# Reconstruction of $\beta$ : analysis of $n=1.03, n=1.05$ and $\mathrm{n}=1.33$ radiators 

## C. Delgado

- Setup and runs description.
- PMTs calibration and alignment procedures.
- Event selection.
- Number of hits.
- $\beta$ reconstruction algorithms.
- Reconstruction results.
- Towards data/MC agreement for aerogel.
- Conclusions/Petitions.


## Setup description

- Wire chambers and scintillators planes on top of radiator.
- PMT array geometry:

341 mm



## Runs descriptions

The trigger rate is $\sim 0.5 \mathrm{~Hz}$
For the MC the simulated spectrum is muons at sea level.

- Radiator: 3 cm of aerogel $\mathrm{n}=1.03$
- Data: run21: ~200000 evts
- MC: 100 Kevts

Clarity: $0.0041 \mathrm{~cm}^{-1} \mu \mathrm{~m}^{4}$

- Radiator: 2 cm of aerogel $\mathrm{n}=1.05$
- Data: run 26: ~150000 evts
- MC: 100 Kevts

Clarity: $0.0091 \mathrm{~cm}^{-1} \mu \mathrm{~m}^{4}$
n simulated: 1.0488

- Radiator: 0.5 cm of $\mathrm{NaF} \mathrm{n}=1.33$
(10×10 $\mathrm{cm}^{2}$ tile)
- Data: run29: ~700000 evts
- MC: 200 Kevts
- All the aerogels with polyester supporting foil.
- NaF without supporting foil.
- All the runs aligned.
- Common PMT calibration for all the runs with noisy channel suppression for each run.


## PMTs calibration



- Gain: $\approx S(1+\lambda)$
- Single P.E. resolution: $\approx S \sqrt{\lambda}$
- Mean number of p.e.: $\mu$

Fit to the convolution of several distributions:

1. Pedestal: Gaussian.
2. Single p.e.: Described by the approximate distribution

$$
P(x) \simeq e^{-\lambda} \delta(x)+\left(1-e^{-\lambda}\right) \frac{e^{-\lambda}}{S \lambda} \frac{\lambda^{\frac{x}{S}}}{\Gamma\left(\frac{x}{S}\right)} \theta(x)
$$

The calibration resulted in

- Mean pedestal width: $\sim 4$ ADC counts.
- Mean gain $(\times 5): \sim 67$ ADC counts.
- $\sim 1 \%$ of negative pedestals
- $\sim 9 \%$ of double peaked pedestals


- Pedestals stability: in 19 hours is in the range $\pm 1 \mathrm{ADC}$ counts.
- PMTs dark current: yields $\sim 4 \times 10^{-5}$ hits per event per channel.
- Pedestal tails $(>4 \sigma)$ : yields $\sim 8 \times 10^{-5}$ hits per event per channel.




## Wire Chamber alignment

## Procedure:

- Wire chamber signals are fitted to a line.
- The residues for each plane are computed.
- The local position or all planes is shifted according to the peak of the residues associated to it.




## Wire Chambers/RICH alignment

- It relies in the RICH capacity to determine the light guide crossed by the particle.
- From MC: $\sigma_{x}^{R I C H} \simeq \sigma_{y}^{R I C H} \approx 0.4 \mathrm{~cm}$.
- It present a small bias if the tracks are not homogeneously distributed in $\phi$.



## Procedure

- Compute the extrapolation of the track to the PMT matrix.
- Within a given region around this point ( $\sim$ $1 / 2$ PMT size), choose the hit with largest number of p.e.
- If this number is large enough ( $\geq 3$ p.e.), compute:

1. Distance in $X$ to the track point.
2. Distance in $Y$ to the track point.
3. Difference in $\tan ^{-1} \frac{Y}{X}$ for the track point and the chosen one.




- All the aligned differences are compatible with zero.
- The $R I C H \oplus W C$ resolution is $\sim 0.6 \mathrm{~cm}$.


## Alignment stability



- From one run to other the alignment parameters change in the range $\pm 1 \mathrm{~mm}$.


## Track selection

- To ensure a well reconstructed track we proceed as in the alignment study, but we are more careful in choosing the signal threshold for the hit matching the track: signal>6 p.e.
- Only tracks with a hit with a signal above the threshold matching the track are selected.
- Strong matching criteria: Distance from track extrapolated point to selected hit $<1 \sigma_{R I C H \oplus W C}$.
- Finally a loose $\chi^{2}$ cut in the track is applied.




## Event selection

Further refinements are achieved by event selection:


- Events with possibility of confusion of the hit due to the crossing of the particle are flagged.


## Number of hits

- Very good agreement between data and MC for aerogel.
- For NaF there is an extra 30\% amount of hits respect MC.





## $\beta$ reconstruction algorithm

Two closely related algorithms implemented in the prototype reconstruction:

- Reconstruction with track parameters.
- Implemented in AMS software too.
- Very tested.
- Robust and fast as far as the tracks parameters are well known.
$\triangleright \approx 140$ times faster than Tracker reconstruction for protons.
- Reconstruction without track parameters.
- Not implemented int the AMS software.
- Not very tested (but working).
- $\approx 146$ times slower than the reconstruction with known track parameters


## Reconstruction with track parameters

1. Back trace: Find all the photons trajectories compatible with each hit and assign a $\beta$ to each trajectory.

- Semi analytical solving of propagation equation.
- Assume a common emission point for all the photons: mean emission point of the detected photons.

2. Fast search of most the probable common $\beta$ for all the reconstructed trajectories with noise reduction and ambiguity solving:

- Look for the cluster of $\beta$ values such that:
- For each hit, only the trajectory with the associated $\beta$ closer to the cluster center is considered.
- Only the $\beta$ closer than $3 \times \sigma(\beta$ hit $)$ to the cluster center are considered.
- Only the cluster with the bigger number of $\beta$ values is retained.
- The reconstructed particle $\beta$ is the mean value of the cluster.


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Reconstruction without track parameters
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1. For a set of points over the radiator estimates the particle direction assuming that it passes by this point.

- Currently it assumes that the particle passes through the light guide associated with the channel with bigger number of p.e.

2. If the estimated track is within the geometrical acceptance: reconstruct $\beta$.
3. Select the track with the best reconstructed $\beta$.

The uncertainty of the PMT behaviour when it is crossed by a charged particle is a source of uncertainty in the resolution for the current implementation of the algorithm.

## Track known: all $\beta_{h i t}$ spectrum

For aerogel there are two main differences between data and MC:

- The background in data is larger than in MC.
- The Čerekov peak width is larger in data than in MC.

This disagreement is larger for $\mathrm{n}=1.05$. Apparently this does not happen for NaF .



## Track known: reconstruction efficiency

- Good agreement with MC for 1.03: - 1.03 : $\quad$ Small disagreement
- Large disagreement for 1.05, to be investigated further.
- Disagreement in NaF compatible with excess of hits in the ring, could be due to geometry (still to be checked).



ext. track distance to center
ext. track distance to center


## Track known: hits in ring

- There is an important defect of hits associated to a ring for aerogel in data respect MC. The disagreement for $\mathrm{n}=1.05$ is quite worse than for $n=1.03$.
- For NaF this difference is a small excess of hits in the ring of data respect MC.




## Track known: resolution per hit

- Estimated as

$$
\lim _{n_{\text {used }} \rightarrow \infty} \sqrt{\frac{n_{\text {used }}}{n_{\text {used }}-1}} \sigma\left(\beta_{\text {event }}-\beta_{\text {hit }}\right)
$$

- $\approx 25 \%$ of disagreement for $n=1.03$
- Roughly $40 \%$ for $n=1.05$, but not conclusive due to the lack of hits.
- Very good agreement for NaF





## Track unknown: correlation with 'track known' rec.




- New reconstruction only working for aerogel runs.
- Reconstructions difference are within $\approx 2 \sigma$ of the reconstruction using the wire chambers for data and MC.


## Track unknown: reconstruction efficiency




- The reconstruction without wire chambers has a strong dependence in the fraction of ring detected. This explain the decrease for large $\beta$ in the efficiency.
- The reconstruction efficiency for $\mathrm{n}=1.05$ is smaller than the expected. This is due to the small mean number of hits per ring on data respect the MC.

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Track unknown: resolution per hit (%)
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- For $\mathrm{n}=1.03$, the disagreement between data and $M C$ is similar to the one of the reconstruction using the wire chambers.
- For $n=1.05$, apparently the disagreement is similar too. However the small number of hits per rings makes it difficult to give a precise number.


## Summary of differences

We can summarize the previous results as:

- For aerogel the background is larger than the expected.
This difference is such that:
- The total number of hits agrees in data and MC.
- The number of hits in the ring is smaller in data.
thus pointing to a migration from the population of hits in the ring to the background.
Apparently this does not happen for NaF.
- The resolution per hit for aerogel is worse than expected.


## Background excess




- The background and (partly) the difference in the resolution can be parametrized as a scattering process over imposed to the Rayleigh with:

$$
\mathrm{n}=1.03
$$

$$
\mathrm{n}=1.05
$$

$$
\begin{aligned}
& \text { - Mean free path: } 2 \mathrm{~cm} \\
& \text { - } \frac{d \sigma}{d \Omega} \propto \frac{G_{\sigma=0.5}(\theta)}{\sin \theta}
\end{aligned}
$$

$$
\begin{aligned}
& \text { - Mean free path: } 1 \mathrm{~cm} \\
& -\frac{d \sigma}{d \Omega} \propto \frac{G_{\sigma=0.5}(\theta)}{\sin \theta}
\end{aligned}
$$

The necessary scattering lengths ratio is compatible with the values obtained from the clarity:

$$
L_{\text {scattering }}(\lambda)=\lambda^{4} \text { Clarity }^{-1}
$$

- $\mathrm{n}=1.03(\mathrm{Cl}=0.0042): L_{\text {scattering }}(400 \mathrm{~nm}) \approx 6 \mathrm{~cm}$
- $\mathrm{n}=1.05(\mathrm{Cl}=0.0091): L_{\text {scattering }}(400 \mathrm{~nm}) \approx 3 \mathrm{~cm}$

However the cross section is different:

$$
\frac{d \sigma_{\text {rayleigh }}}{d \Omega} \propto 1+\cos ^{2} \theta
$$

The deterioration of the resolution is, at least partly, due to this effect.

- NaF resolution agrees with MC. So light guide angular inefficiency is well reproduced in MC. not clear.
- NaF photon yield is larger than expected. This to be understood.
- Agreement for $\mathrm{n}=1.03$ is better than for $\mathrm{n}=1.05$ (but note the acceptance and radiator thickness differences). However the source of both disagreements seems to be similar:
- Aparently it is not the wire chambers, as we see it using the independent algorithm.
- It is not the light guides or PMT, as we do not see it for NaF.
- The stacking cannot be responsible for all the effect as for $\mathrm{n}=1.05$ the stacking is smaller but the disagreement is larger.
- The disagreement could be due to a very forward scattering process.
- We have not idea of the contribution due to the polyester foil!!


## Petitions

- Set of measurements of the aerogels forward scattering angular distribution to determine if the new scattering process is really there.
- Set of runs without the polyester.
- Run of $\mathrm{n}=1.03$ (at least) without stacking.
- Run with $\mathrm{n}=1.05$ with 3 cm of thickness and/or smaller drift distance.
- Measurement of light guides properties.

