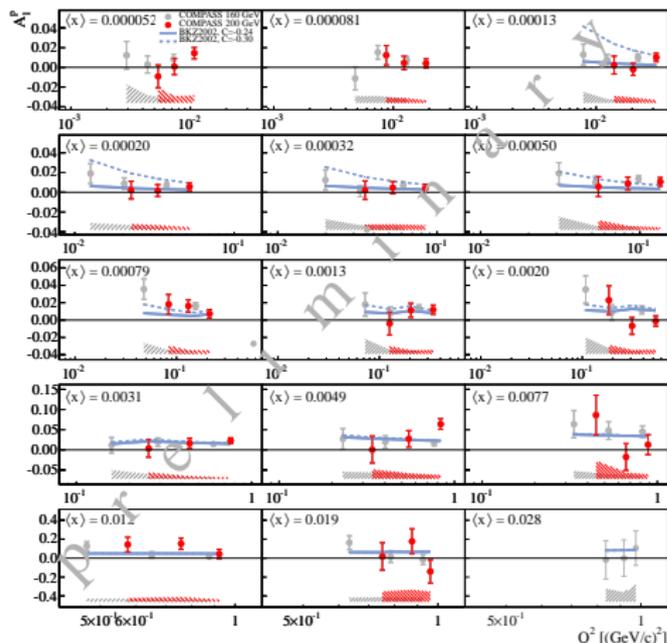
$A_1^P(x)$ & comparison with previous experiments



- results for the two beam energies are compatible within errors
- systematic errors are similar to the statistical errors (not shown here)
- A_1^P is **significantly positive**
- no dependence on x is seen (nor on ν , not shown here)
- the COMPASS results **improve the precision** of the measurement

A_1^P and g_1^P at low x and low Q^2 : results for the grid (x, Q^2)

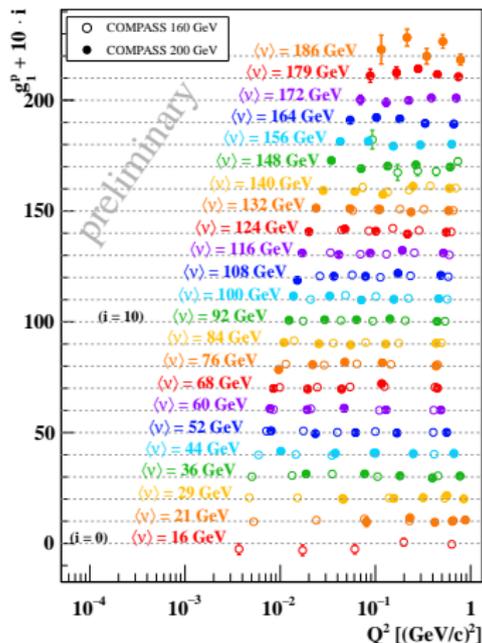
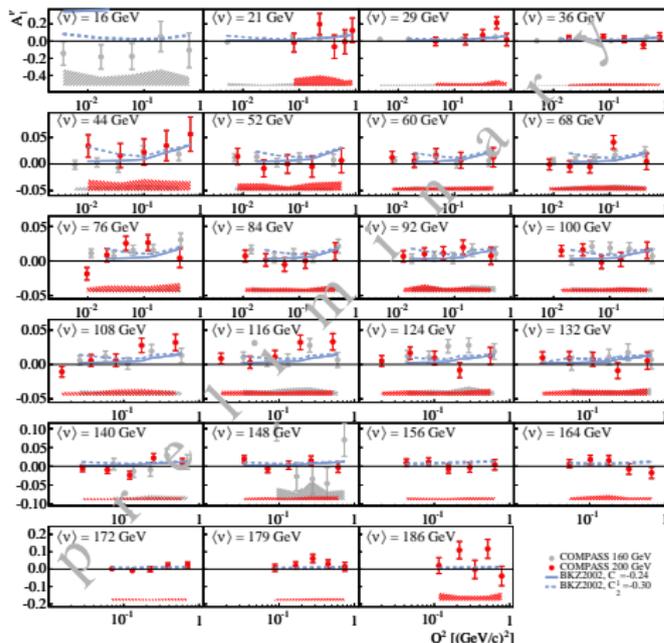
Data: 2007&2011, $\mu^+p \rightarrow \mu^+X$



- **no strong dependence** on x or Q^2
- results **compatible with theoretical model (GVMD)** [Eur.Phys.J. C26 (2002) 45]

A_1^P and g_1^P at low x and low Q^2 : results for the grid (ν , Q^2)

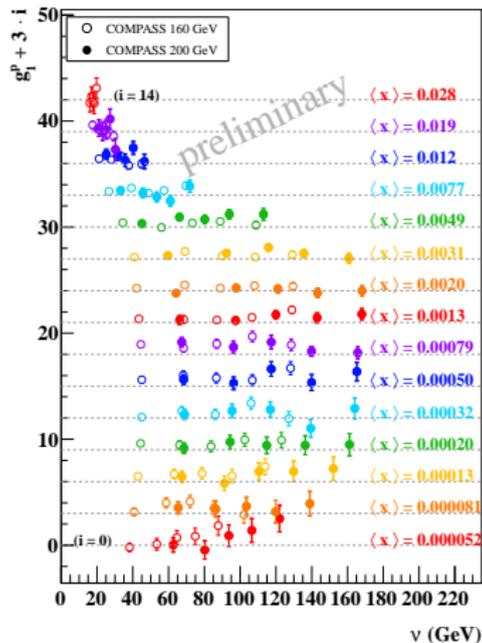
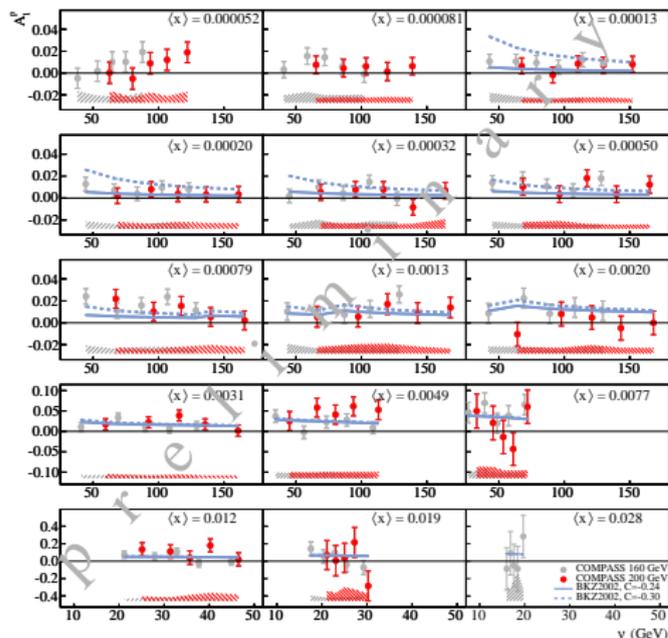
Data: 2007&2011, $\mu^+p \rightarrow \mu^+X$



- **no strong dependence** on ν or Q^2
- results **compatible with theoretical model (GVMD)** [Eur.Phys.J. C26 (2002) 45]

A_1^P and g_1^P at low x and low Q^2 : results for the grid (ν, x)

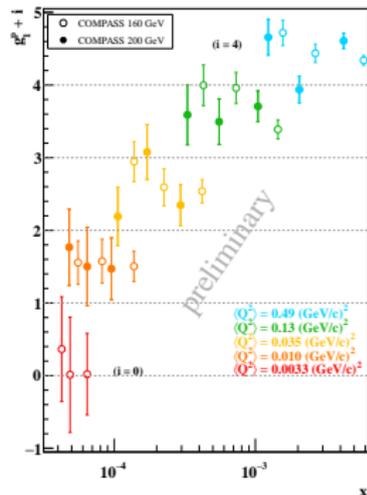
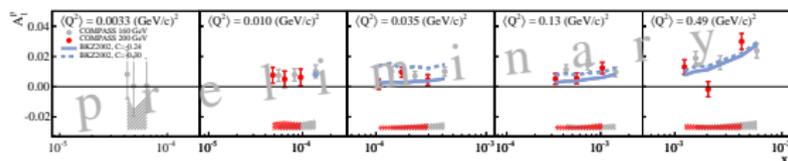
Data: 2007&2011, $\mu^+p \rightarrow \mu^+X$



- **no strong dependence** on ν or x
- results **compatible with theoretical model (GVM)** [Eur.Phys.J. C26 (2002) 45]

A_1^p and g_1^p at low x and low Q^2 : results for the grid (Q^2, x)

Data: 2007&2011, $\mu^+p \rightarrow \mu^+X$



- **no strong dependence** on x or Q^2
- results **compatible with theoretical model (GVMD)** [Eur.Phys.J. C26 (2002) 45]

Summary and outlook

- Longitudinal double spin asymmetries A_1^P and the spin dependent structure function g_1^P extracted in 4 two-dimensional grids:
 - ▶ (x, Q^2)
 - ▶ (ν, Q^2)
 - ▶ (x, Q^2)
 - ▶ (Q^2, x)
- **Positive** spin asymmetries at very low x
- **No significant dependence** on studied kinematic variables
- **Compatibility with GVMD model** predictions [Eur.Phys.J. C26 (2002) 45]

BACKUP

GVMD model [Eur.Phys.J. C26 (2002) 45]

[Badelek et al., Eur.Phys.J. C26 (2002) 45]

$$\begin{aligned}g_1(x, Q^2) &= g_1^\perp(x, Q^2) + g_1^{\text{AS}}(\bar{x}, Q^2 + Q_0^2) \\&= C \left[\frac{4}{9} (\Delta u_{\text{val}}^{(0)}(x) + \Delta \bar{u}^{(0)}(x)) \right. \\&\quad \left. + \frac{1}{9} (\Delta d_{\text{val}}^{(0)}(x) + \Delta \bar{d}^{(0)}(x)) \right] \frac{M_\rho^4}{(Q^2 + M_\rho^2)^2} \\&\quad + C \left[\frac{1}{9} (2\Delta \bar{s}^{(0)}(x)) \right] \frac{M_\phi^4}{(Q^2 + M_\phi^2)^2} \\&\quad + g_1^{\text{AS}}(\bar{x}, Q^2 + Q_0^2).\end{aligned}\tag{5}$$

To obtain the value of C from (12), the contribution of resonances was evaluated using the preliminary data taken at ELSA/MAMI by the GDH Collaboration [16] at the photoproduction, for $W_i = 1.8$ GeV. The asymptotic part of g_1 was parametrized using the GRSV2000 fit for the “standard scenario” of polarized parton distributions with a flavor symmetric light sea, $\Delta \bar{u} = \Delta \bar{d} = \Delta s = \Delta \bar{s}$, at the NLO accuracy [9]. The non-perturbative parton distributions, $\Delta p_j^{(0)}(x)$, in the light vector meson component of g_1 , (3), were evaluated at fixed $Q^2 = Q_0^2$, using, either

- (i) the GRSV2000 fit, or
- (ii) a simple, “flat” input:

$$\Delta p_i^{(0)}(x) = N_i(1-x)^{\eta_i},\tag{13}$$

with $\eta_{u_v} = \eta_{d_v} = 3$, $\eta_u = \eta_s = 7$ and $\eta_g = 5$. The normalization constants N_i were determined by imposing the Bjorken sum rule for $\Delta u_v^{(0)} - \Delta d_v^{(0)}$, and requiring that the first moments of all other distributions are the same as those determined from the QCD analysis [18]. It was checked that the parametrization (13) combined with the unified equations gives a reasonable description of the SMC data on $g_1^{\text{NS}}(x, Q^2)$ [19] and on $g_1^p(x, Q^2)$ [5]. This fit was also used to investigate the magnitude of the double logarithmic corrections, $\ln^2(1/x)$, to the spin structure function of the proton at low x [20]. We have assumed $Q_0^2 = 1.2$ GeV², cf. (1) and (3), in accordance with the analysis of F_2 [7, 8]. As a result the constant C was found to be -0.30 in case (i) and -0.24 in case (ii). These values change at most by 13% when Q_0^2 changes in the interval $1.0 < Q_0^2 < 1.6$ GeV².

GVMD model predictions [Eur.Phys.J. C26 (2002) 45]

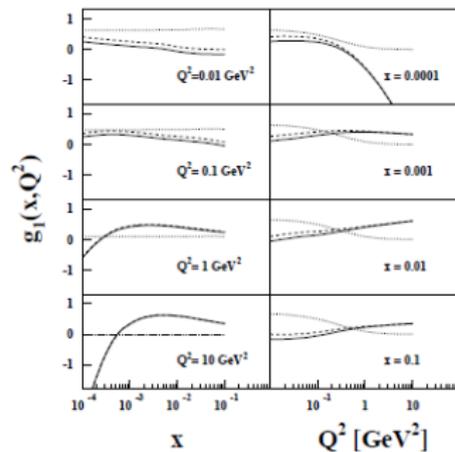


Fig. 1. Values of g_1 for the proton as a function of x and Q^2 . The asymptotic contribution, g_1^{ASYM} , is marked with broken lines, the VMD part, g_1^{VMD} , with dotted lines and the continuous curves mark their sum, according to (5)

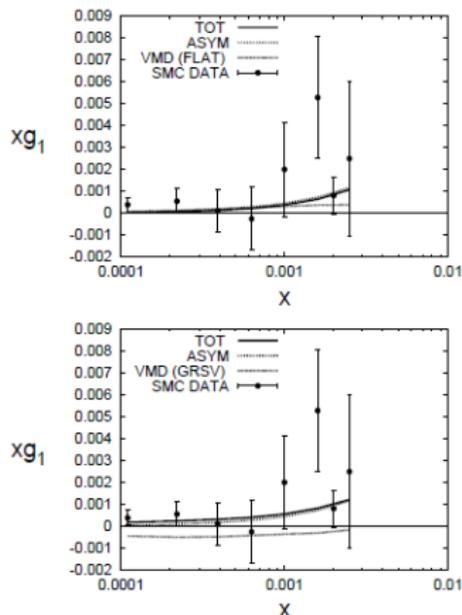


Fig. 2. Values of xg_1 for the proton as a function of x at the measured values of Q^2 in the non-resonant region, $x < x_t = Q^2/2Mv_t(Q^2)$. The upper plot corresponds to the VMD part parametrized using (13), the lower plot corresponds to the GRSV parametrization [9] of the VMD input. The g_1^{ASYM} in both plots has been calculated using the GRSV fit for standard scenario at the NLO accuracy. The contributions of the VMD and of the xg_1^{ASYM} are shown separately. Points are the SMC measurements at $Q^2 < 1 \text{ GeV}^2$ [3]; errors are total. The curves have been calculated at the measured x and Q^2 values